

Física do spin

F.S. Navarra

navarra@if.usp.br

Richard Terra (monitor)

richard.terra@usp.br

edisciplinas.if.usp.br

(buscar: física do spin)

Plano do curso

- 13/03 aula 1: Partículas elementares e idéias da física quântica
- 20/03 aula 2: Átomo de Bohr, quantização do momento angular
- 27/03 aula 3: Momento de dipolo magnético, Stern - Gerlach
- 10/04 aula 4: Efeito Zeeman anômalo
- 17/04 1ª Prova
- 24/04 aula 5: Equações de autovalores, matrizes de Pauli
- 08/05 aula 6: Comutadores, medidas SG sequenciais
- 15/05 aula 7: Medidas e valores médios
- 22/05 aula 8: Precessão e adição de spins
- 29/05 aula 9: Adição de spins

O momento angular é quantizado !

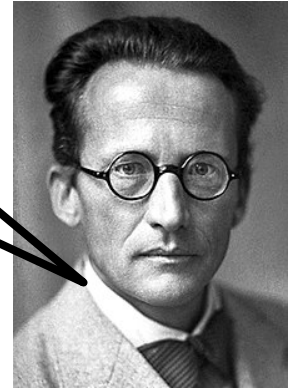


Eu chutei que era assim :

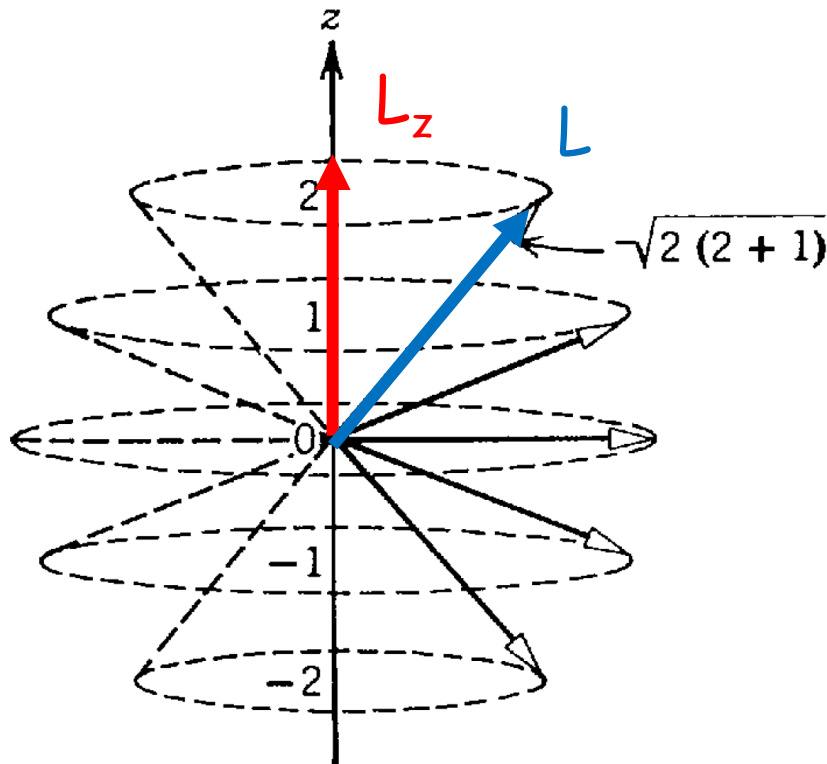
$$L = n \hbar$$

Bohr

Mas tava errado !
Certo é assim:



Schrödinger



$$L^2 = l(l + 1) \hbar^2$$

$$l = 0, 1, 2, \dots$$

$$L_z = m_l \hbar$$

$$-l \leq m_l \leq l$$

"Rotação quântica"

Aula 3

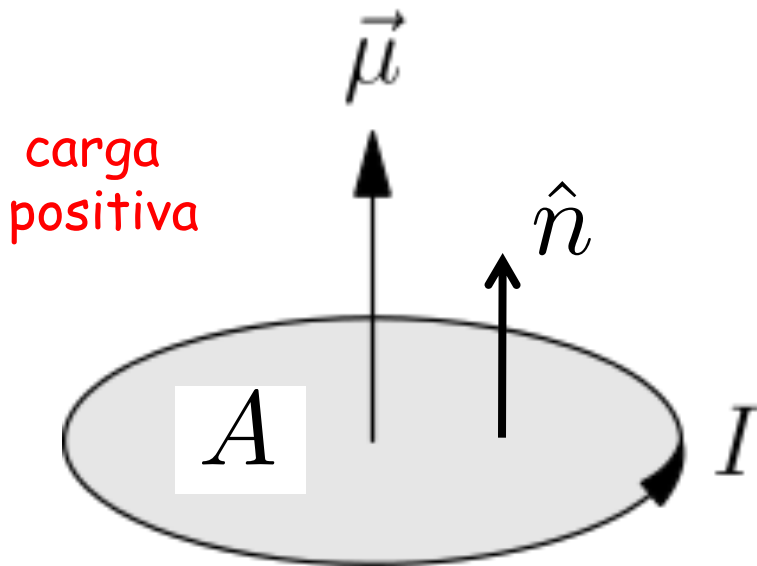
Momento de dipolo magnético

Espira num campo magnético

Experiência de Stern-Gerlach

Momento de dipolo magnético

Espira de corrente

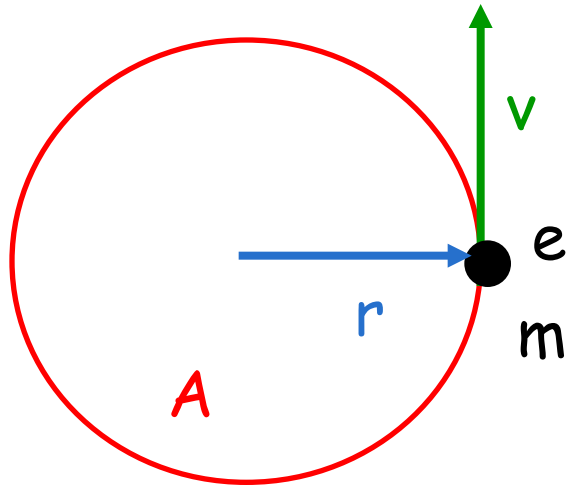


regra da
mão direita

$$\vec{A} = A \hat{n}$$

$$\vec{\mu} = I \vec{A}$$

Momento magnético e momento angular



corrente = carga / tempo

$$i = \frac{e}{T}$$

$$v = \frac{2\pi r}{T} \longrightarrow \frac{1}{T} = \frac{v}{2\pi r}$$

$$A = \pi r^2$$

$$i = \frac{ev}{2\pi r}$$

$$\left. \begin{aligned} \mu &= iA = \frac{ev}{2\pi r} \pi r^2 = \frac{evr}{2} \\ L &= mvr \end{aligned} \right\} \frac{\mu}{L} = \frac{e}{2m}$$

$$\frac{\mu}{L} = \frac{e}{2m} \frac{\hbar}{\hbar}$$

$$\mu_b = \frac{e \hbar}{2m}$$

magneton de Bohr

$$g_l = 1$$

fator g orbital

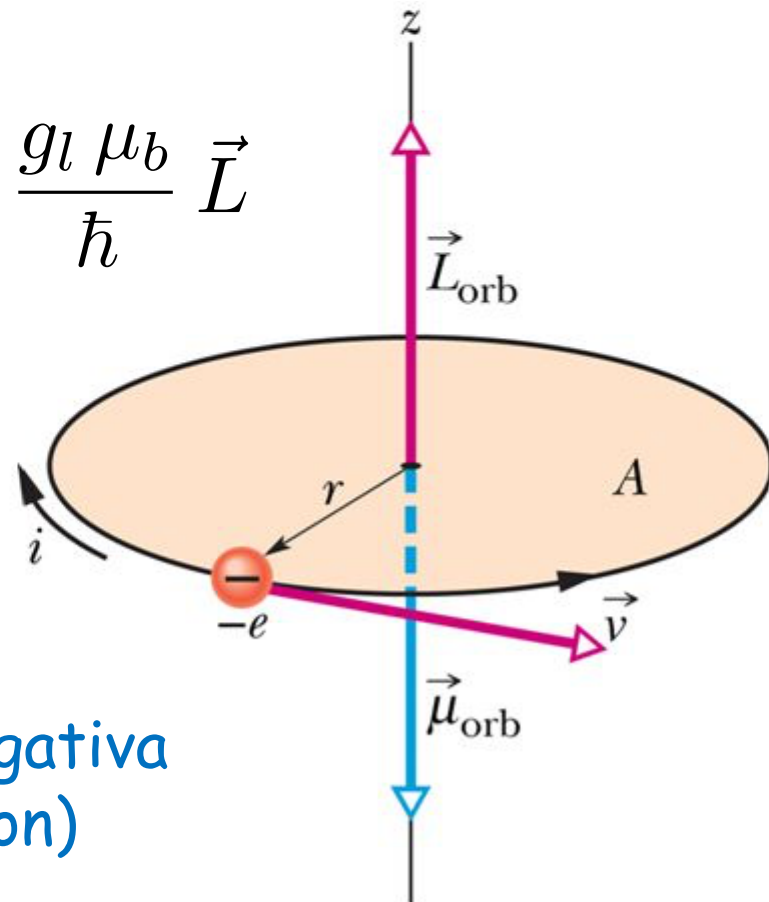
$$\frac{\mu}{L} = \frac{g_l \mu_b}{\hbar}$$

$$\vec{\mu} = - \frac{g_l \mu_b}{\hbar} \vec{L}$$

↓

$$\vec{\mu} = \frac{g_l \mu_b}{\hbar} \vec{L}$$

carga positiva



carga negativa
(eletron)

Momento magnético do elétron é quantizado !

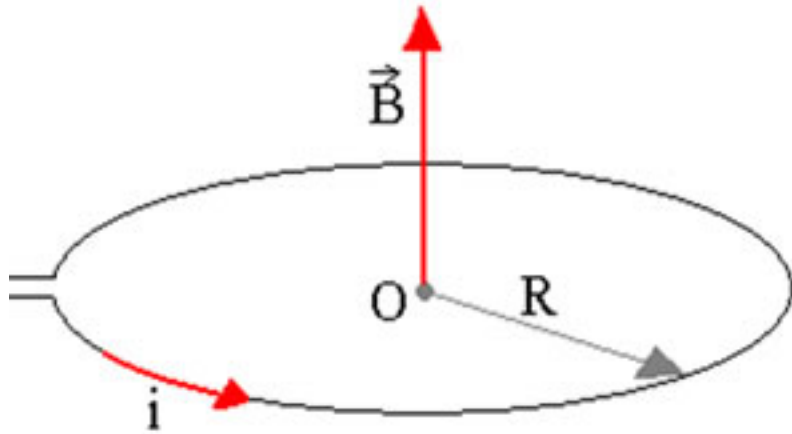
$$L = \sqrt{l(l+1)} \hbar$$

$$\mu = \frac{g_l \mu_b}{\hbar} L = \frac{g_l \mu_b}{\hbar} \sqrt{l(l+1)} \hbar = g_l \mu_b \sqrt{l(l+1)}$$

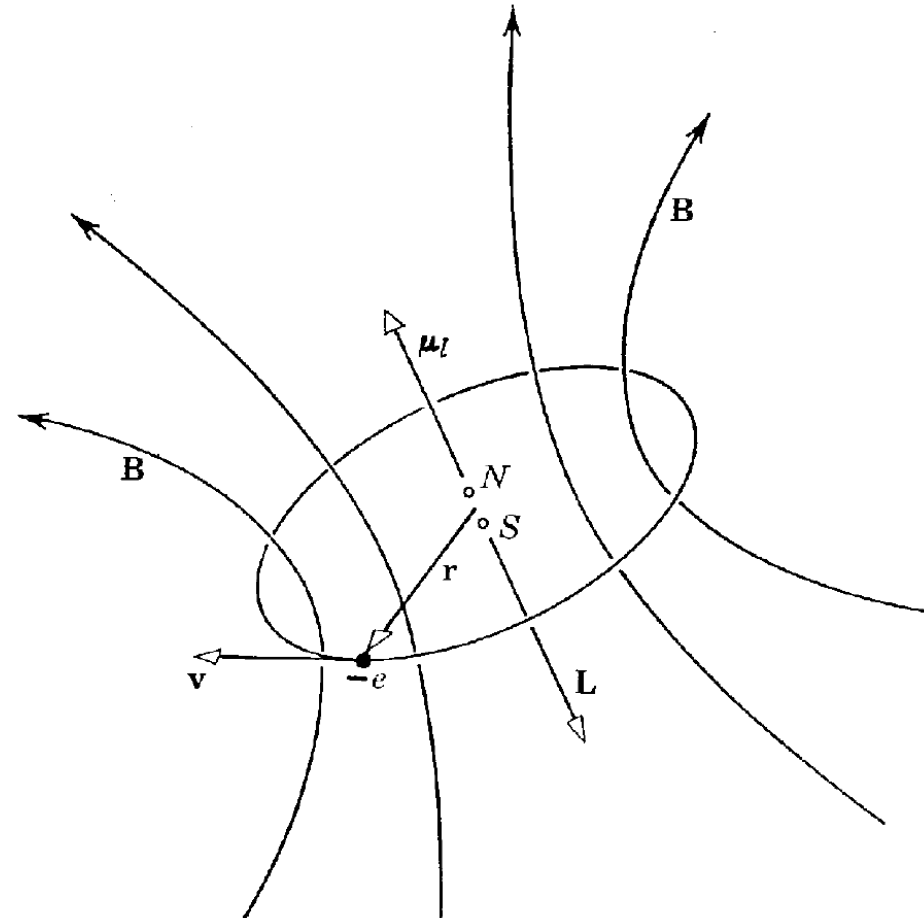
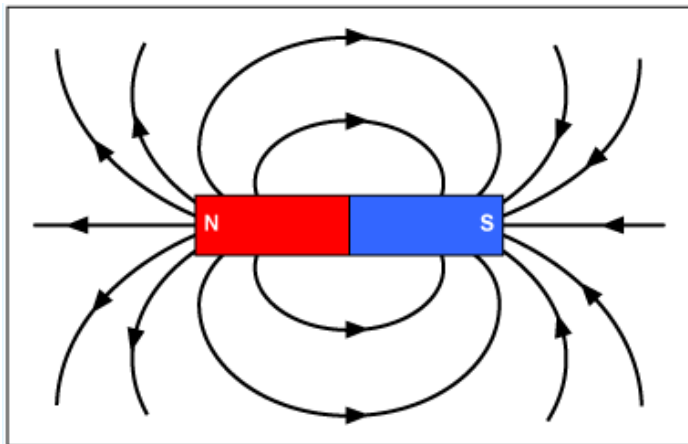
$$\mu_z = \frac{g_l \mu_b}{\hbar} L_z = \frac{g_l \mu_b}{\hbar} m_l \hbar = g_l \mu_b m_l$$

$$\left\{ \begin{array}{l} l = 0, 1, 2, \dots \quad \text{inteiro} \\ -l \leq m_l \leq l \quad \text{inteiro} \end{array} \right.$$

Espira é um dipolo magnético



Espira gera campo magnético semelhante ao de um imã :

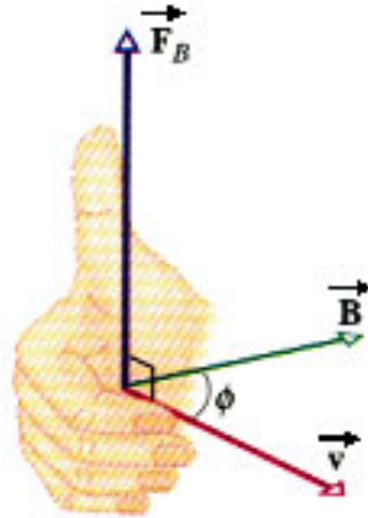
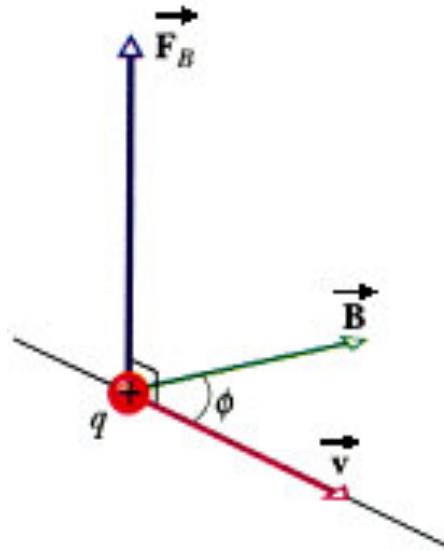


Dipolo magnético !

Espira num campo magnético uniforme

Força magnética sobre uma carga em movimento

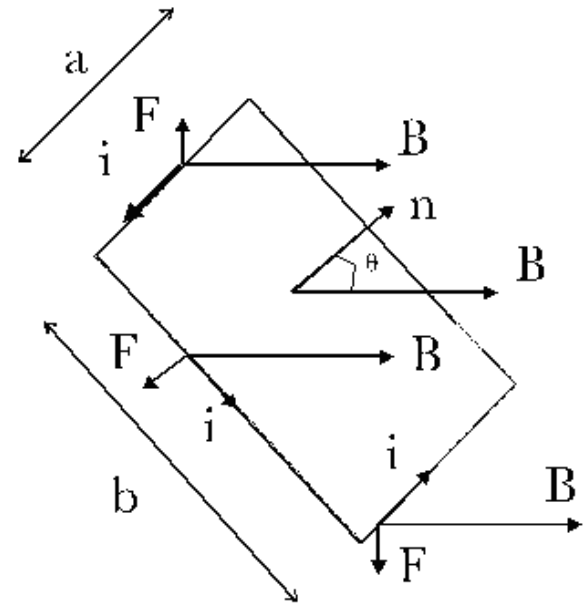
$$\vec{F} = q \vec{v} \times \vec{B}$$



Força magnética sobre uma corrente

$$\vec{F} = I \vec{L} \times \vec{B}$$

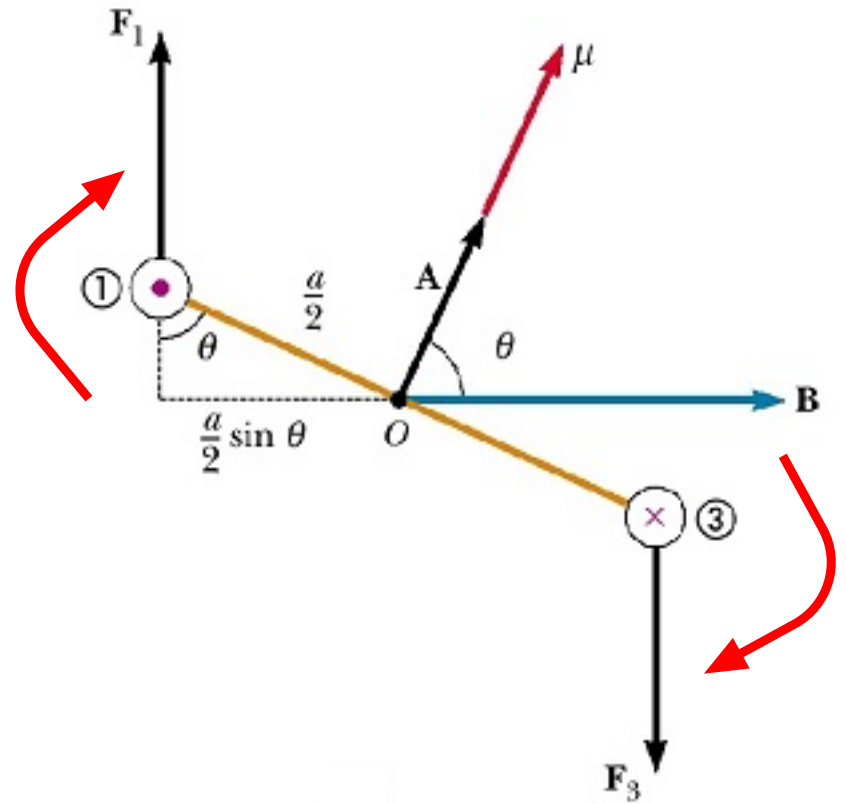
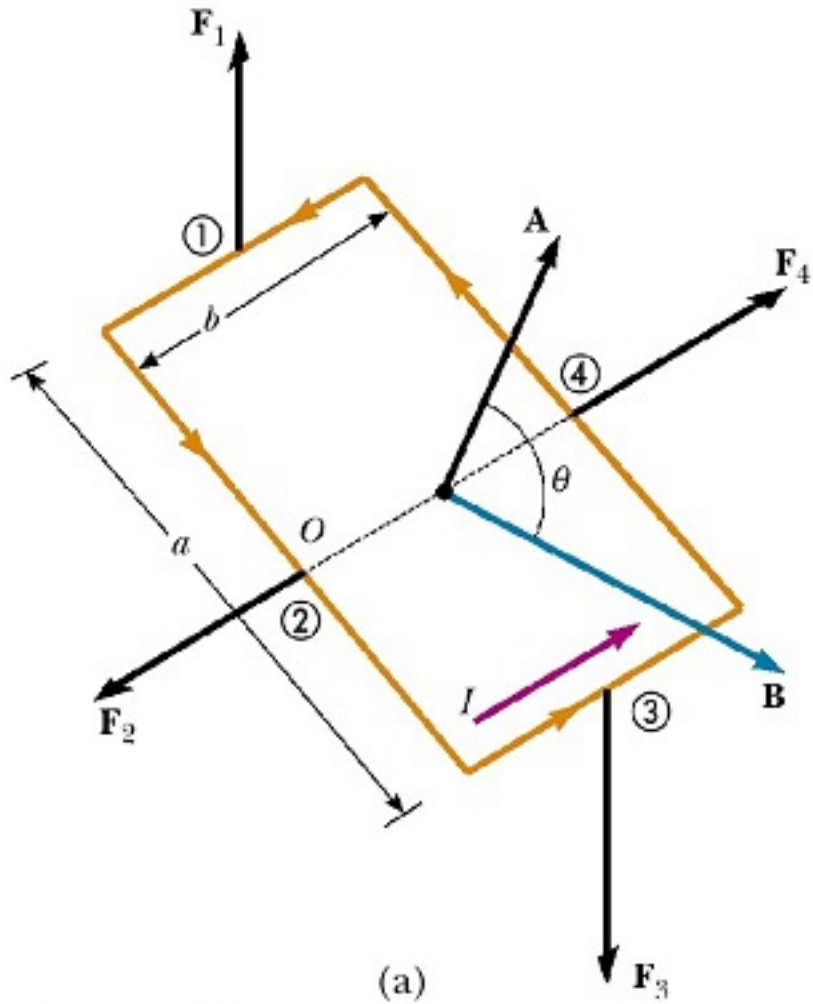
(\vec{L} no sentido da corrente)



Espira sofre um torque !

$$\vec{\mu} = I \vec{A}$$

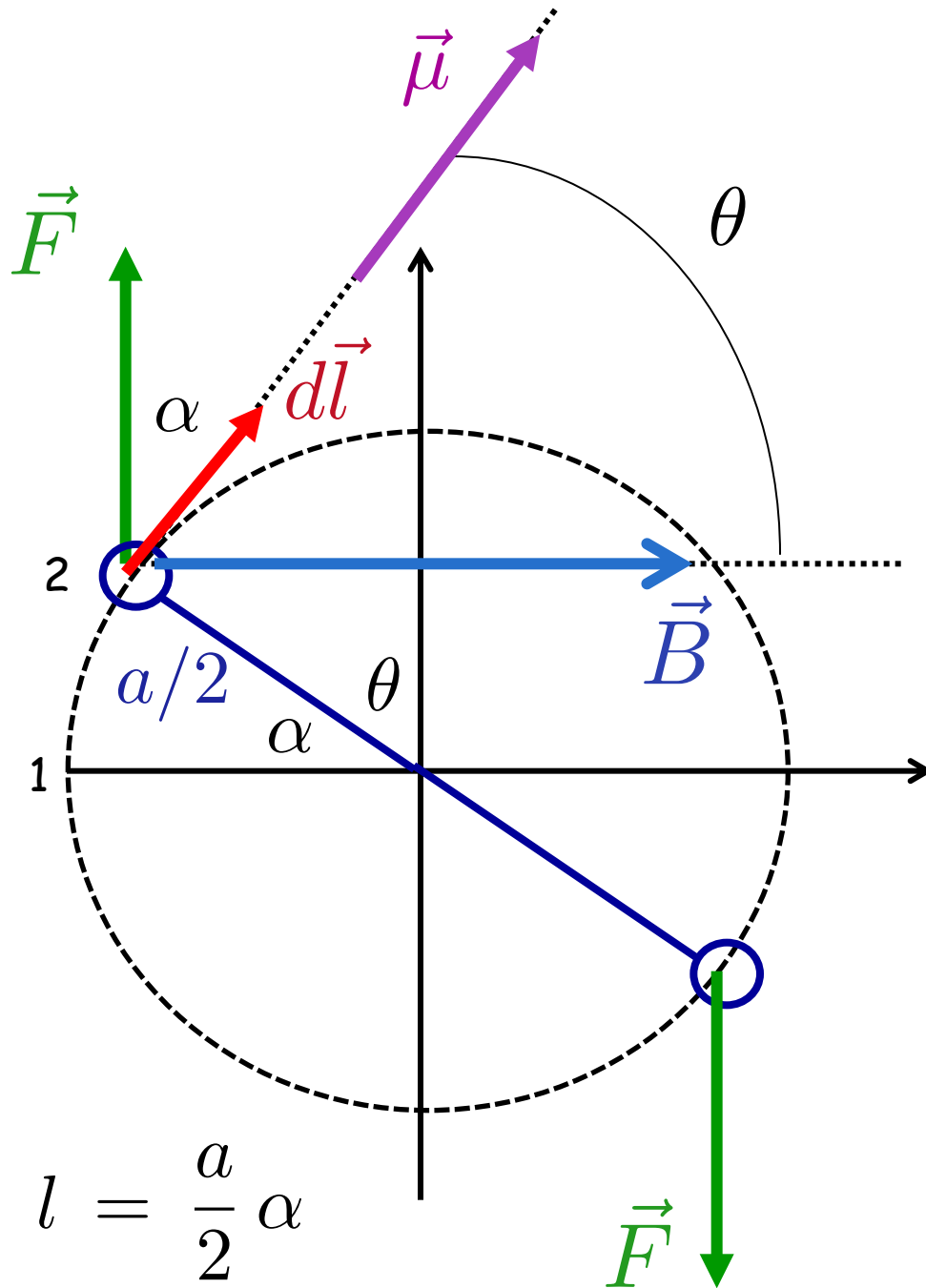
$$\mu = I a b$$



$$\vec{F} = I \vec{L} \times \vec{B}$$

$$F = I b B$$

Trabalho da força F quando a espira vai de 1 a 2.:



$$dW = \vec{F} \cdot d\vec{l}$$

$$= F dl \cos \alpha$$

$$= F \frac{a}{2} d\alpha \cos \alpha$$

$$W = \int_0^\alpha F \frac{a}{2} \cos \alpha d\alpha$$

$$W = F \frac{a}{2} \sin \alpha$$

$$W = F \frac{a}{2} \cos \theta$$

$$W = F \frac{a}{2} \cos \theta$$

Somando a contribuição da barra inferior (3) :

$$W = F a \cos \theta \qquad F = I b B$$

$$W = I a b B \cos \theta \qquad \mu = I a b$$

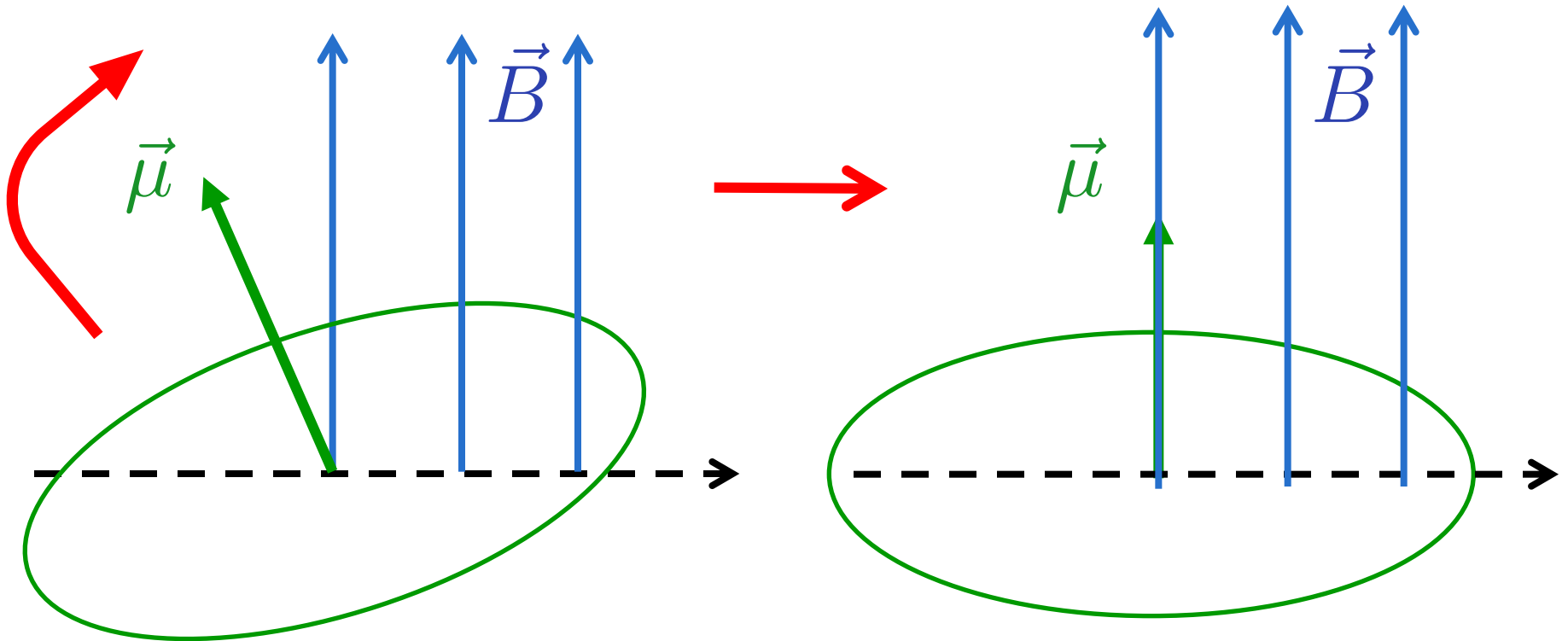
$$W = \mu B \cos \theta$$

$$W = \vec{\mu} \cdot \vec{B} \qquad W = -U$$

$$U = -\vec{\mu} \cdot \vec{B}$$

$$U = -\mu B \cos \theta$$

Espira tenta se alinhar como uma bússola !



mais energia !

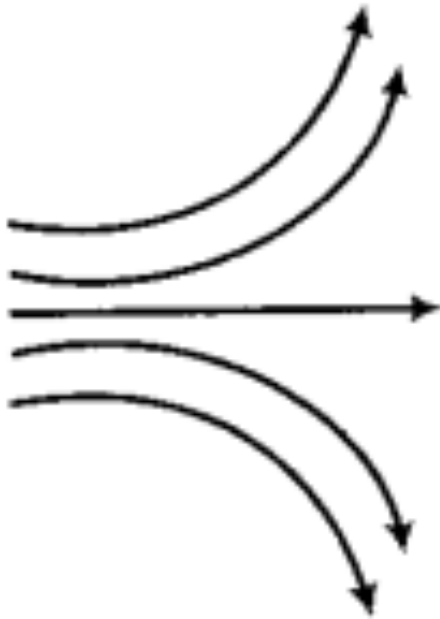
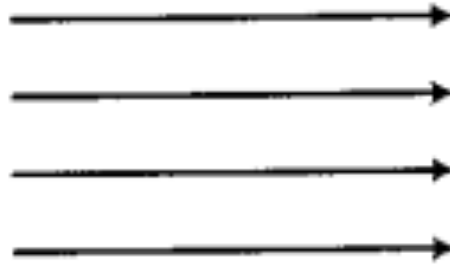
menos energia !

Energia potencial de orientação :

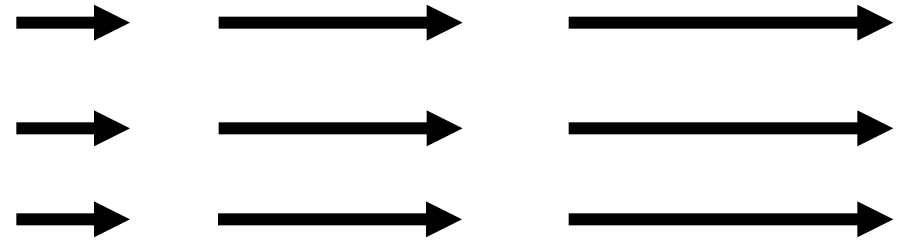
$$U = -\vec{\mu} \cdot \vec{B}$$

$$U = -\mu B \cos \theta$$

Espira num campo magnético não uniforme

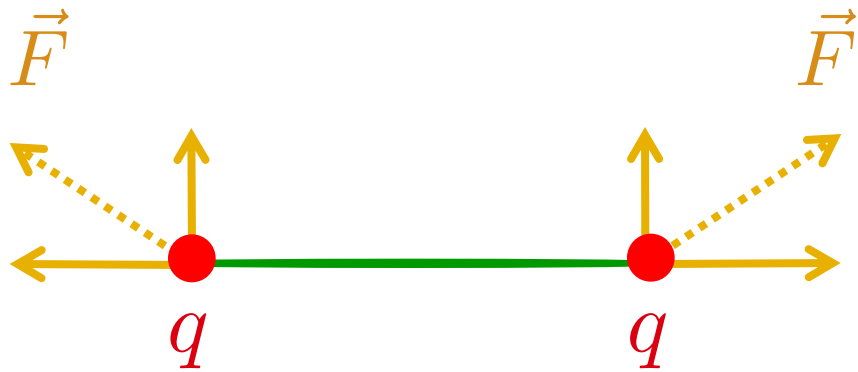
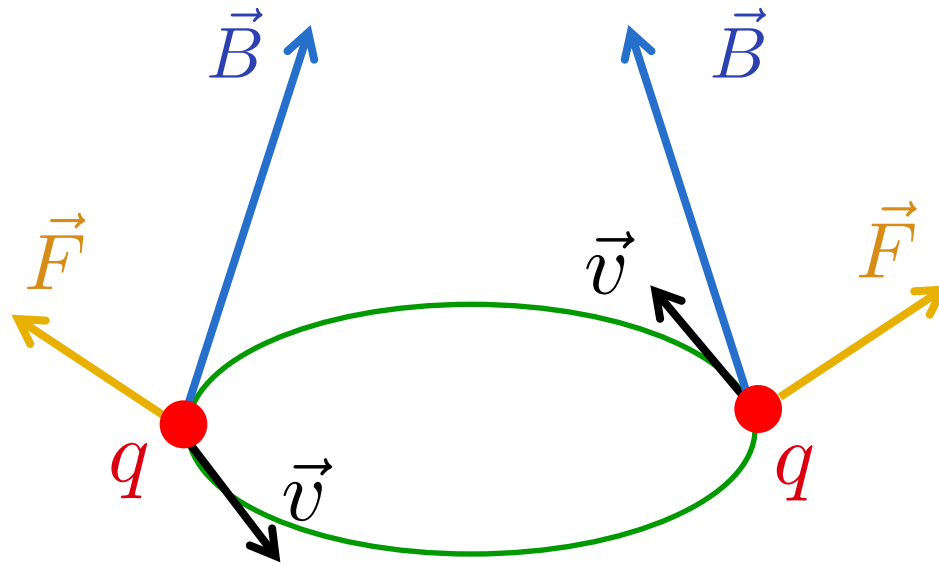


não uniforme

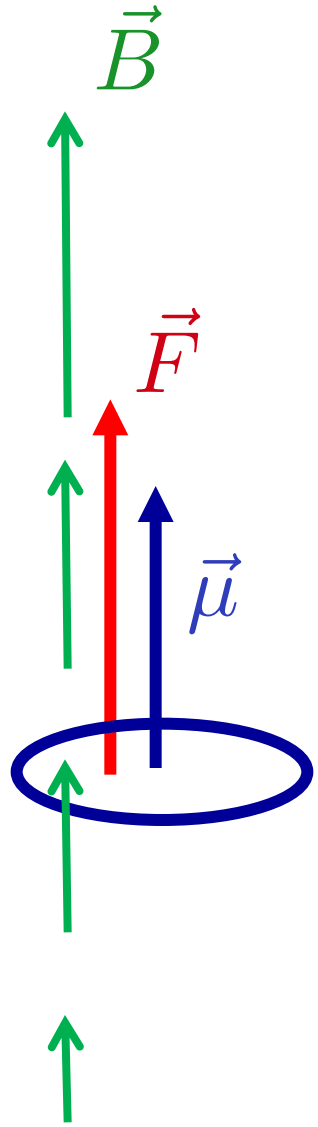


não uniforme

Espira num campo magnético não uniforme



Campo magnético não uniforme



$$\vec{B} = B_z \hat{k} \quad \vec{\mu} = \mu_z \hat{k}$$

$$U = -\vec{\mu} \cdot \vec{B} = -\mu_z B_z \quad \text{energia potencial mínima}$$

$$\vec{F} = -\vec{\nabla} U = - \left[\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right] U$$

$$\vec{F} = \mu_z \frac{\partial B_z}{\partial z} \hat{k} \quad \mu_z = \frac{g_l \mu_b}{\hbar} L_z$$

$$F = \frac{g_l \mu_b L_z}{\hbar} \frac{\partial B_z}{\partial z}$$

Mecânica clássica :

L_z assume
valores contínuos !

$$F = \frac{g_l \mu_b L_z}{\hbar} \frac{\partial B_z}{\partial z}$$

F assume
valores contínuos

Mecânica quântica :

$$L_z = m_l \hbar$$

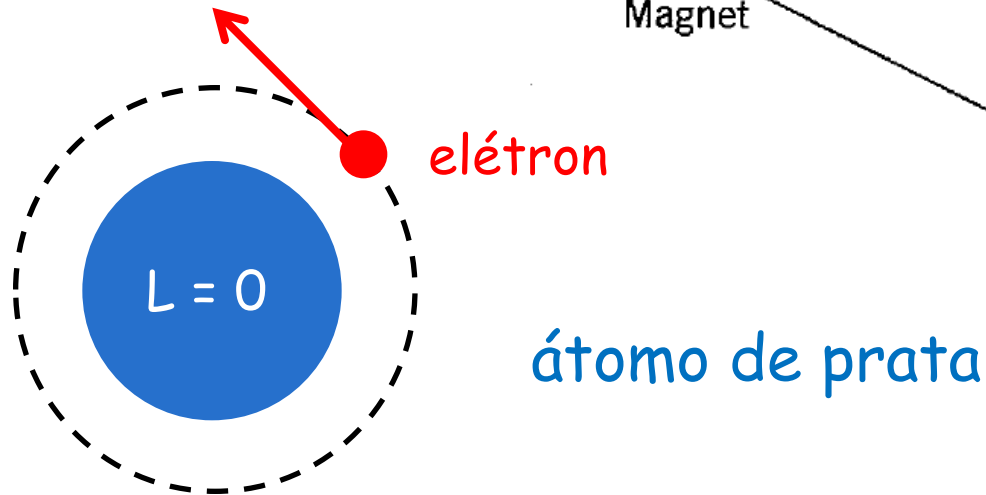
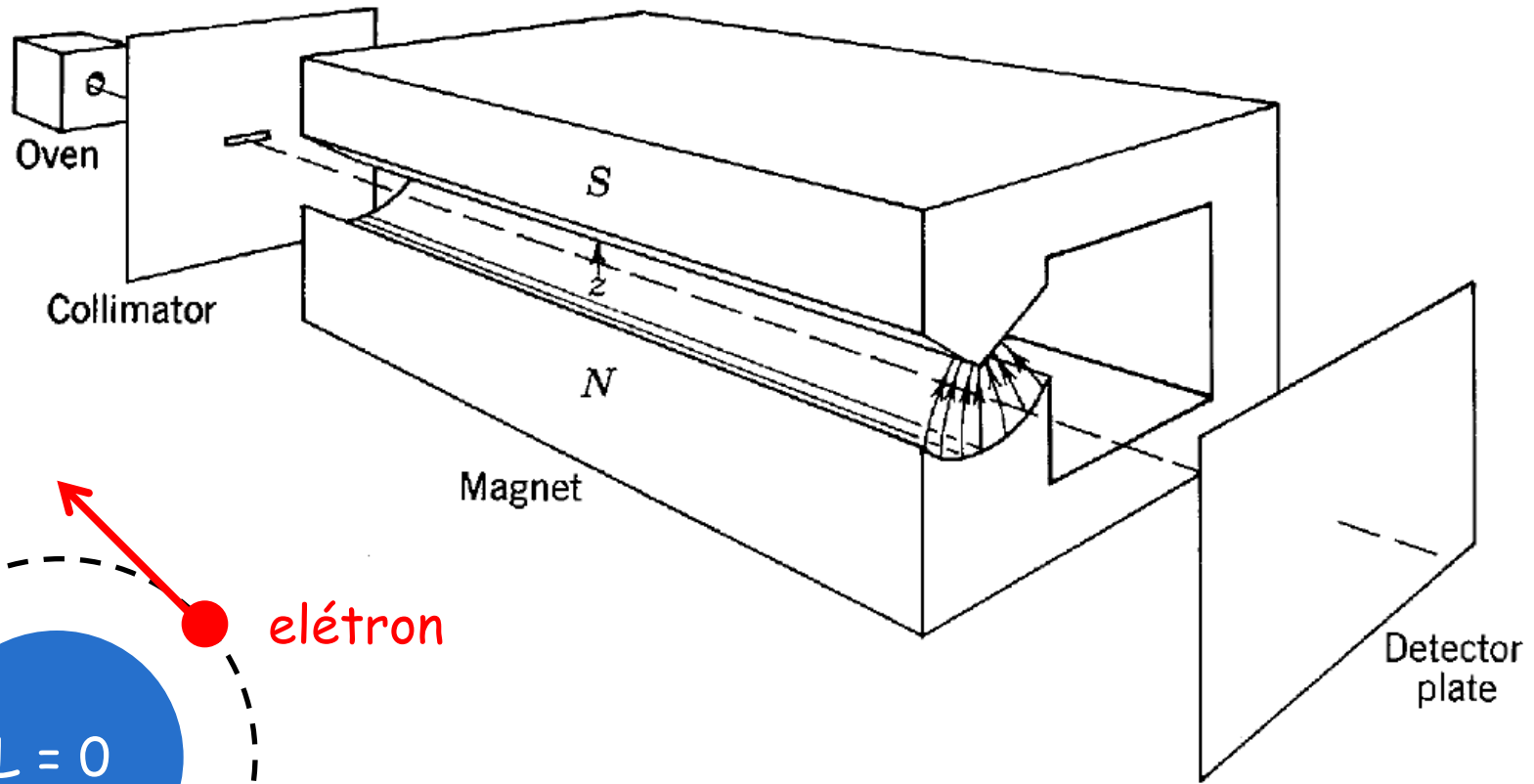
$$\left\{ \begin{array}{ll} l = 0, 1, 2, \dots & \text{inteiro} \\ -l \leq m_l \leq l & \text{inteiro} \end{array} \right.$$

$$F = g_l \mu_b m_l \frac{\partial B_z}{\partial z}$$

F assume
valores discretos

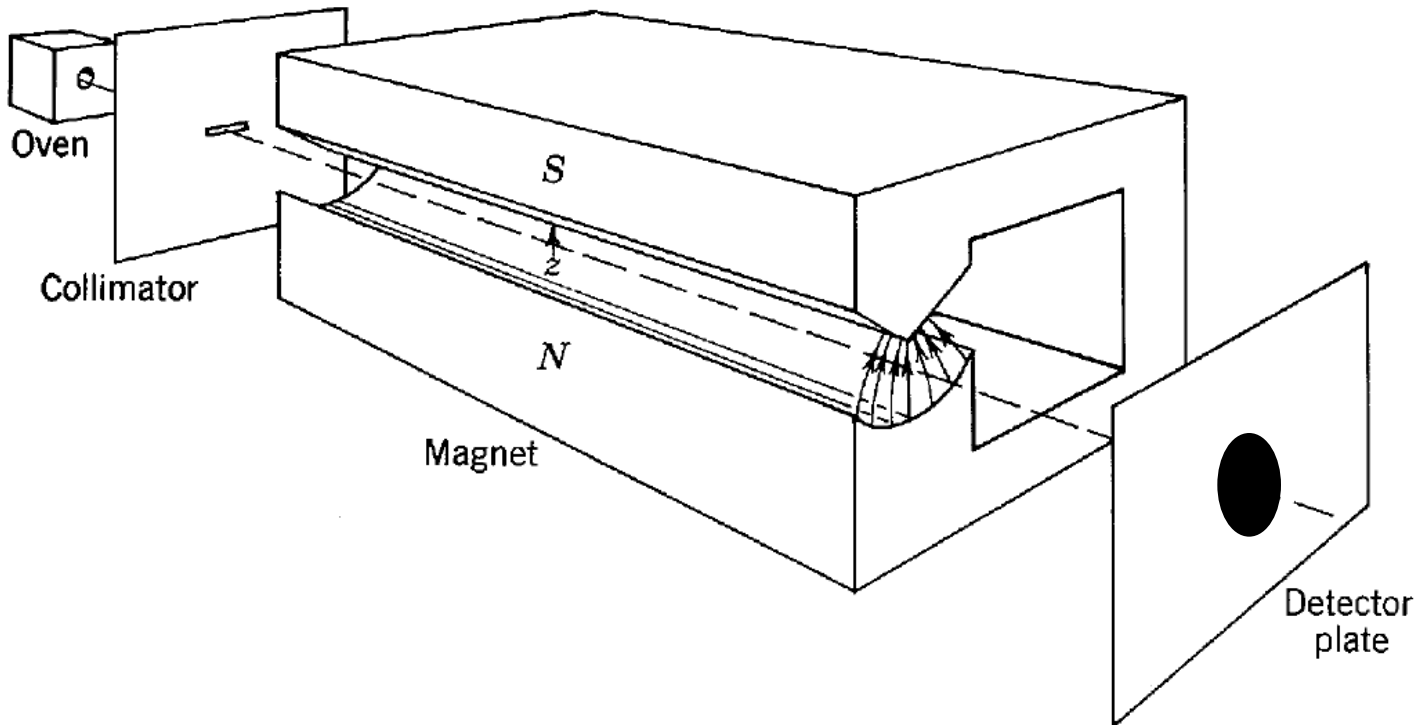
Experiência de Stern-Gerlach (1922)

Medida do momento magnético de átomos de prata



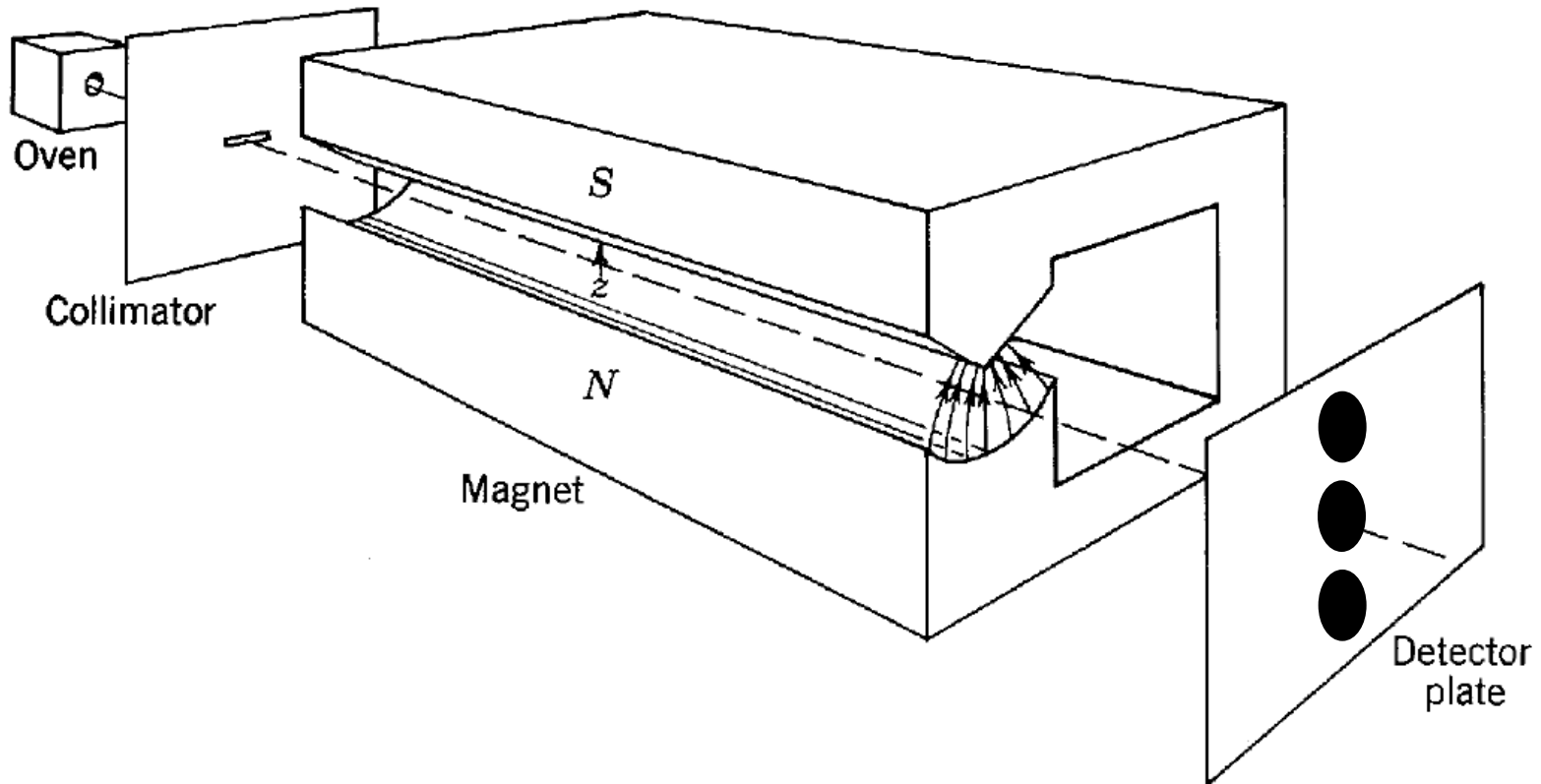
Se o elétron tiver momento angular zero:

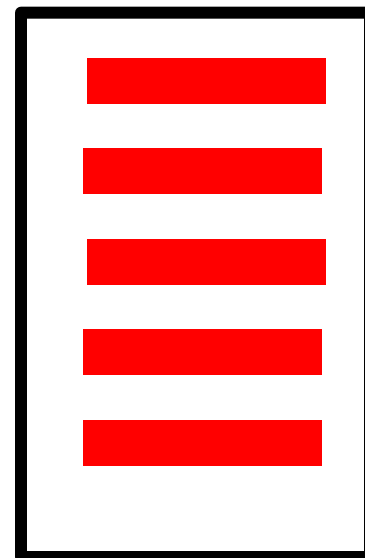
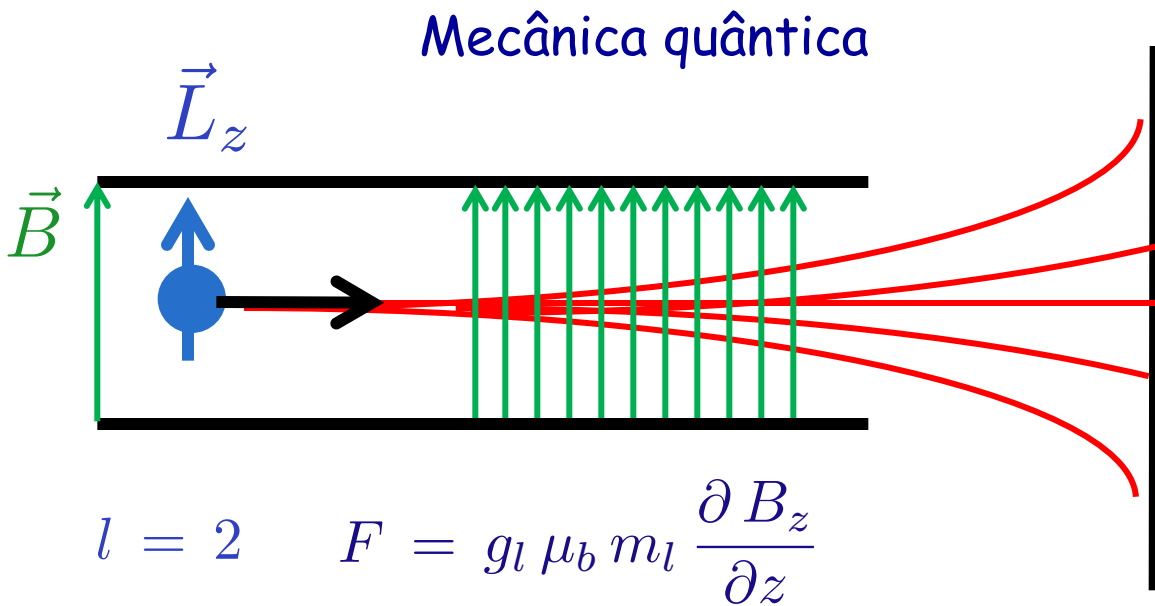
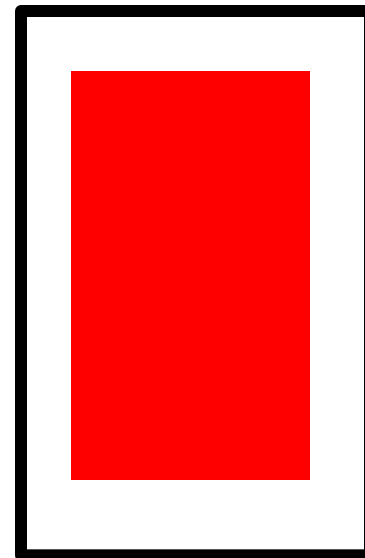
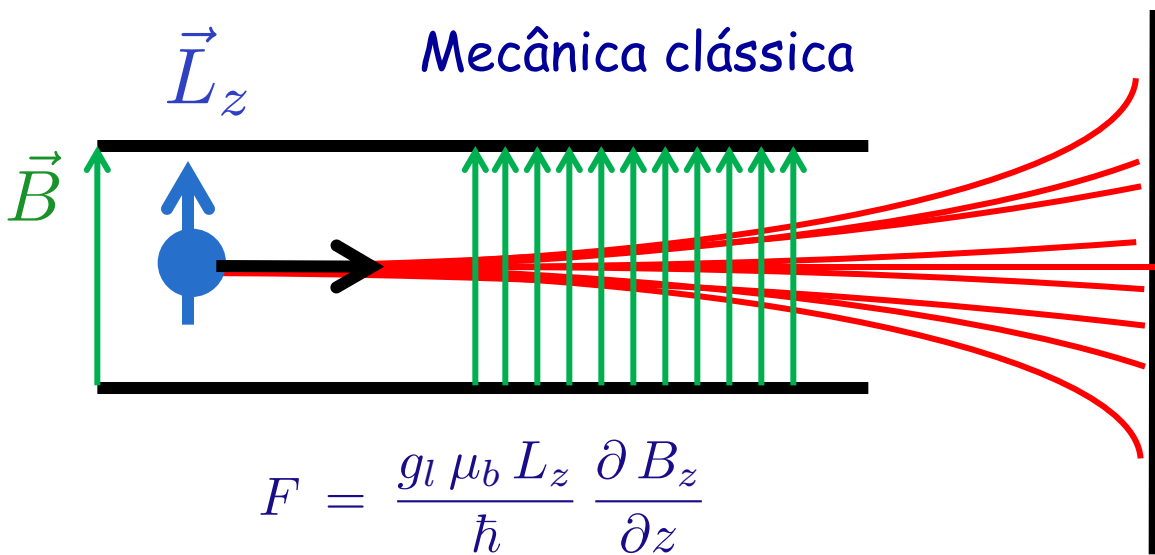
$$l = 0 \quad m_l = 0 \quad F = g_l \mu_b m_l \frac{\partial B_z}{\partial z} \quad F = 0$$



Se o elétron tiver momento angular "um" :

$$l = 1 \left\{ \begin{array}{l} m_l = -1 \\ m_l = 0 \\ m_l = 1 \end{array} \right. \quad F = g_l \mu_b m_l \frac{\partial B_z}{\partial z} \quad \left\{ \begin{array}{l} F = -g_l \mu_b \frac{\partial B_z}{\partial z} \\ F = 0 \\ F = +g_l \mu_b \frac{\partial B_z}{\partial z} \end{array} \right.$$





$$m_l = -2$$

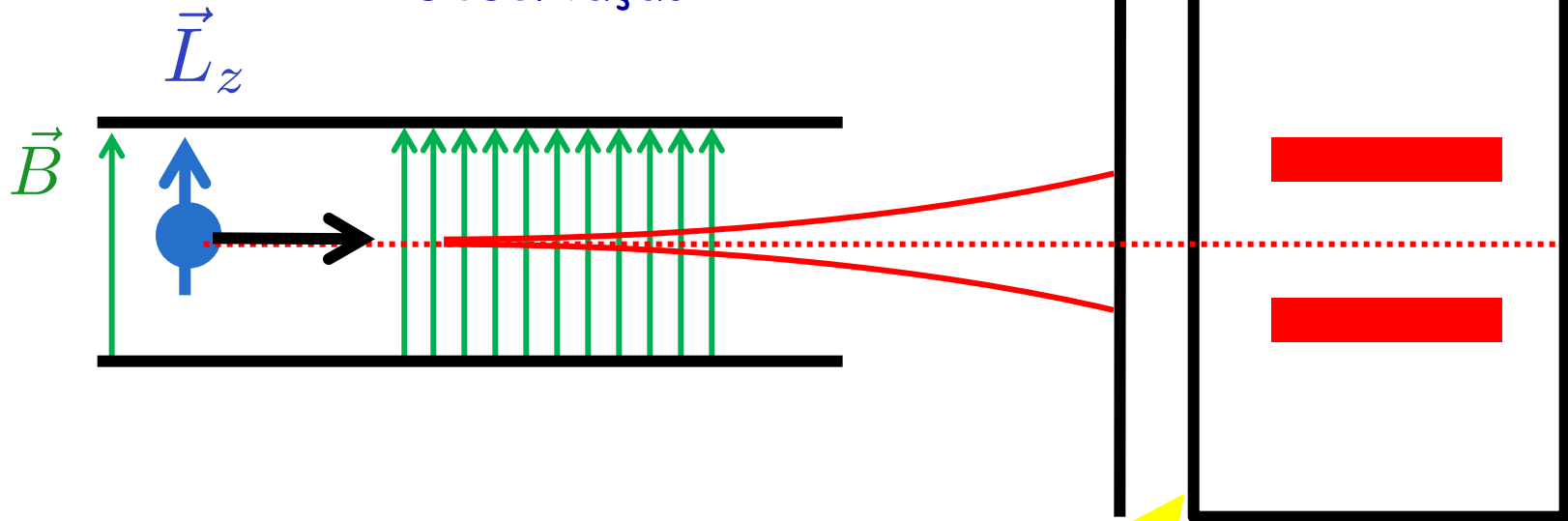
$$m_l = -1$$

$$m_l = 0$$

$$m_l = 1$$

$$m_l = 2$$

Observação :



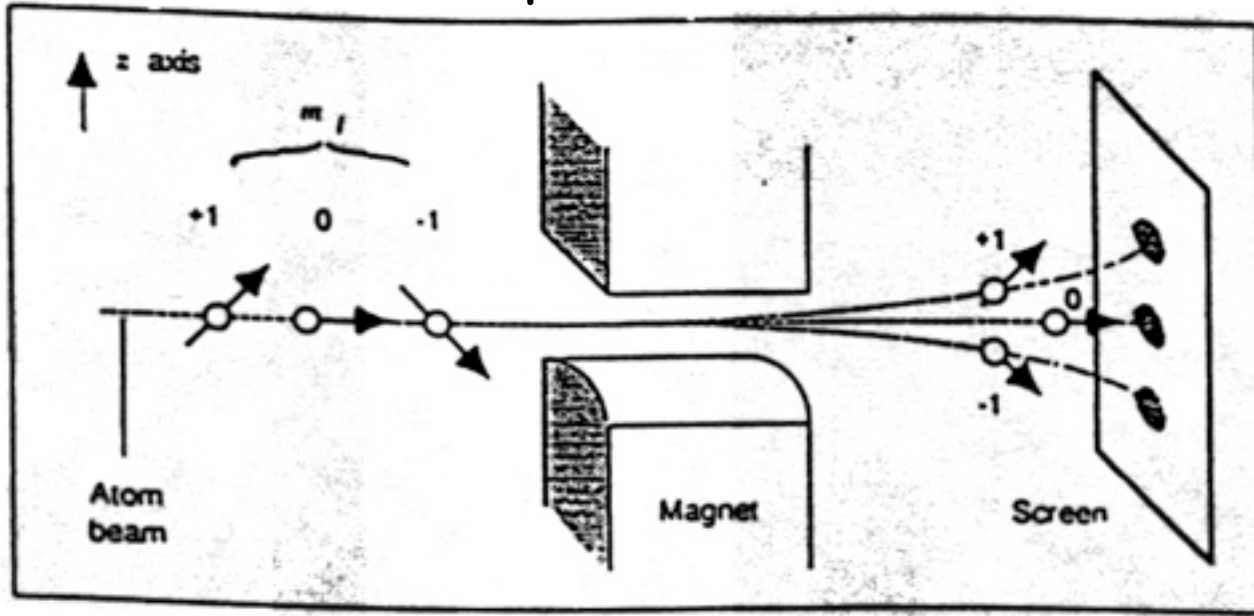
Só duas franjas sempre ! Dois valores de m : $+m$ e $-m$

Não existe $m = 0$!

clássico

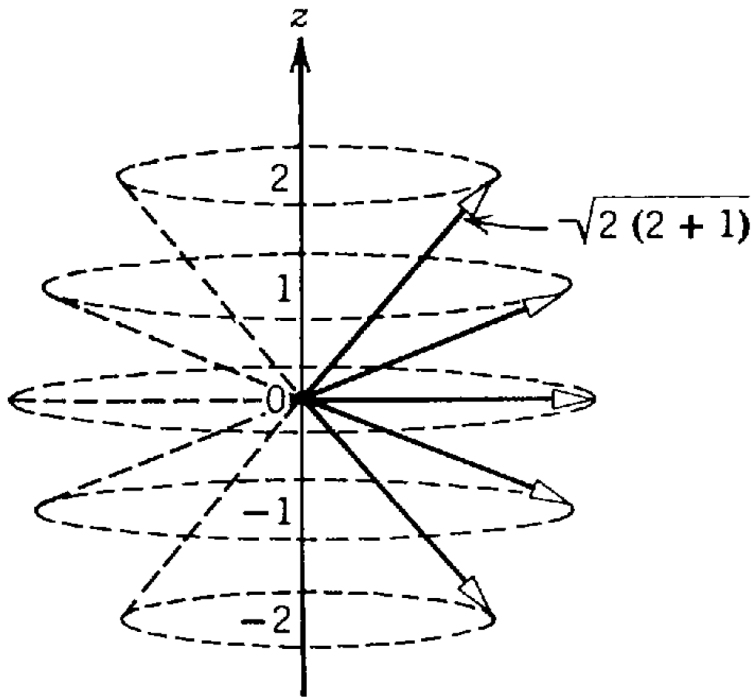
quântico

observado



Elétron tem momento angular intrínseco : spin !

Copiamos a teoria do momento angular

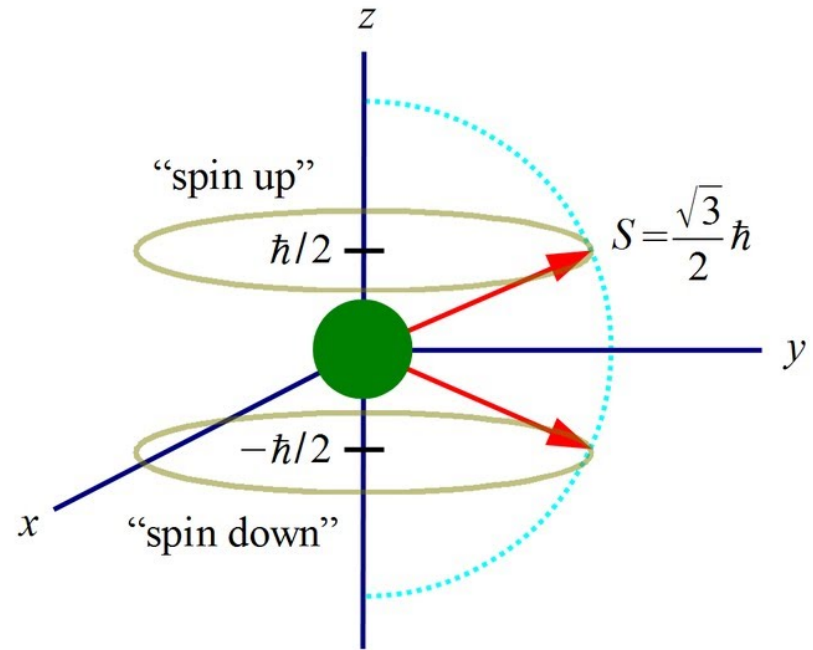


$$L = \sqrt{l(l+1)} \hbar$$

$$L_z = m_l \hbar$$

$$-l \leq m_l \leq l$$

$$l = 0, 1, 2, \dots$$



$$S = \sqrt{s(s+1)} \hbar$$

$$S_z = m_s \hbar$$

$$m_s = -s, -s+1, \dots, s$$

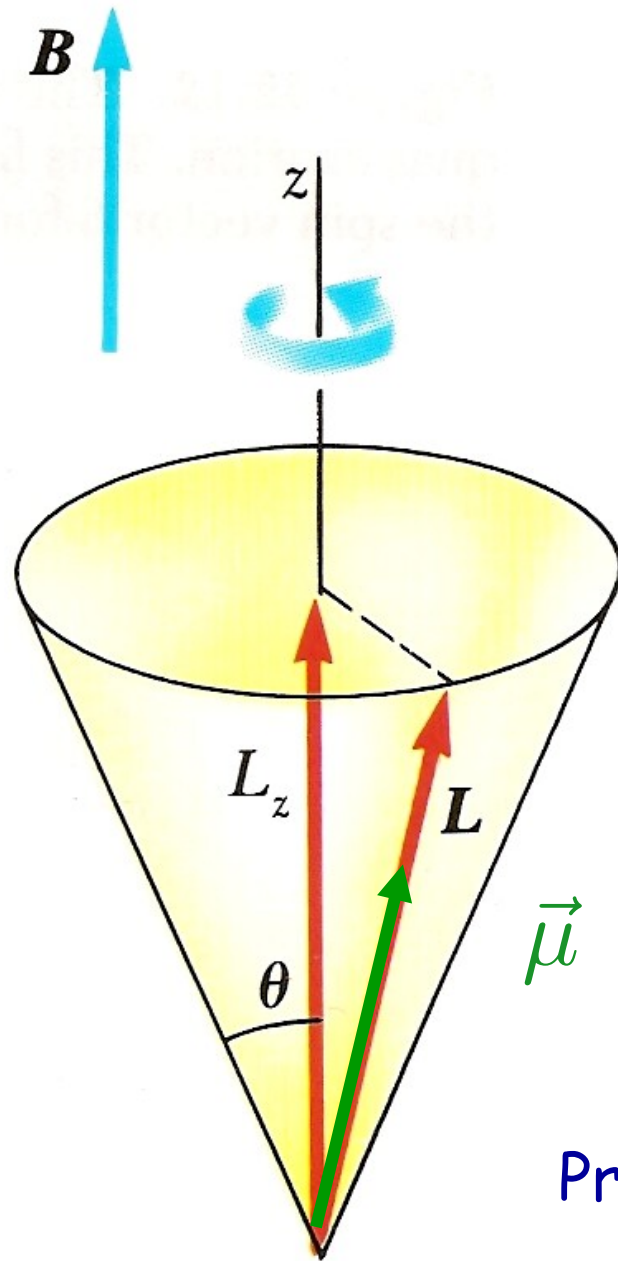
$$s = \frac{1}{2} \quad m_s = -\frac{1}{2}, +\frac{1}{2}$$

FIM

A Little History of Spin in Quantum Mechanics

- 1922 – Otto Stern & Walter Gerlach – The existence of spin angular momentum is inferred from their experiment, in which particles (Ag atoms) are observed to possess angular momentum that cannot be accounted for by orbital angular momentum alone.
- 1924 – Wolfgang Pauli – proposed a new quantum degree of freedom (or quantum number) with two possible values and formulated the Pauli exclusion principle.
- 1925 – Ralph Kronig, George Uhlenbeck & Samuel Goudsmit – identified Pauli's new degree of freedom as electron spin and suggested a physical interpretation of particles spinning around their own axis.
- 1926 – Enrico Fermi & Paul Dirac – formulated (independently) the Fermi-Dirac statistics, which describes distribution of many identical particles obeying the Pauli exclusion principle (fermions with half-integer spins – contrary to bosons satisfying the Bose-Einstein statistics)
- 1926 – Erwin Schrödinger – formulated his non-relativistic Schrödinger equation, but it incorrectly predicted the magnetic moment of H to be zero in its ground state.
- 1927 – T.E. Phipps & J.B. Taylor – reproduced the effect using H atoms in the ground state, thereby eliminating any doubts that may have been caused by the use of Ag atoms.
- 1927 – Wolfgang Pauli – worked out on mathematical formulation of spin (2×2 matrices).
- 1928 – Paul Dirac – showed that spin comes naturally from his relativistic Dirac equation.

Mas nunca se alinha! Precessiona!

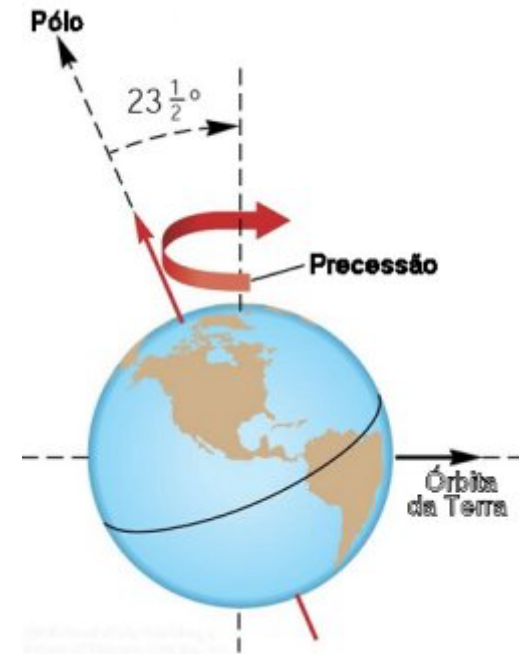
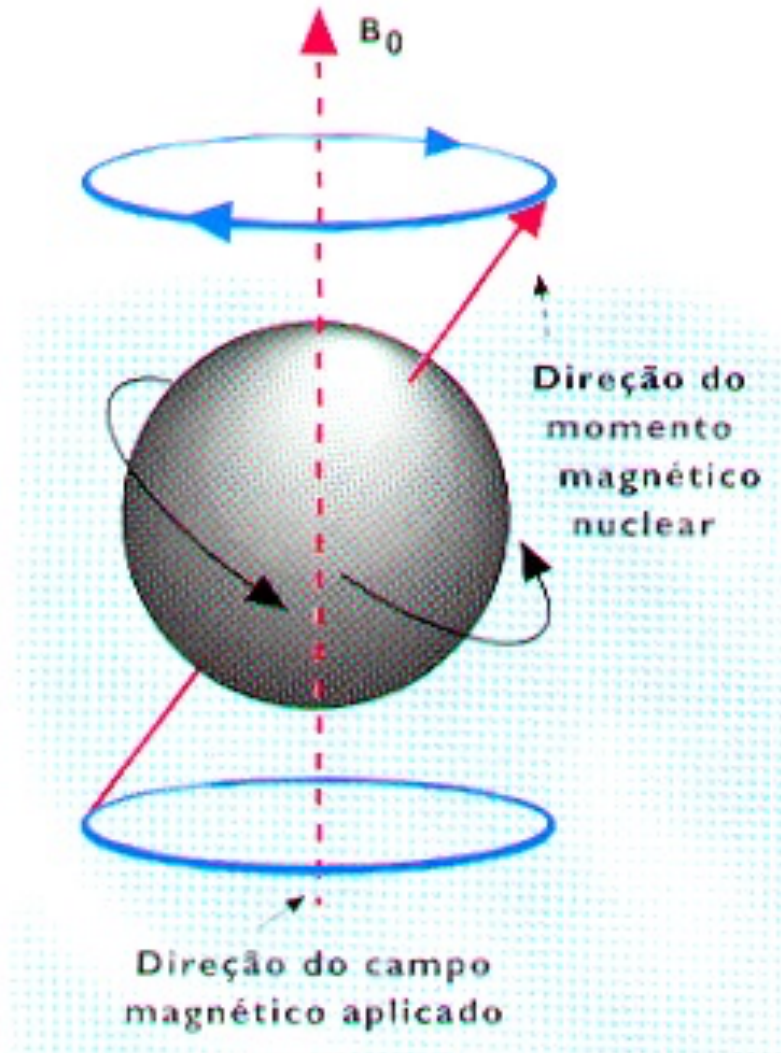


$$\vec{\mu} = \frac{g_l \mu_b}{\hbar} \vec{L}$$

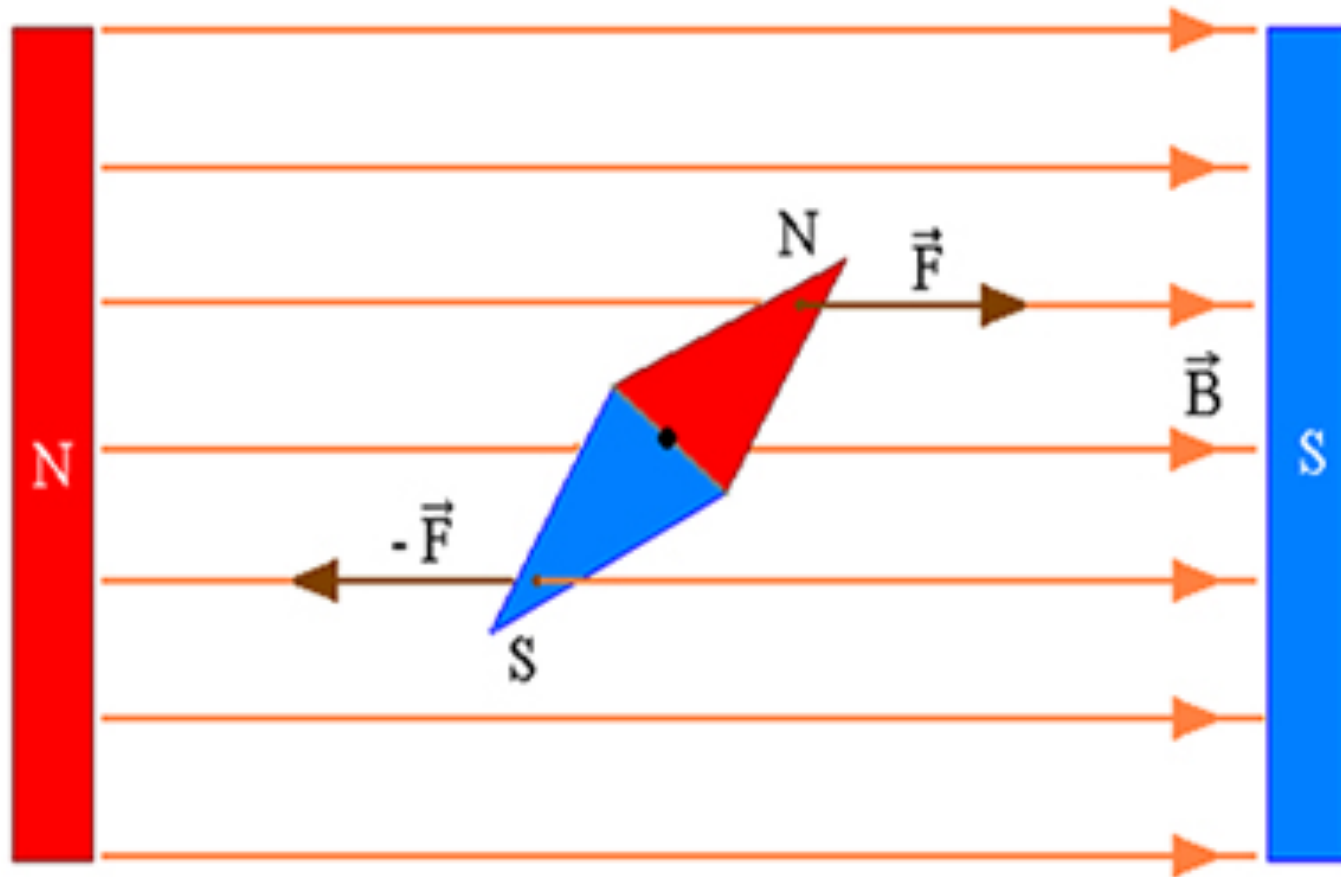
Precessão de Larmor

Mas nunca se alinha ! Precessiona !

Precessão de Larmor

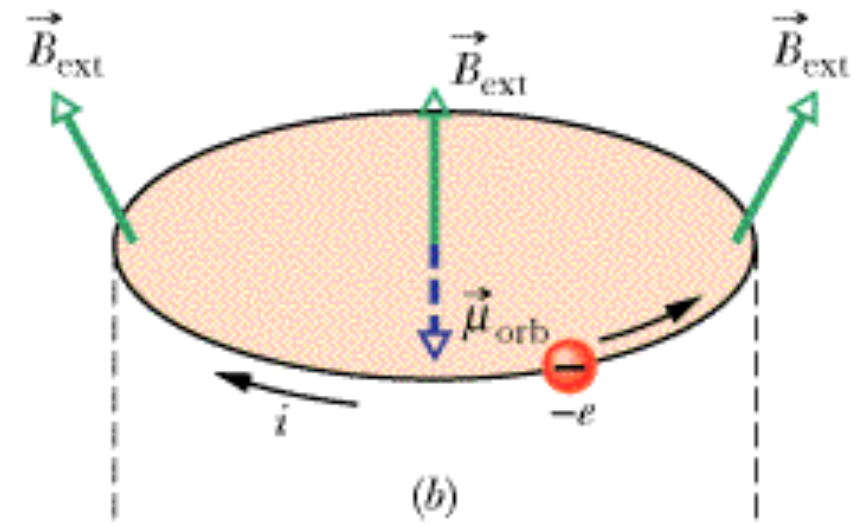
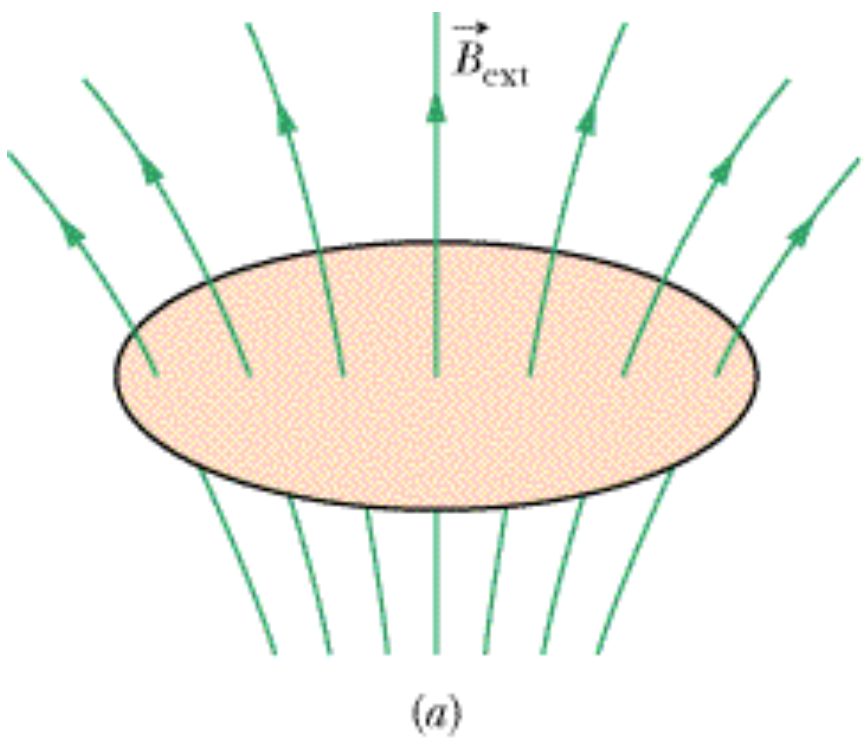


Espira se alinha como uma bússola !



Um imã é uma espira ?

Espira num campo magnético não uniforme



Levitação magnética!

maglev



**How to
levitate
a frog?**



sapo



sapo no youtube