

- Qual a função de onda esférica?
- Existe uma função de onda esférica para ondas transversais?
- O que é uma fonte puntiforme para o eletromagnetismo?
- Onde nasce uma onda?

# Ondas esféricas: Acústica x Ótica

Som: onda de pressão

$$\text{Solução esférica: } p(\vec{r}, t) = p_0 \frac{1}{r} \cos(kr - \omega t + \varphi) =$$

$$= \text{Re} \left[ \frac{\tilde{p}_0}{r} e^{i(kr - \omega t)} \right] = \text{Re}[\tilde{p}]$$

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\frac{\partial \tilde{p}}{\partial x} = \frac{\partial \tilde{p}}{\partial r} \cdot \frac{\partial r}{\partial x}$$

$$\frac{\partial r}{\partial x} = \frac{\partial \sqrt{x^2 + y^2 + z^2}}{\partial x} = \frac{1}{2} \frac{2x}{\sqrt{x^2 + y^2 + z^2}} = \frac{x}{r}$$

$$= \left[ -\frac{\tilde{p}_0}{r^2} e^{i(kr - \omega t)} + ik \tilde{p} \right] \frac{x}{r}$$

$$= \left( ik - \frac{1}{r} \right) \frac{x}{r} \tilde{p}$$

$$\frac{\partial^2}{\partial x^2} \tilde{p} = \frac{\partial}{\partial x} \tilde{p} \cdot \left( ik - \frac{1}{r} \right) \cdot \frac{x}{r} + \tilde{p} \frac{\partial}{\partial x} \left( ik \frac{x}{r} - \frac{x}{r^2} \right)$$

$$\frac{\partial}{\partial x} \left( ik \frac{x}{r} \right) = ik \frac{1}{r} + ikx \frac{\partial}{\partial x} \left( \frac{1}{r} \right)$$

$$= ik \frac{1}{r} - ik \frac{x^2}{r^3}$$

$$\frac{\partial}{\partial x} \left( \frac{1}{r} \right) = \frac{\partial}{\partial r} \left( \frac{1}{r} \right) \frac{\partial r}{\partial x} =$$

$$= -\frac{1}{r^2} \cdot \frac{x}{r} = -\frac{x}{r^3}$$

$$\frac{\partial}{\partial x} \left( \frac{x}{r^2} \right) = \frac{1}{r^2} + x \frac{\partial}{\partial x} \left( \frac{1}{r^2} \right)$$

$$= \frac{1}{r^2} - 2 \frac{x^2}{r^4}$$

$$\frac{\partial}{\partial x} \left( \frac{1}{r^2} \right) = \frac{\partial}{\partial r} \left( \frac{1}{r^2} \right) \cdot \frac{\partial r}{\partial x}$$

$$= -\frac{2}{r^3} \cdot \frac{x}{r}$$

$$\frac{\partial^2}{\partial x^2} \tilde{p} = \tilde{p} \left[ \left( ik - \frac{1}{r} \right)^2 \frac{x^2}{r^2} + \frac{ik}{r} - ik \frac{x^2}{r^3} - \frac{1}{r^2} + 2 \frac{x^2}{r^4} \right]$$

$$= \tilde{p} \left[ -k^2 \frac{x^2}{r^2} - 2ik \frac{x^2}{r^3} + \frac{x^2}{r^4} + \frac{ik}{r} - ik \frac{x^2}{r^3} - \frac{1}{r^2} + 2 \frac{x^2}{r^4} \right]$$

$$= \tilde{p} \left[ -k^2 \frac{x^2}{r^2} - 3ik \frac{x^2}{r^3} + 3 \frac{x^2}{r^4} + \frac{ik}{r} - \frac{1}{r^2} \right]$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \tilde{p} = \tilde{p} \left[ -k^2 - 3ik \frac{1}{r} + \frac{3}{r^2} + 3ik \frac{1}{r} - \frac{3}{r^2} \right] = -k^2 \tilde{p}$$



Coordenadas esféricas:  $\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2}$

$$\nabla^2 \tilde{\psi}(r) = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \tilde{\psi}(r) \right) = \left[ \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \left( \frac{e^{ikr}}{r} \right) \right) \right] \tilde{p}_0 e^{-i\omega t}$$

$$\frac{\partial}{\partial r} \left( \frac{e^{ikr}}{r} \right) = -\frac{1}{r^2} e^{ikr} + ik \frac{e^{ikr}}{r}$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left[ r^2 \frac{\partial}{\partial r} \left( \frac{e^{ikr}}{r} \right) \right] = \frac{1}{r^2} \frac{\partial}{\partial r} \left[ -e^{ikr} + ik r e^{ikr} \right]$$

$$= \frac{1}{r^2} \left[ -ik e^{ikr} + ik e^{ikr} + (ik)^2 r e^{ikr} \right] = -k^2 \frac{e^{ikr}}{r}$$

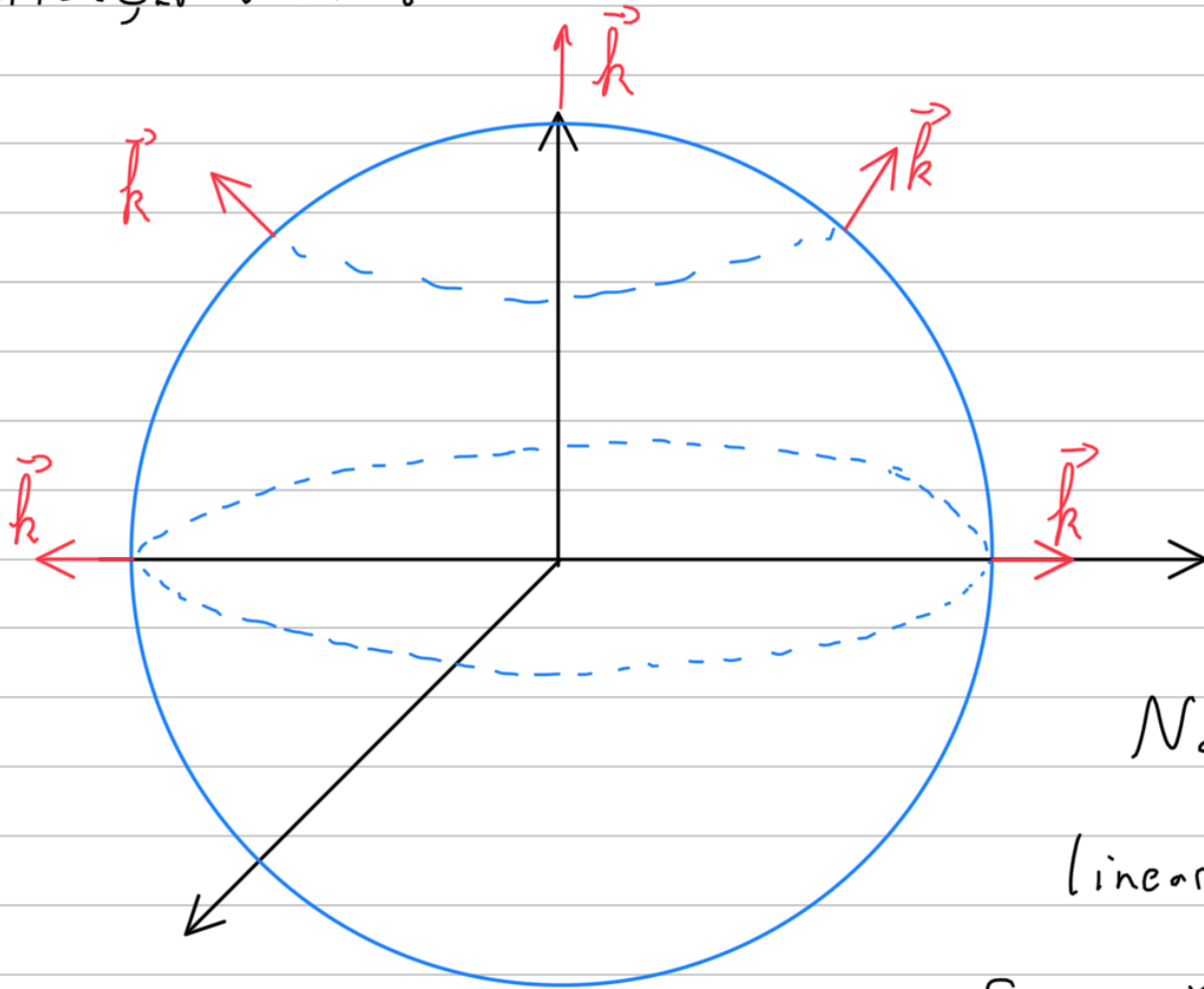
$$\therefore \nabla^2 \tilde{\psi}(r) = -k^2 \tilde{\psi}(r)$$

$$\frac{\partial^2}{\partial t^2} \tilde{\psi}(r) = -\omega^2 \tilde{\psi}(r)$$

$$\left. \begin{array}{l} \nabla^2 \tilde{\psi}(r) = -k^2 \tilde{\psi}(r) \\ \frac{\partial^2}{\partial t^2} \tilde{\psi}(r) = -\omega^2 \tilde{\psi}(r) \end{array} \right\} \left( \nabla^2 - \frac{1}{v^2} \frac{\partial^2}{\partial t^2} \right) \tilde{\psi} = 0$$

Ondas transversais:  $\vec{E} = \frac{E}{r} \hat{n} e^{i(\vec{k} \cdot \vec{r} - \omega t)}$  ?

polarização varia!

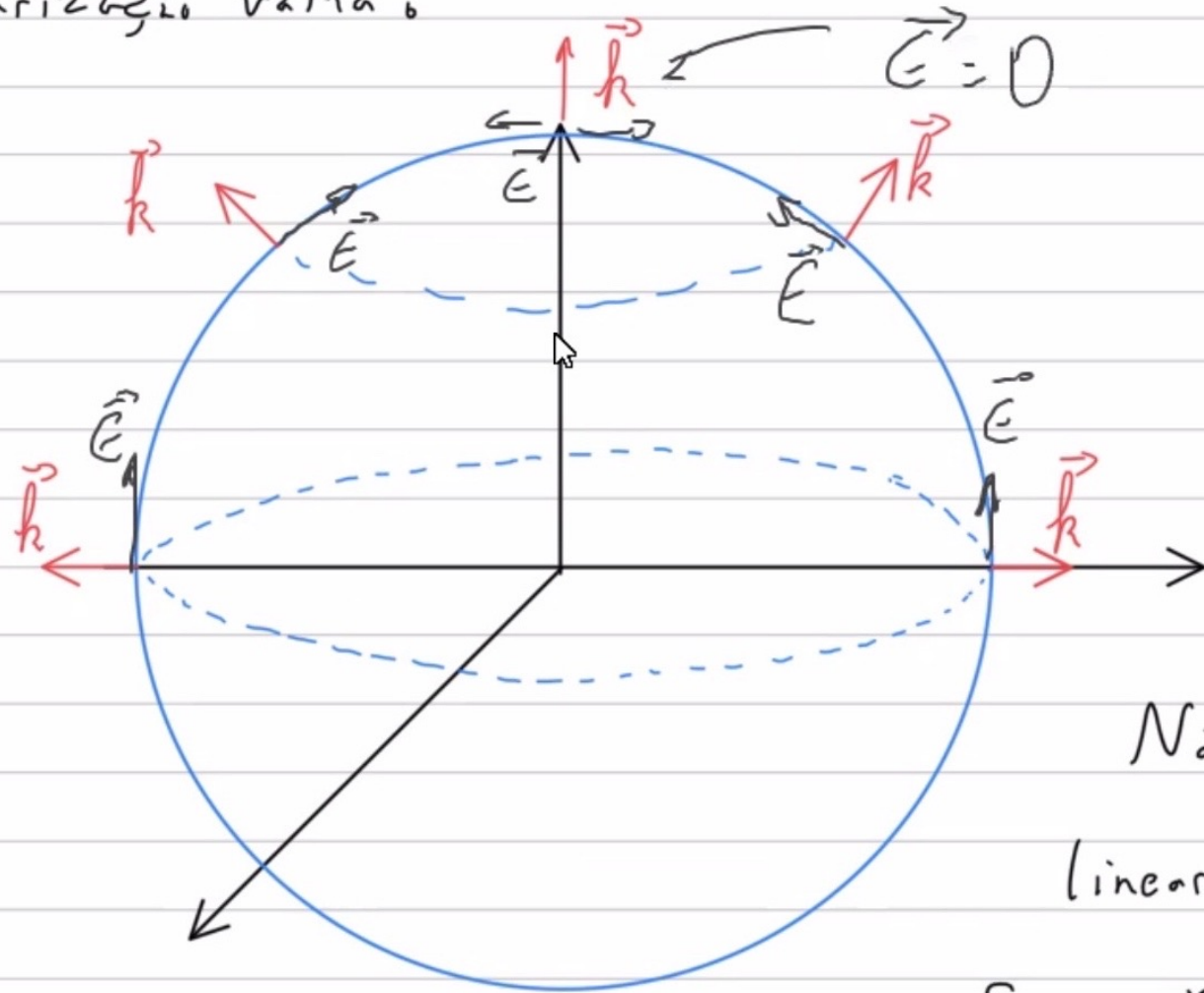


Não pode ser  
linear em toda a  
esfera  $\rightarrow \nabla \cdot \vec{E} \neq 0$  no pólo!

Veremos o truque mais tarde...

Ondas transversais:  $\vec{E} = \frac{E}{r} \hat{n} e^{i(\vec{k} \cdot \vec{r} - \omega t)}$  ?

polarização varia!



Não pode ser

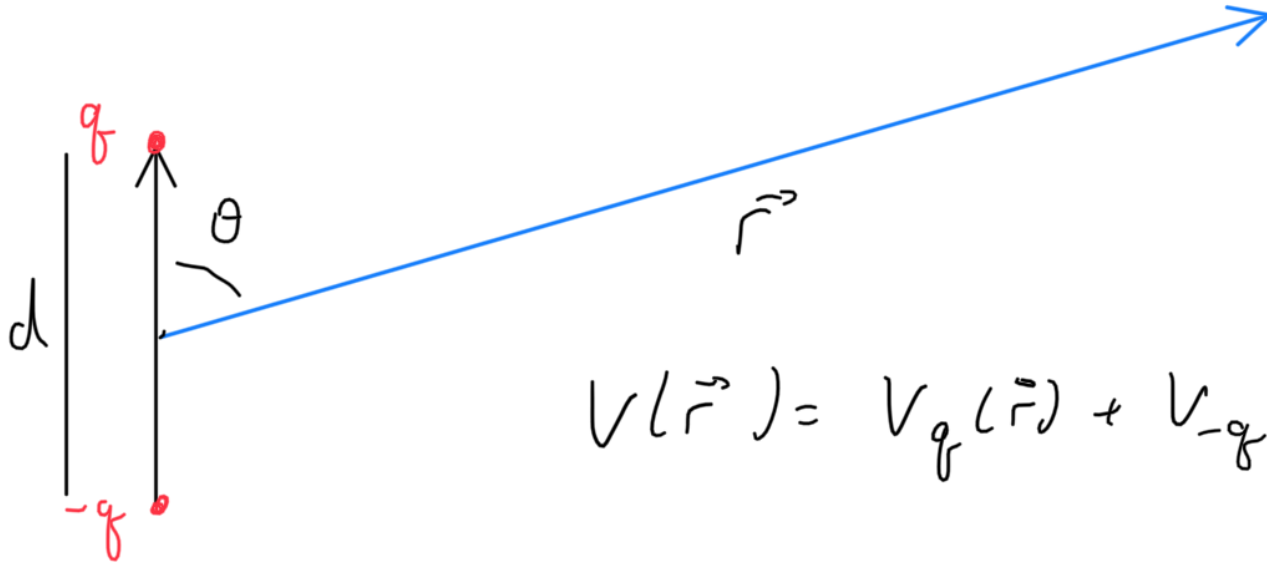
linear em toda a

esfera  $\rightarrow \nabla \cdot \vec{E} \neq 0$  no polo!

Veremos o truque mais tarde...



# Potencial elétrico de um dipolo



$$V(\vec{r}) = V_q(\vec{r}) + V_{-q}(\vec{r})$$

$$= \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{|\vec{r} - \vec{d}/2|} - \frac{1}{|\vec{r} + \vec{d}/2|} \right]$$

$$= \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{|\vec{r} - \vec{d}/2|} - \frac{1}{|\vec{r} + \vec{d}/2|} \right]$$

$$(\vec{r} - \vec{d})^2 = r^2 \sin^2 \theta + (r \cos \theta - \frac{d}{2})^2 = r^2 \sin^2 \theta + r^2 \cos^2 \theta - d r \cos \theta + \frac{d^2}{4}$$

$$= r^2 + \frac{d^2}{4} - d r \cos \theta$$

$$(\vec{r} + \vec{d})^2 = r^2 + \frac{d^2}{4} + d r \cos \theta$$

$$r \gg d \quad (\vec{r} - \vec{d})^2 \approx r^2 \left( 1 - \frac{d}{r} \cos \theta \right)$$

$$(\vec{r} + \vec{d})^2 \approx r^2 \left( 1 + \frac{d}{r} \cos \theta \right)$$

$$V(\vec{r}) = \frac{q}{4\pi\epsilon_0 r} \left[ \frac{1}{\sqrt{1 - \frac{d}{r} \cos \theta}} - \frac{1}{\sqrt{1 + \frac{d}{r} \cos \theta}} \right]$$

$$V(\vec{r}) = \frac{q}{4\pi\epsilon_0 r} \left[ \frac{1}{\sqrt{1 - \frac{d}{r} \omega \cos\theta}} - \frac{1}{\sqrt{1 + \frac{d}{r} \omega \cos\theta}} \right]$$

$$\frac{1}{\sqrt{1-\alpha}} \approx 1 - \frac{1}{2}\alpha$$

$$\Rightarrow V(\vec{r}) = \frac{q}{4\pi\epsilon_0 r} \cdot \frac{d \omega \cos\theta}{r} = \frac{p \omega \cos\theta}{4\pi\epsilon_0 r^2} = \frac{\vec{p} \cdot \vec{r}}{4\pi\epsilon_0 r^3}$$

$$\text{com } \vec{p} = q \cdot \vec{d}$$

Lembra o que vimos na aula passada:

esfera metálica em campo elétrico

$$\vec{A}(\vec{r}) \simeq \frac{\mu_0}{4\pi r^2} \int \vec{J}(\vec{r}') r' \sin\theta \, r'^2 dr' d\theta d\varphi$$

$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{r}}{r^2}$$

$$\vec{m} = \frac{1}{2} \int_V [\vec{r}' \times \vec{J}(\vec{r}')] dV$$

Potencial vetor  
de dipolo magnético

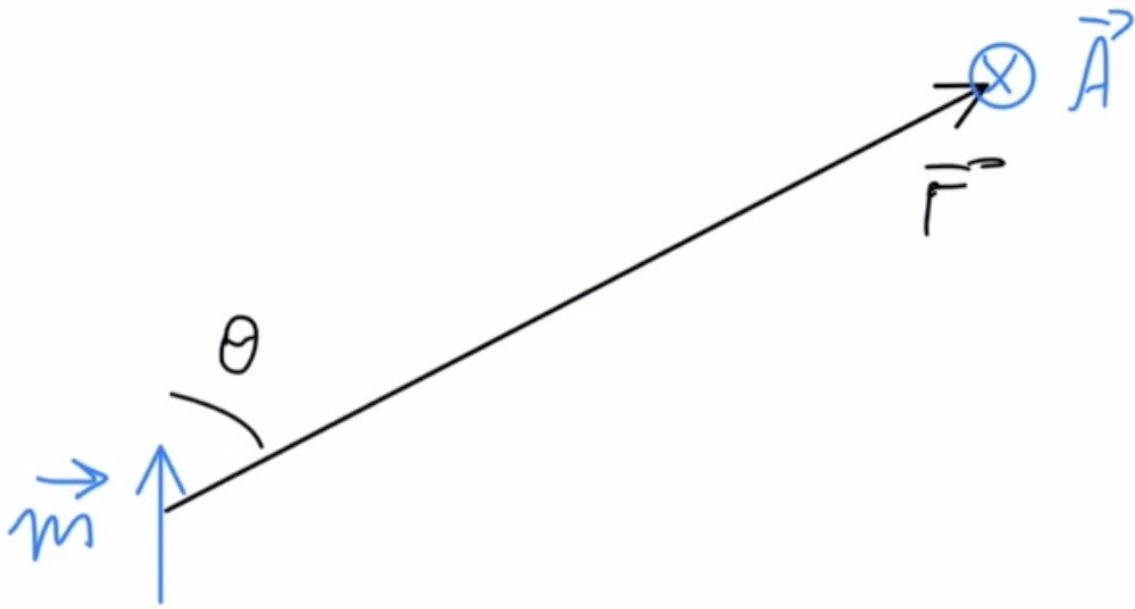
Momento de dipolo  
magnético

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \vec{r}}{r^2}$$

$$\vec{p} = \int \vec{r}' \cdot \rho(\vec{r}') dV$$

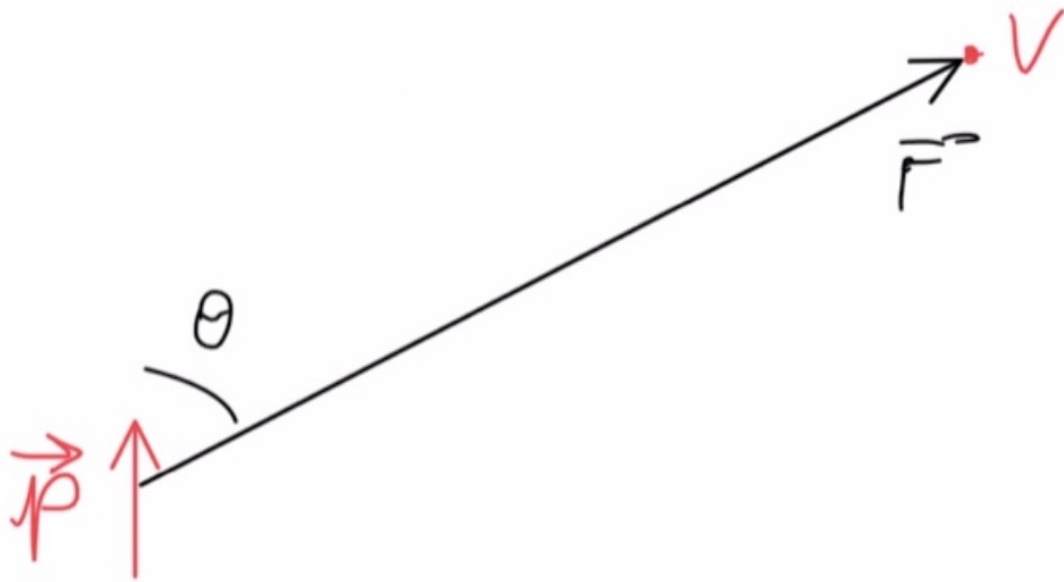
Potencial de dipolo elétrico

Momento de dipolo elétrico



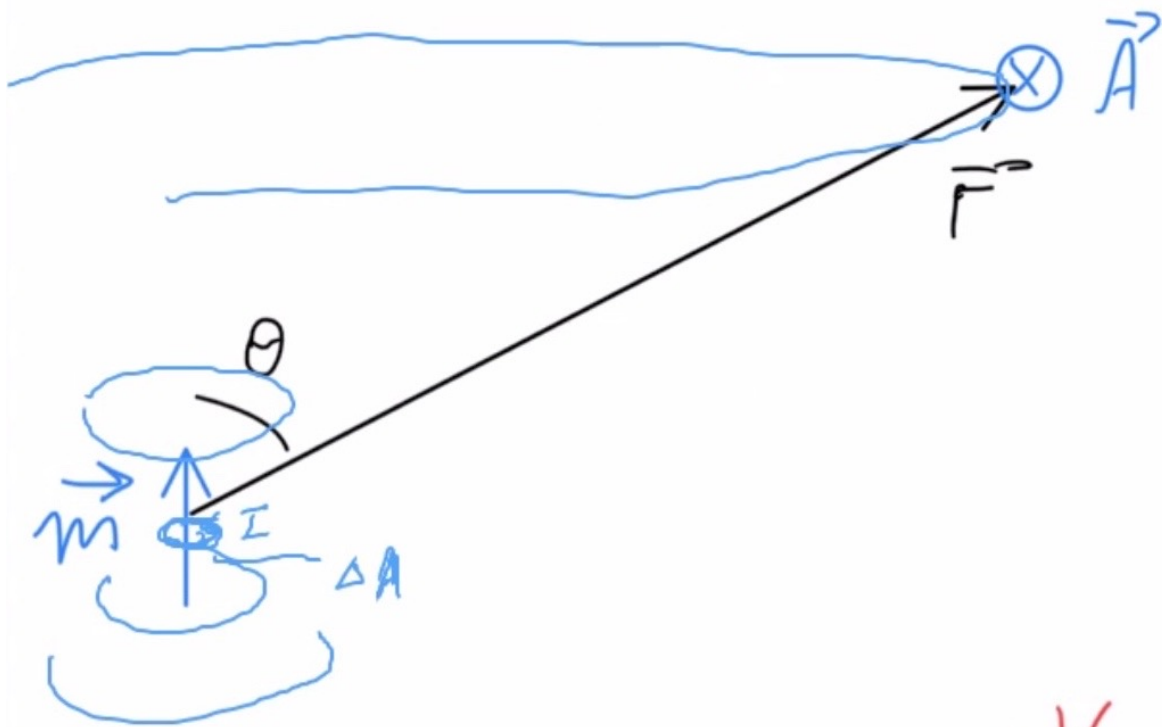
$$A = \frac{\mu_0 m}{4\pi r^2} \sin \theta$$

(anéis)



$$V = \frac{p}{4\pi \epsilon_0 r^2} \cos \theta$$





$$A = \frac{\mu_0 m}{4\pi r^2} \sin \theta$$

(anéis)



$$V = \frac{p}{4\pi \epsilon_0 r^2} \cos \theta$$

# Aula 27: Radiação

Fontes: Cargas em movimento

É mais fácil usar os potenciais:

$$\vec{E} = -\nabla V - \frac{\partial}{\partial t} \vec{A} \quad \vec{B} = \nabla \times \vec{A}$$

Calibre de Lorentz:  $\nabla \cdot \vec{A} = -\frac{1}{c^2} \frac{\partial}{\partial t} V$

$\Rightarrow$  Das eqs. de  
Maxwell

$$\left\{ \begin{array}{l} \nabla^2 V - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} V = -\frac{1}{\epsilon_0} \rho \\ \nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \vec{A} = -\mu_0 \vec{J} \end{array} \right.$$

Caso estacionário: eqs. de Poisson:  $\nabla^2 V = -\frac{\rho}{\epsilon_0}$ ;  $\nabla^2 \vec{A} = -\mu_0 \vec{J}$

$$\Rightarrow V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}')}{r} d\tau, \quad r = |\vec{r} - \vec{r}'|$$

↳ fonte  
↳ esp.  $\forall$

$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{r} d\tau$$

Cargas se movem: o potencial visto se refere à configuração antiga - tempo regresso  $t_r$

$$t_r = t - \frac{r}{c}$$

Generalizando:  $V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{r} d\tau$

$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t_r)}{r} d\tau$$

Generalizando: 
$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{r} dV$$

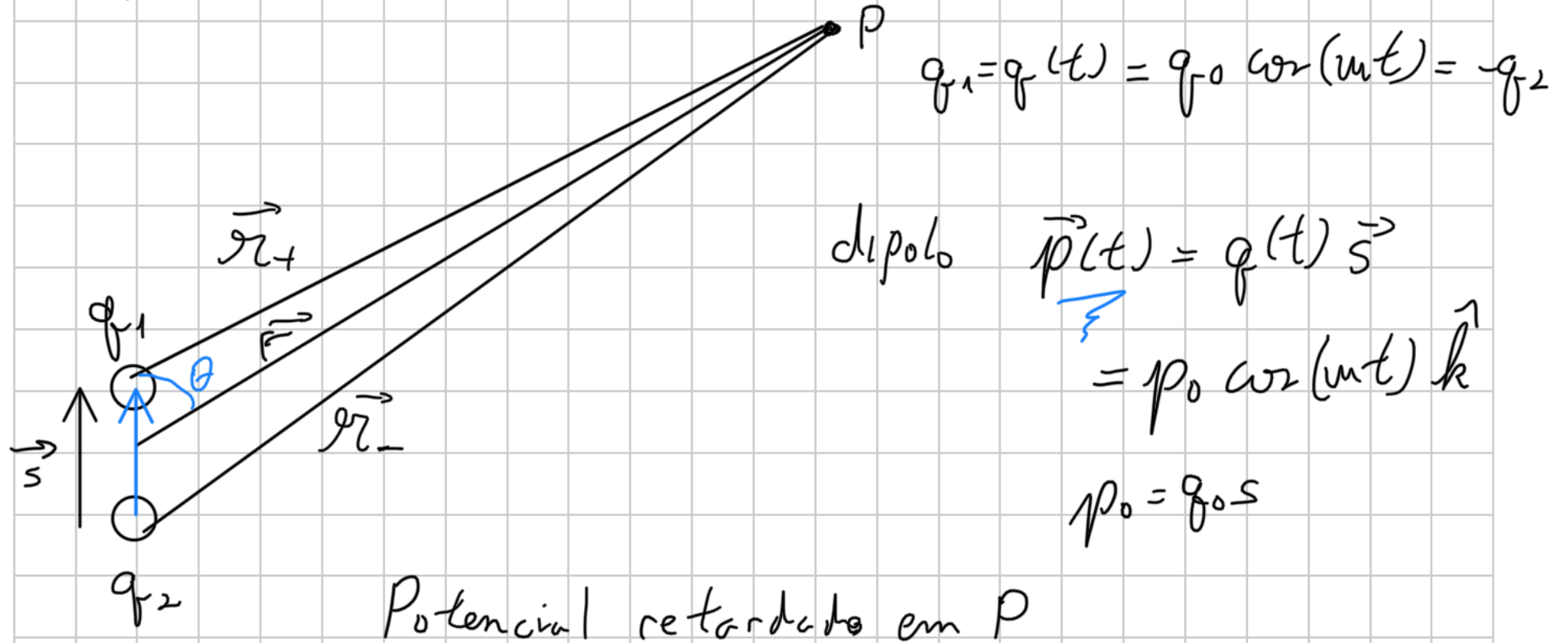
$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t_r)}{r} dV$$

Potenciais retardados de um tempo  $r/c$  no integrando

Precisamos conferir se isso satisfaz o Calibre de Lorentz

Leis de Coulomb e Biot-Savart não satisfazem!

# Aplicando ao potencial de dipolo



$$V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \left[ \frac{q_0 \cos(\omega(t - r_+/c))}{r_+} - \frac{q_0 \cos(\omega(t - r_-/c))}{r_-} \right]$$

$$r_{\pm} = \sqrt{r^2 \mp r s \cos \theta + (s/2)^2}$$

$$r_{\pm} = \sqrt{r^2 \mp r s \cos \theta + (s/2)^2}$$

Dipolo ideal:  $s \ll r \Rightarrow r_{\pm} \approx r \left( 1 \mp \frac{1}{2} \frac{s}{r} \cos \theta \right)$

$$\frac{1}{r_{\pm}} \approx \frac{1}{r} \left( 1 \pm \frac{1}{2} \frac{s}{r} \cos \theta \right)$$

$$\cos m \left( t - \frac{r_{\pm}}{c} \right) \approx \cos \left[ m \left( t - \frac{r}{c} \right) \pm \frac{m s \cos \theta}{2c} \right]$$

$$= \cos m \left( t - \frac{r}{c} \right) \cos \left( \frac{m s \cos \theta}{2c} \right) \mp \sin m \left( t - \frac{r}{c} \right) \sin \left( \frac{m s \cos \theta}{2c} \right)$$

$$\text{Fonte } \ll \text{ comprimento de onda} : s \ll \lambda = \frac{2\pi c}{\omega}$$

$$\Rightarrow \frac{ms}{c} \ll 1$$

$$\cos\left(\frac{ms}{2c} \omega \theta\right) \simeq 1; \quad \sin\left(\frac{ms}{2c} \omega \theta\right) \simeq \frac{ms}{2c} \omega \theta$$

$$\Rightarrow \cos m\left(t - \frac{r_{\pm}}{c}\right) \simeq \cos m\left(t - r/c\right) \pm \frac{ms}{2c} \omega \theta \sin m\left(t - r/c\right)$$

$$\frac{1}{r_{\pm}} \simeq \frac{1}{r} \left(1 \pm \frac{1}{2} \frac{s}{r} \cos \theta\right)$$

$$V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \left[ \frac{q_0 \cos m\left(t - r_+/c\right)}{r_+} - \frac{q_0 \cos m\left(t - r_-/c\right)}{r_-} \right]$$

$$\simeq \frac{q_0}{4\pi\epsilon_0} \left[ -\frac{ms}{c} \cos \theta \sin m\left(t - r/c\right) + \frac{s}{r^2} \cos \theta \cos m\left(t - r/c\right) \right]$$

$$V(\vec{r}, t) = \frac{\rho_0 \cos \theta}{4\pi \epsilon_0 r} \left[ -\frac{u}{c} \sin u(t - r/c) + \frac{1}{r} \cos u(t - r/c) \right]$$

$$u=0 \rightarrow V(\vec{r}) = \frac{\rho_0 \cos \theta}{4\pi \epsilon_0 r^2} \rightarrow \text{de volta ao começo } \nabla$$

Distante da fonte:  $r \gg \frac{c}{u} = 2\pi \lambda$


$$V(\vec{r}, t) = -\frac{\rho_0 u}{4\pi \epsilon_0 c} \left( \frac{\cos \theta}{r} \right) \sin u(t - r/c)$$



Potencial vector:  $\vec{d} = \frac{dq}{dt} \hat{k} = -q_0 v \text{ sen}(\omega t) \hat{k}$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int_{-s/2}^{s/2} \frac{-q_0 v \text{ sen}(\omega(t - r/c))}{r} \hat{k} dz$$

$$\approx \frac{\mu_0}{4\pi} \cdot \frac{-q_0 v \text{ sen}(\omega(t - r/c))}{r} \hat{k} \int_{-s/2}^{s/2} dz$$



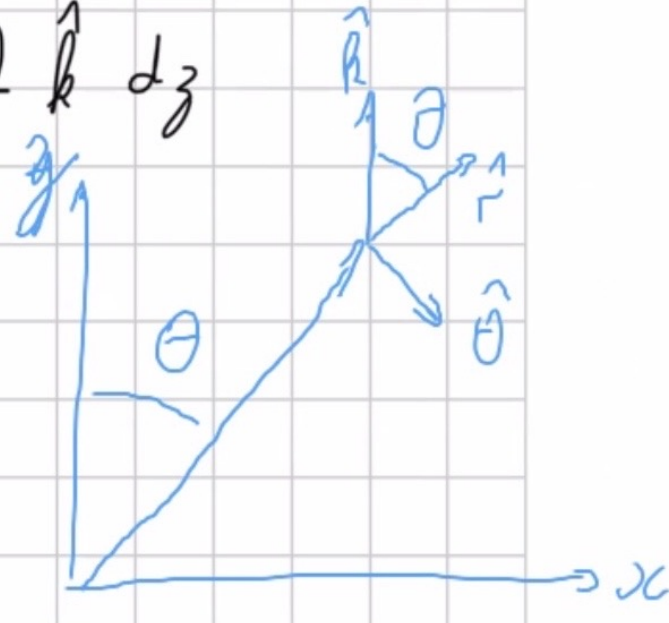
$$\vec{A}(\vec{r}, t) = -\frac{\mu_0 q_0 v}{4\pi r} \text{ sen}(\omega(t - r/c)) \hat{k}$$

$$\hat{k} = \cos \theta \hat{r} - \text{sen} \theta \hat{\theta}$$

Potencial vector:  $\vec{d} = \frac{dq}{dt} \hat{k} = -q_0 m \text{sen}(\omega t) \hat{k}$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int_{-s/2}^{s/2} \frac{-q_0 m \text{sen}(\omega(t - r/c))}{r} \hat{k} dz$$

$$\approx \frac{\mu_0}{4\pi} \cdot \frac{-q_0 m \text{sen}(\omega(t - r/c))}{r} \hat{k} \int_{-s/2}^{s/2} dz$$



$$\vec{A}(\vec{r}, t) = -\frac{\mu_0 p_0 m}{4\pi} \frac{1}{r} \text{sen}(\omega(t - r/c)) \hat{k}$$

$$\hat{k} = \cos \theta \hat{r} - \text{sen} \theta \hat{\theta}$$

Dados os potenciais, vamos aos campos!

$$\nabla V = \left( \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \right) V$$

$$= - \frac{\rho_0 \omega}{4\pi \epsilon_0 c} \left[ \hat{r} \frac{\partial}{\partial r} \left( \frac{\cos \theta}{r} \right) \sin \omega(t - r/c) + \hat{\theta} \frac{1}{r^2} \sin \omega(t - r/c) \frac{\partial}{\partial \theta} \cos \theta \right]$$

$$= - \frac{\rho_0 \omega}{4\pi \epsilon_0 c} \left\{ \cos \theta \left[ -\frac{1}{r^2} \sin \omega(t - r/c) - \frac{\omega}{rc} \cos \omega(t - r/c) \right] \hat{r} - \frac{\sin \theta}{r^2} \sin \omega(t - r/c) \hat{\theta} \right\}$$

$$\approx \frac{\rho_0 \omega^2}{4\pi \epsilon_0 c^2} \left( \frac{\cos \theta}{r} \right) \cos \omega(t - r/c) \hat{r}$$

$$\frac{\partial A}{\partial t} = - \frac{\mu_0 \rho_0 m}{4\pi r} \frac{\cos\theta \hat{r} - \sin\theta \hat{\theta}}{r} \frac{\partial}{\partial t} \sin m(t - r/c)$$

$$= - \frac{\rho_0 m^2}{4\pi \epsilon_0 c^2} \frac{\cos\theta \hat{r} - \sin\theta \hat{\theta}}{r} \cos m(t - r/c)$$

$$\Rightarrow \vec{E} = -\nabla V - \frac{\partial}{\partial t} \vec{A} = - \frac{\mu_0 \rho_0 m^2}{4\pi} \frac{\sin\theta}{r} \cos m(t - r/c) \hat{\theta}$$

$$\nabla \times \vec{A} = \frac{1}{r} \left[ \frac{\partial}{\partial r} (r A_\theta) - \frac{\partial}{\partial \theta} A_r \right] \hat{\varphi}$$

$$= - \frac{\rho_0 m}{4\pi \epsilon_0 c^2} \frac{1}{r} \left[ \frac{m}{c} \sin\theta \cos m(t - r/c) + \cancel{\frac{\sin\theta}{r}} \sin m(t - r/c) \right] \hat{\varphi}$$

$$\vec{B} = \nabla \times \vec{A} = - \frac{\mu_0 \rho_0 m^2}{4\pi c} \frac{\sin\theta}{r} \cos m(t - r/c) \hat{\varphi}$$

$$\vec{k} = \frac{m}{c} \hat{r}$$

$$\hat{\theta} \times \hat{\varphi} = \hat{k}$$

$$\vec{E} = -\frac{\mu_0 p_0 m^2}{4\pi r} \frac{\sin \theta}{r} \cos(\vec{k} \cdot \vec{r} - \omega t) \hat{\theta}$$

$$\vec{B} = -\frac{\mu_0 p_0 m^2}{4\pi r c} \frac{\sin \theta}{r} \cos(\vec{k} \cdot \vec{r} - \omega t) \hat{\varphi}$$

Ondas esféricas moduladas por sen θ

Vetor de Poynting

$$\vec{S} = \frac{\vec{E} \times \vec{B}}{\mu_0} = \frac{\mu_0}{c} \left( \frac{p_0 m^2}{4\pi} \right)^2 \frac{\sin^2 \theta}{r^2} \cos^2(\vec{k} \cdot \vec{r} - \omega t) \hat{r}$$

$$\langle \vec{S} \rangle = \frac{\mu_0 p_0^2 m^4}{32 \pi^2 c} \frac{\sin^2 \theta}{r^2} \hat{r}$$

Potência irradiada pelo dipolo

$$\langle P \rangle = \int \langle \vec{S} \rangle \cdot d\vec{a} = \frac{\mu_0 p_0^2 \omega^4}{32 \pi^2 c} \int \frac{\sin^2 \theta}{r^2} r^2 \sin \theta d\theta d\varphi$$

$$\int_0^{2\pi} \int_0^{\pi} \sin^2 \theta \sin \theta d\theta d\varphi = \int_{-1}^1 (1-x^2) dx = \left( x - \frac{1}{3} x^3 \right) \Big|_{-1}^1 = 2 - \frac{2}{3} = \frac{4}{3}$$

$$x = \cos \theta \quad \frac{dx}{d\theta} = -\sin \theta$$

$$\langle P \rangle = \frac{\mu_0 p_0^2 \omega^4}{32 \pi^2 c} \cdot \frac{4}{3} \cdot 2\pi = \frac{\mu_0 p_0^2 \omega^4}{12 \pi c}$$

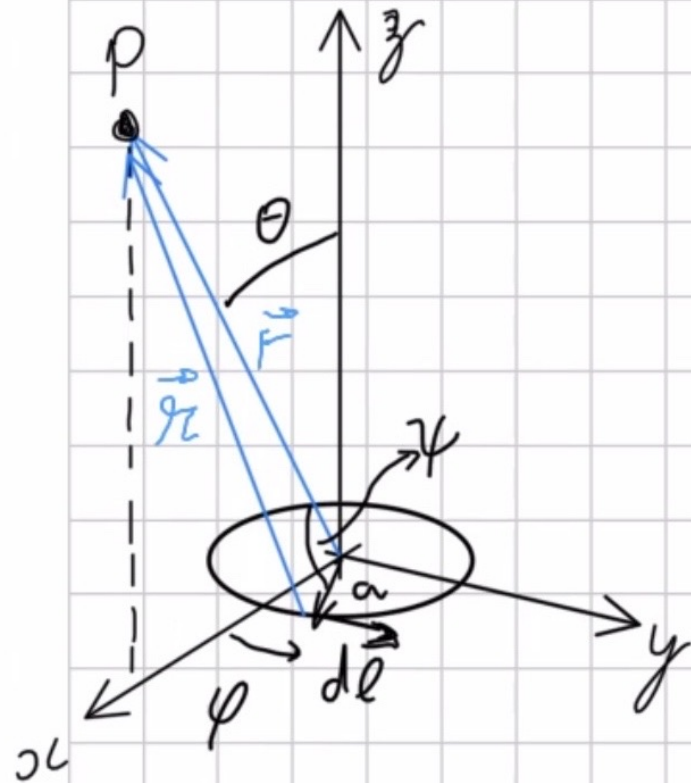
Qualquer dipolo oscilante vai ter dissipação por irradiação!

$\langle P \rangle \propto \omega^4 \rightarrow$  espalhamento maior em altas frequências

espalhamento Rayleigh  $\rightarrow$  céu azul



# Dipolo magnético



$$I = I_0 \cos(\omega t)$$

$$\vec{m}(t) = \pi a^2 I(t) \hat{k} = m_0 \cos \omega t \hat{k}$$

$$\rho = 0 \Rightarrow V = 0$$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t')}{r} dV'$$

$$= \frac{\mu_0}{4\pi} \int I_0 \frac{\cos \omega(t - r'/c)}{r} d\vec{\ell}'$$

Em  $\vec{P} = (x, y, z) = (x, 0, z)$  ;  $\vec{A} = A \hat{j}$  (componentes em  $i^1$  se cancelam)

$$\Rightarrow \vec{A}(\vec{r}, t) = \frac{\mu_0 I_0}{4\pi R} \cdot \hat{j} \cdot a \int_0^{2\pi} \frac{wr \sin(t - r/c)}{r^2} \cos\varphi d\varphi$$

$$r^2 = \sqrt{r^2 + a^2 - 2ra \cos\psi}$$

$$\vec{r} = r \sin\theta \hat{i} + r \cos\theta \hat{k}; \quad \vec{a} = a \cos\varphi \hat{i} + a \sin\varphi \hat{j}$$

$$\vec{r} \cdot \vec{a} = r \cdot a \cdot \cos\varphi = r a \sin\theta \cos\varphi$$

$$r^2 = \sqrt{r^2 - 2ra \cos\varphi \sin\theta + a^2} \approx r \left( 1 - \frac{a}{r} \cos\varphi \sin\theta \right)$$

se  $a \ll r$

$$\frac{1}{r^2} = \frac{1}{r} \left( 1 + \frac{a}{r} \cos\varphi \sin\theta \right)$$



$$\cos \omega(t - r/c) = \cos \left[ \omega(t - r/c) + \frac{\omega a \sin \theta \cos \varphi}{c} \right]$$

$$= \cos \omega(t - r/c) \cos \left( \frac{\omega a \sin \theta \cos \varphi}{c} \right) - \sin \omega(t - r/c) \sin \left( \frac{\omega a \sin \theta \cos \varphi}{c} \right)$$

$$\approx \cos \omega(t - r/c) - \sin \omega(t - r/c) \cdot \frac{\omega a \sin \theta \cos \varphi}{c}$$

$$\text{se } a \ll \lambda = \frac{2\pi c}{\omega}$$

$$\Rightarrow \vec{A}(\vec{r}, t) = \frac{\mu_0 I_0 a}{4\pi r} \hat{j} \left\{ \int_0^{2\pi} \cos \omega(t-r/c) \cos \varphi d\varphi \right.$$

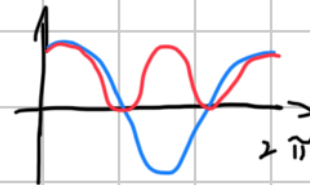
$$+ \int_0^{2\pi} \frac{a}{r} \sin \theta \cos \omega(t-r/c) \cos^2 \varphi d\varphi$$

$$\left. - \int_0^{2\pi} \frac{\mu_0 a}{c} \sin \theta \sin(\omega t - r/c) \cos^2 \varphi d\varphi \right\}$$

retendo termos em 1ª ordem

$$\int_0^{2\pi} \cos \varphi d\varphi = 0$$

$$\int_0^{2\pi} \cos^2 \varphi d\varphi = \pi$$



$$\Rightarrow \vec{A}(\vec{r}, t) = \frac{\mu_0 \tilde{\rho} a^2 I_0}{4\pi r} \left( \frac{\sin \theta}{r} \right) \left[ \frac{1}{r} \cos m(t - r/c) - \frac{m}{c} \sin m(t - r/c) \right] \hat{\varphi}$$

onde, por simetria, passamos  $\hat{j} \rightarrow \hat{\varphi}$

$$m=0 \Rightarrow \vec{A}(r, \theta) = \frac{\mu_0}{4\pi} \frac{m_0 \sin \theta}{r^2} \hat{\varphi}$$

$$r \gg \lambda = \frac{2\pi c}{\omega} \Rightarrow \vec{A}(r, \theta, t) = -\frac{\mu_0 m_0 m}{4\pi c} \frac{\sin \theta}{r} \sin \left( \omega t - \frac{\omega r}{c} \right) \hat{\varphi}$$

$$\Rightarrow \vec{E} = -\frac{\partial \vec{A}}{\partial t} = \frac{\mu_0 m_0 m^2}{4\pi c} \left( \frac{\sin \theta}{r} \right) \cos(\vec{k} \cdot \vec{r} - \omega t) \hat{\varphi}$$

$$\vec{B} = \nabla \times \vec{A} = -\frac{\mu_0 m_0 m^2}{4\pi c^2} \left( \frac{\sin \theta}{r} \right) \cos(\vec{k} \cdot \vec{r} - \omega t) \hat{\theta}$$

$$\vec{S} = \frac{\vec{E} \times \vec{B}}{\mu_0} = \frac{\mu_0}{c} \left[ \frac{m_0 m^2}{4\pi c} \left( \frac{\sin \theta}{r} \right) \omega (\vec{k} \cdot \vec{r} - \omega t) \right]^2 \hat{r}$$

$$\langle \vec{S} \rangle = \frac{\mu_0 m_0^2 m^4}{32 \pi^2 c^3} \left( \frac{\sin^2 \theta}{r^2} \right) \hat{r} ; \quad \langle P \rangle = \left( \frac{\mu_0 m_0^2}{12 \pi c^3} \right) m^4$$

$\vec{E}, \vec{B} \rightarrow$  trocam orientações  $\rightarrow$  similares!

Novo ponto:  $\frac{P_{\text{mag}}}{P_{\text{el}}} = \left( \frac{m_0}{\rho_0 c} \right)^2 = \left( \frac{\hat{\pi} a^2 I_0}{q_0 s} \cdot \frac{1}{c} \right)^2$

se  $I_0 \sim q_0 m$ ,  $s \sim \pi a$

$$\frac{P_{\text{mag}}}{P_{\text{el}}} = \left( \frac{a m}{c} \right)^2 = \left( 2\pi \frac{a}{\lambda} \right)^2$$

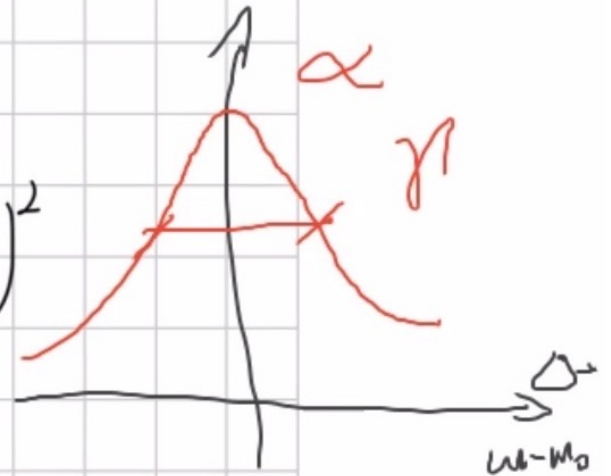
Acuplamento de dipolo magnético  $\ll$  Acuplamento dipolo elétrico!

$$\vec{S} = \frac{\vec{E} \times \vec{B}}{\mu_0} = \frac{\mu_0}{c} \left[ \frac{m_0 m^2}{4\pi c} \left( \frac{\sin \theta}{r} \right) \omega (\vec{k} \cdot \vec{r} - \omega t) \right]^2 \hat{r}$$

$$\langle \vec{S} \rangle = \frac{\mu_0 m_0^2 m^4}{32 \pi^2 c^3} \left( \frac{\sin^2 \theta}{r^2} \right) \hat{r} ; \quad \langle P \rangle = \left( \frac{\mu_0 m_0^2}{12 \pi c^3} \right) m^4$$

$\vec{E}, \vec{B} \rightarrow$  trocam orientações  $\rightarrow$  similares?

Nem tanto:  $\frac{P_{\text{mag}}}{P_{\text{el}}} = \left( \frac{m_0}{\rho_0 c} \right)^2 = \left( \frac{\hat{n} a^2 I_0}{q_0 s} \cdot \frac{1}{c} \right)^2$



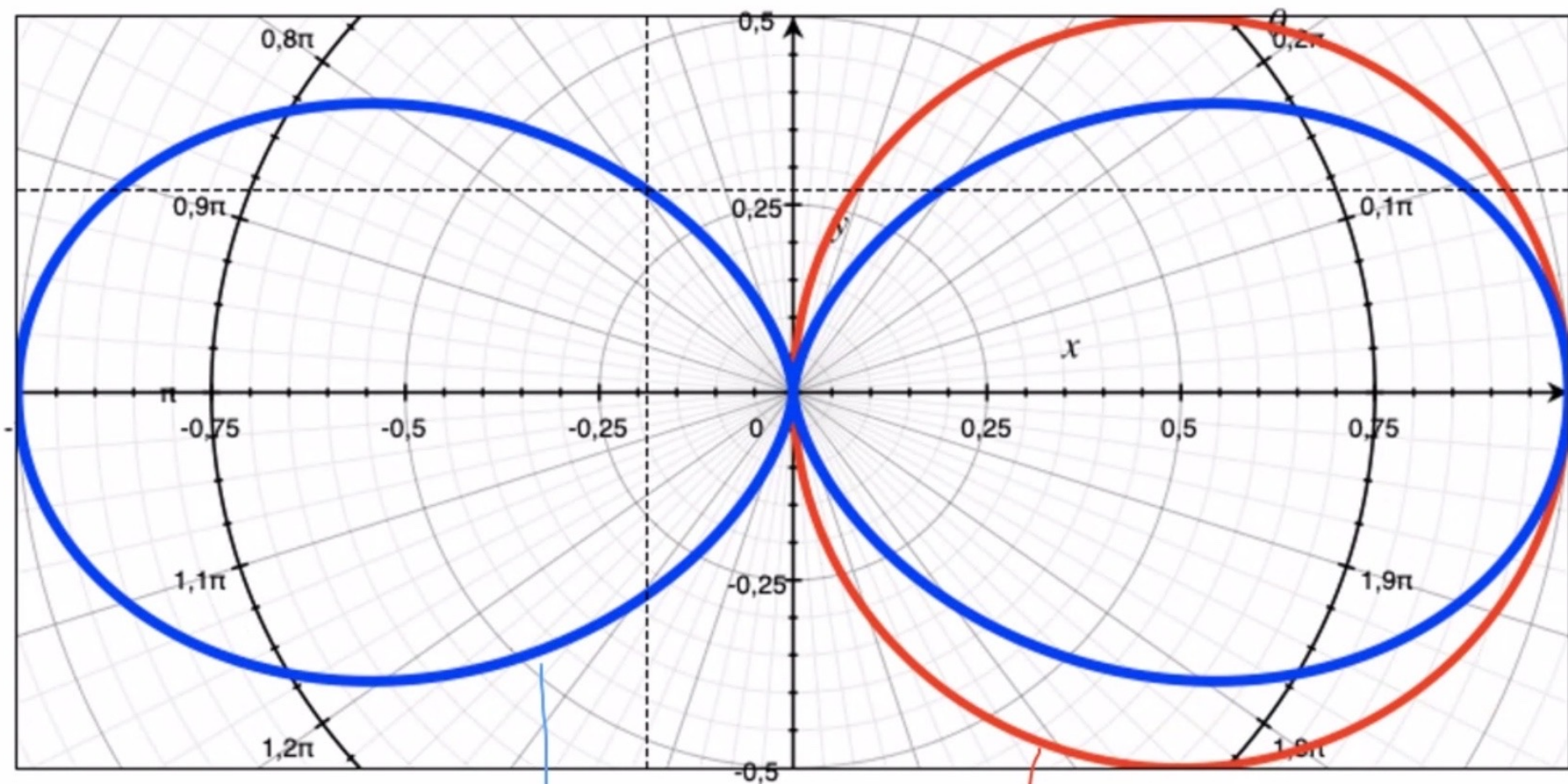
se  $I_0 \sim q_0 v$  ,  $s \sim \pi a$

$$\frac{P_{\text{mag}}}{P_{\text{el}}} = \left( \frac{a v}{c} \right)^2 = \left( 2\pi \frac{a}{\lambda} \right)^2$$

$$\frac{dE}{dt} = -\gamma E$$

$$\frac{P}{E} = \gamma$$

Acuplamentos de dipolo magnético  $\ll$  Acuplamentos dipolo elétrico!

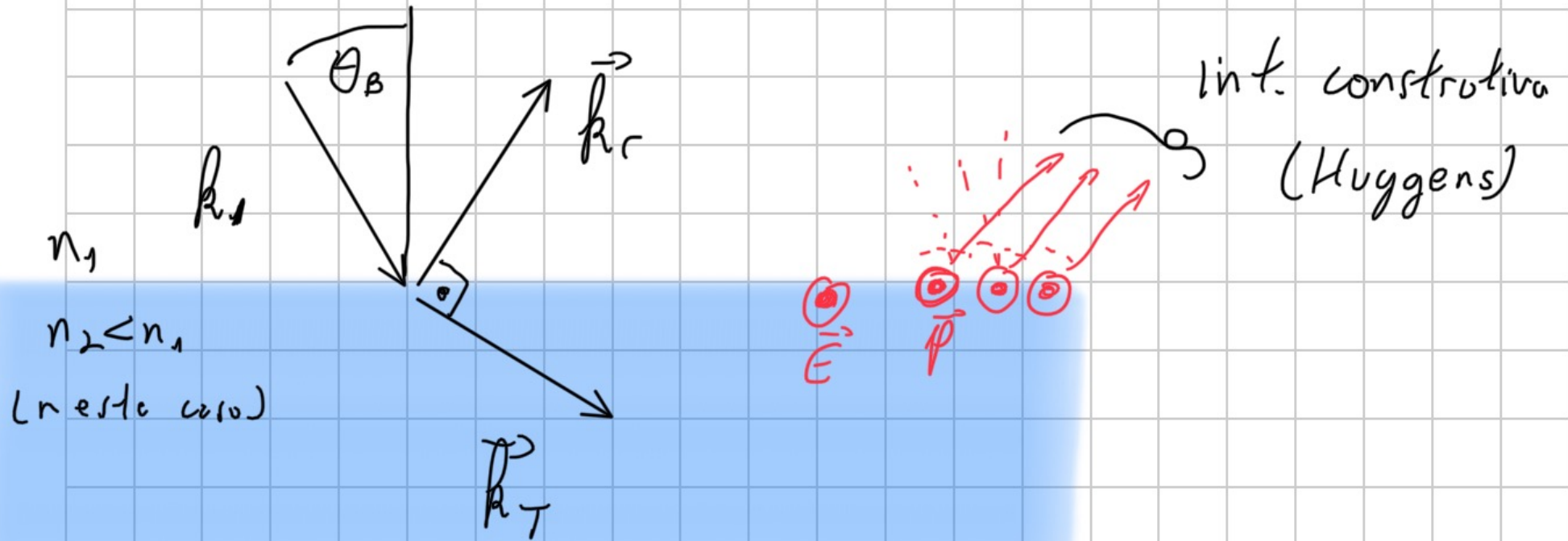


$|S|$

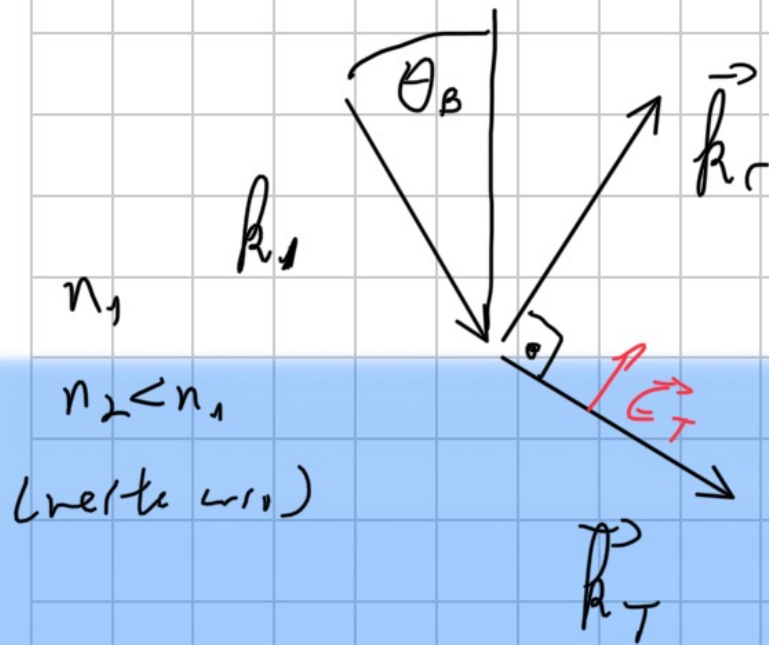
$|\vec{E}|, |\vec{B}|$



# Reflexão Polarizada



# Reflexão Polarizada



Int. construtiva

$\langle \vec{S} \rangle = 0 \rightarrow$  sem irradiação de dipolo



# Espejamiento Rayleigh

Sol  
 $\vec{k}_s$



$\downarrow |\vec{S}|_{\max}$



$\downarrow \vec{S} = 0$

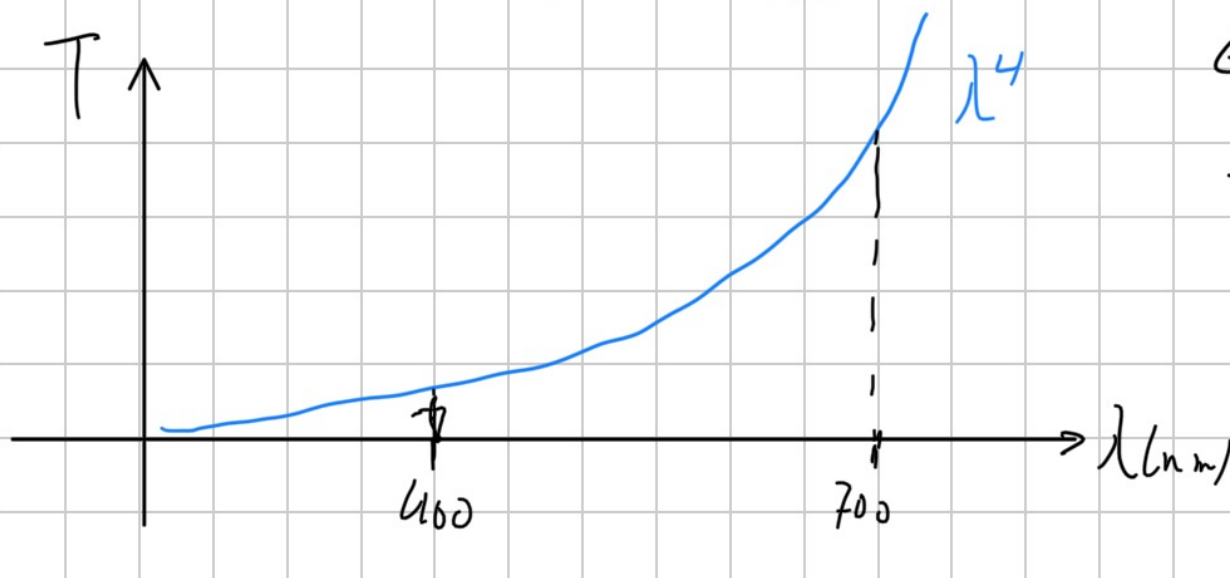
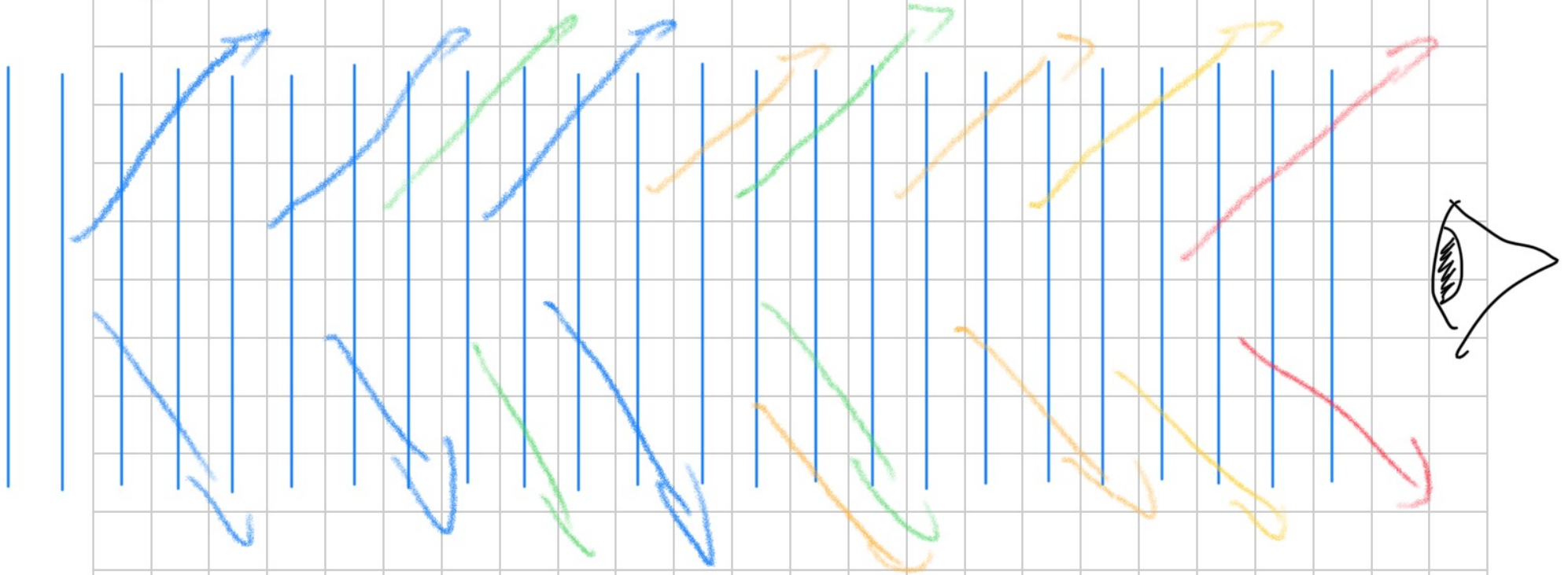


$\vec{E} \odot \downarrow \vec{k}_0$



$\rightarrow$  dipolos transversos

Por do Sol



Espectro de ondas  $\propto \omega^4 = 1 - T$

$$T = 1 - \kappa \cdot \omega^4$$

$$= 1 - \frac{\kappa' c^4}{\lambda^4}$$

$$= \frac{\lambda^4 - \kappa' c^4}{\lambda^4}$$