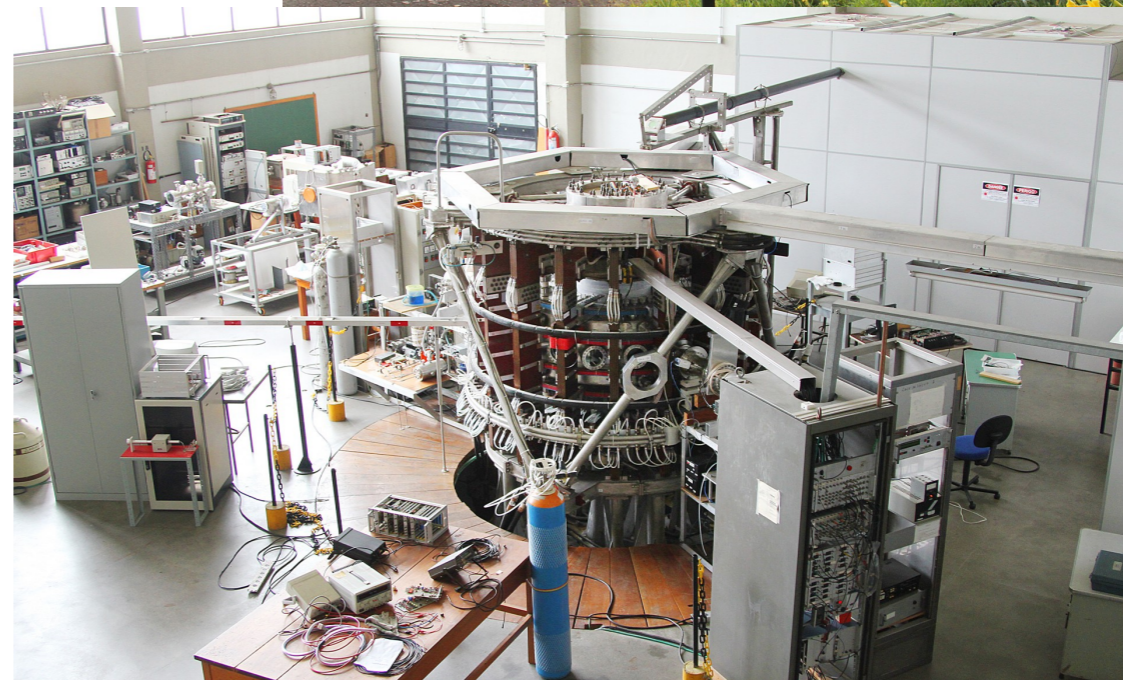


PGF5112 - Plasma Physics I

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São Paulo - SP, 30 March 2023



- **Single particle orbits: the motion of charged particles in electromagnetic fields**
 - *Introduction*
 - *Uniform and static electric field*
 - *Uniform and static magnetic field*
 - *Uniform and static electric and magnetic fields*
 - *Non-uniform and static magnetic field (physical insight)*
 - *Non-uniform and static electric field (physical insight)*
 - *Non-uniform and time-dependent electric and magnetic fields (next lecture)*

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The single particle orbit theory

- Knowing the trajectory of charged particles in special field configurations is important as it provides a good physical insight into some dynamic processes
- Here, we are interested in the motion of charged particles in the presence of electric (\mathbf{E}) and magnetic (\mathbf{B}) fields, which are known as functions of \mathbf{r} and t
 - Therefore, the fields are not affected by the charged particles
- The relativistic equation of motion for a charged particle under the action of the Lorentz force due to \mathbf{E} and \mathbf{B} fields is

$$\frac{d\mathbf{p}}{dt} = \mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

- Here, $\mathbf{p} = \gamma m \mathbf{v}$ is the relativistic particle momentum, with $\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$

q and m are the particle charge and rest mass, respectively

$c = 2.99 \times 10^8$ m/s is the speed of light in vacuum

The classical/non-relativistic single particle orbit theory

- **In many situations of practical interest, the term $v^2/c^2 \ll 1$**
 - Therefore, $\gamma \approx 1$ and m can be considered constant (independent of v)
- **Relativistic effects are important only for highly energetic particles**
 - A 1 MeV proton has $v = 1.4 \times 10^7$ m/s, i.e. $v^2/c^2 = 0.002 \ll 1$
 - Radiative effects, which are relativistic effects, will also be neglected here
- **In such a situations, the motion equation reduces to the non-relativistic equation**

$$m \frac{d\mathbf{v}}{dt} = q (\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

- **If the velocity obtained from this equation does not satisfy the condition $v^2/c^2 \ll 1$, then the relativistic equation of motion must be used instead**

The classical/non-relativistic single particle orbit theory

- Let's consider the particle kinetic energy

$$W = \frac{1}{2}mv^2 = \frac{1}{2}m \mathbf{v} \cdot \mathbf{v} \quad \rightarrow \quad \frac{dW}{dt} = \frac{d}{dt} \left(\frac{1}{2}m \mathbf{v} \cdot \mathbf{v} \right) = m \frac{d\mathbf{v}}{dt} \cdot \mathbf{v}$$

- Using the motion equation, this equation becomes

$$\frac{dW}{dt} = m \frac{d\mathbf{v}}{dt} \cdot \mathbf{v} = q (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \mathbf{v} = q (\mathbf{E} \cdot \mathbf{v}) + \cancel{q(\mathbf{v} \times \mathbf{B}) \cdot \mathbf{v}} = q (\mathbf{E} \cdot \mathbf{v})$$

- From the equation above, one concludes that
 - Any change in the particle kinetic energy is done by electric fields
 - Magnetic fields do no work on charged particles, i.e. the particle kinetic energy is conserved when there is only a magnetic field

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Single particle orbits: trajectories of charged particles in a static and uniform electric field

- Charged particles in externally applied electromagnetic fields are subject to the Lorentz force

$$\frac{d\mathbf{v}}{dt} = \frac{q}{m} (\mathbf{E}_{\text{ext}} + \mathbf{v} \times \mathbf{B}_{\text{ext}})$$

- For the case in which $\mathbf{B}_{\text{ext}} = 0$ and $\mathbf{E}_{\text{ext}} = \mathbf{E}_0$ is uniform and static, one has

$$\frac{d\mathbf{v}}{dt} = \frac{q}{m} \mathbf{E}_0$$

- The solution of this equation is obtained by direct integration

$$\int_0^t \frac{d\mathbf{v}}{dt} dt = \int_0^t \frac{q}{m} \mathbf{E}_0 dt \quad \rightarrow \quad \int_{\mathbf{v}_0}^{\mathbf{v}(t)} d\mathbf{v} = \frac{q}{m} \mathbf{E}_0 \int_0^t dt \quad \rightarrow \quad \mathbf{v}(t) = \mathbf{v}_0 + \frac{q}{m} \mathbf{E}_0 t$$

(Uniforme Rectilinear Motion)

$$\int_0^t \frac{d\mathbf{r}}{dt} dt = \int_0^t \mathbf{v}_0 dt + \frac{q}{m} \mathbf{E}_0 \int_0^t t dt \quad \rightarrow \quad \mathbf{r}(t) = \mathbf{r}_0 + \mathbf{v}_0 t + \frac{q}{2m} \mathbf{E}_0 t^2$$

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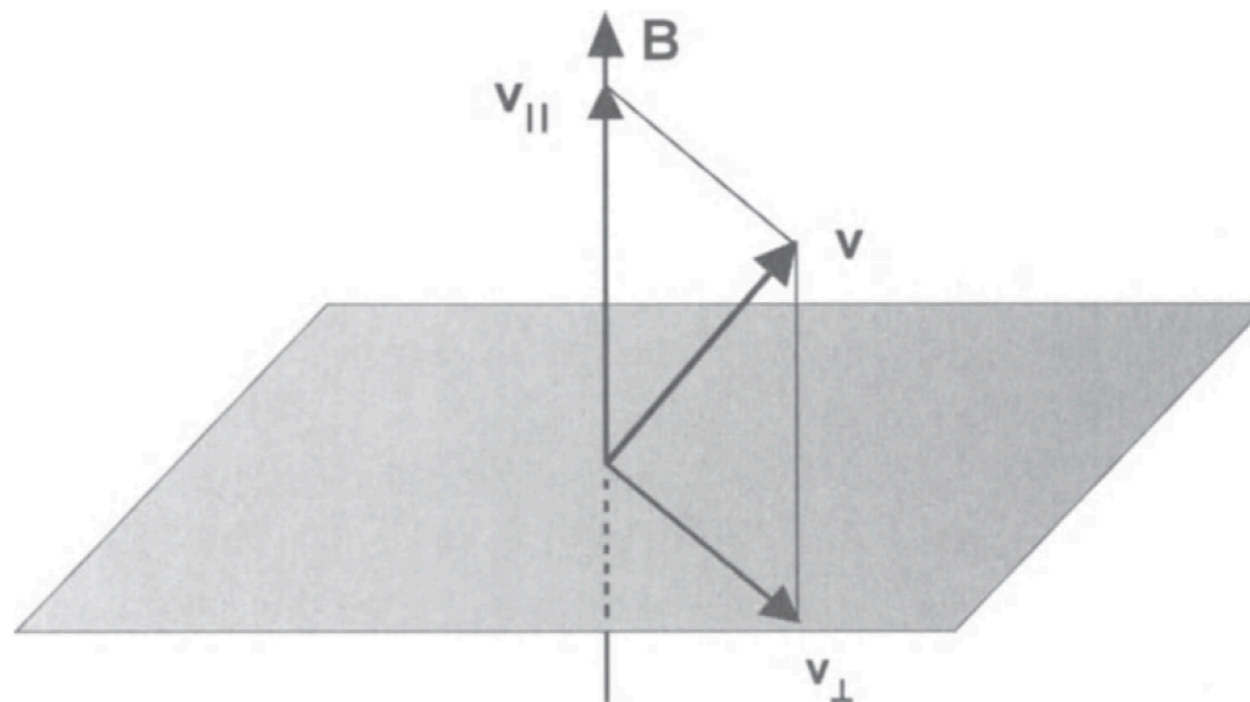
Single particle orbits: trajectories of charged particles in a static and uniform magnetic field

- For the case in which $\mathbf{E}_{\text{ext}} = 0$ and $\mathbf{B}_{\text{ext}} = \mathbf{B}_0$ is uniform and static, one has

$$\frac{d\mathbf{v}}{dt} = \frac{q}{m} \mathbf{v} \times \mathbf{B}_0$$

- Decompose \mathbf{v} in its parallel and perpendicular (to \mathbf{B}_0) components: $\mathbf{v} = \mathbf{v}_{\parallel} + \mathbf{v}_{\perp}$

$$\frac{d\mathbf{v}_{\parallel}}{dt} + \frac{d\mathbf{v}_{\perp}}{dt} = \frac{q}{m} \mathbf{v}_{\perp} \times \mathbf{B}_0 \quad (\text{The term } \mathbf{v}_{\parallel} \times \mathbf{B}_0 = 0 \text{ because } \mathbf{v}_{\parallel} \parallel \mathbf{B}_0)$$



Single particle orbits: trajectories of charged particles in a static and uniform magnetic field

- For the case in which $\mathbf{E}_{\text{ext}} = 0$ and $\mathbf{B}_{\text{ext}} = \mathbf{B}_0$ is uniform and static, one has

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- Decompose \mathbf{v} in its parallel and perpendicular (to \mathbf{B}_0) components: $\mathbf{v} = \mathbf{v}_{\parallel} + \mathbf{v}_{\perp}$

$$\frac{d\mathbf{v}_{\parallel}}{dt} + \frac{d\mathbf{v}_{\perp}}{dt} = \frac{q}{m} \mathbf{v}_{\perp} \times \mathbf{B}_0 \quad (\text{The term } \mathbf{v}_{\parallel} \times \mathbf{B}_0 = 0 \text{ because } \mathbf{v}_{\parallel} \parallel \mathbf{B}_0)$$

- In the parallel direction: uniform rectilinear motion

$$\frac{d\mathbf{v}_{\parallel}}{dt} = 0 \quad \rightarrow \quad \mathbf{v}_{\parallel} = \mathbf{v}_0$$

- In the perpendicular direction: cyclotron motion

$$\frac{d\mathbf{v}_{\perp}}{dt} = \frac{q}{m} \mathbf{v}_{\perp} \times \mathbf{B}_0$$

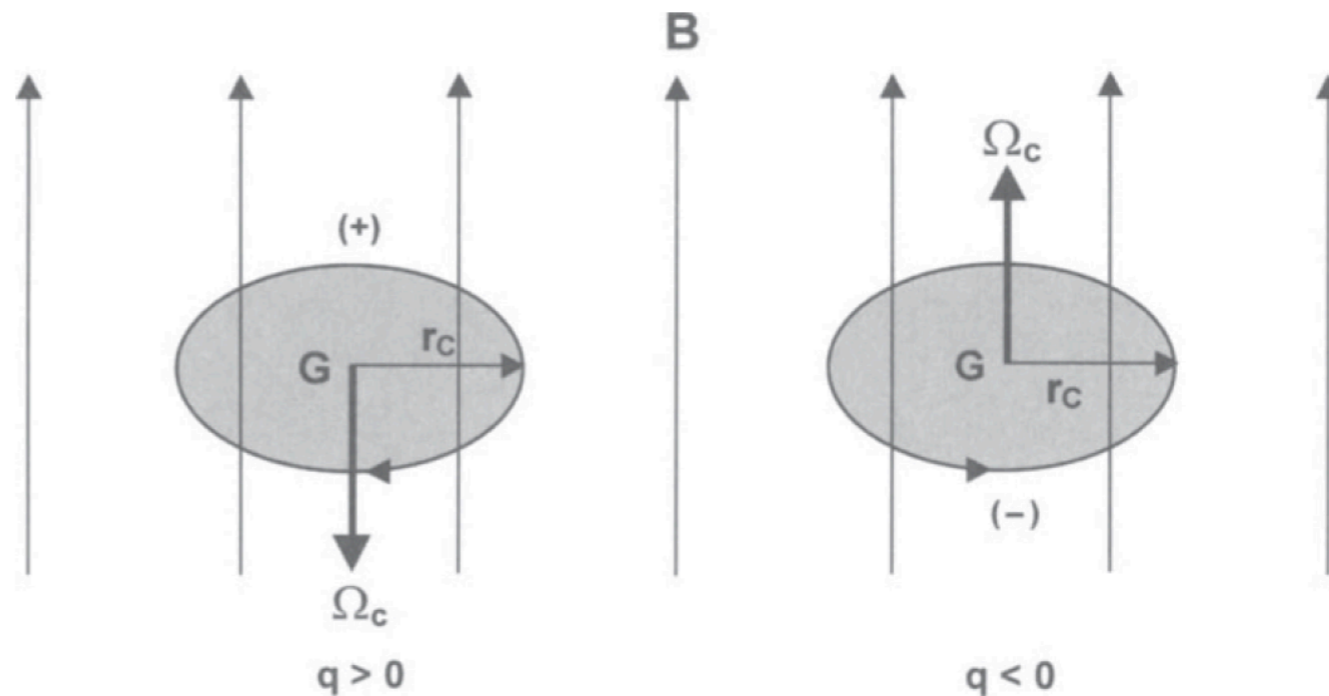
- If one defines $\mathbf{\Omega}_c = -\frac{q}{m} \mathbf{B}_0$, the motion equation becomes: $\frac{d\mathbf{v}_{\perp}}{dt} = \mathbf{\Omega}_c \times \mathbf{v}_{\perp}$

The cyclotron frequency (or the gyrofrequency)

- The quantity $\Omega_c = -\frac{q}{m}\mathbf{B}_0$ is called the cyclotron frequency or gyrofrequency

$$\frac{d\mathbf{v}_\perp}{dt} = \Omega_c \times \mathbf{v}_\perp$$

(Equation of Motion)



- Its direction is chosen as the diamagnetic direction
 - Ω_c is opposite to \mathbf{B}_0 for a positive charge ($q > 0$), which moves such that the magnetic field created by it is opposite to \mathbf{B}_0
 - Ω_c points in the direction of \mathbf{B}_0 for a negative charge ($q < 0$), which also moves such that the magnetic field created by it is opposite to \mathbf{B}_0
 - Note that Ω_c always points in the direction of the particle angular momentum

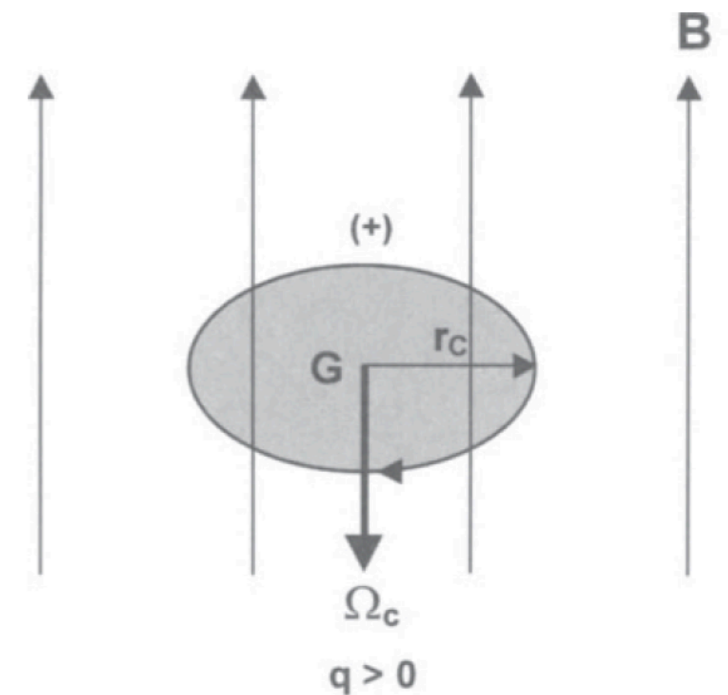
The Larmor radius (or the gyroradius)

- Since Ω_c is constant and, from kinetic energy conservation, $|\mathbf{v}_\perp|$ is also constant, the equation of motion implies that
 - The particle acceleration is constant in magnitude
 - The particle acceleration is perpendicular to both \mathbf{v}_\perp and \mathbf{B}_0

- The equation of motion can be integrated directly

$$\frac{d\mathbf{v}_\perp}{dt} = \mathbf{\Omega}_c \times \mathbf{v}_\perp = \mathbf{\Omega}_c \times \frac{d\mathbf{r}_c}{dt} \quad \rightarrow \quad \mathbf{v}_\perp = \mathbf{\Omega}_c \times \mathbf{r}_c \quad (\text{Circular Motion})$$

- Here, \mathbf{r}_c is the particle vector position measured with respect to a point (G in the figure) in the plane perpendicular to \mathbf{B}_0 and that contains the particle
- Since $|\mathbf{v}_\perp|$ is constant, $|\mathbf{r}_c|$ is also constant



- The vector \mathbf{r}_c is called the Larmor radius or gyroradius

The cyclotron frequency and Larmor radius for some particular cases

- **Larmor radius:** $r_c = v_{\perp} / \Omega_c = \frac{m v_{\perp}}{|q| B_0}$ One can estimate $v_{\perp} \approx v_{th} = \sqrt{\frac{k_B T}{m}}$

- **Cyclotron frequency:** $\Omega_c = |q| B_0 / m$

- Electron cyclotron: $f_{ce} = \Omega_{ce} / 2\pi = 28.0 \times B_0$ (GHz)

- Ion cyclotron: $f_{ci} = \Omega_{ci} / 2\pi = 15.2 \times B_0$ (MHz)

- **Cyclotron frequency and Larmor radius in plasmas**

- Tokamaks ($m_i = 1.67 \times 10^{-27}$ kg; $B_0 = 1.5$ T; $T = 1 \times 10^8$ K)

$$f_{ce} = 42 \text{ GHz}, f_{ci} = 22.8 \text{ MHz}, r_{ce} = 0.15 \text{ mm} \text{ and } r_{ci} = 6.3 \text{ mm}$$

- Solar corona ($m_i = 1.67 \times 10^{-27}$ kg; $B_0 = 0.1$ T; $T = 1 \times 10^6$ K)

$$f_{ce} = 2.8 \text{ GHz}, f_{ci} = 1.5 \text{ MHz}, r_{ce} = 0.22 \text{ mm} \text{ and } r_{ci} = 9.5 \text{ mm}$$

- **When $r_{ce}, r_{ci} \ll L$ (plasma size), the electrons/ions are said to be magnetized**

Single particle orbits: trajectories of charged particles in a static and uniform magnetic field

- The trajectory of the particle is given by the superposition of a uniform motion along B_0 and a cyclotron motion perpendicular to B_0

- The particle trajectory describes a helix

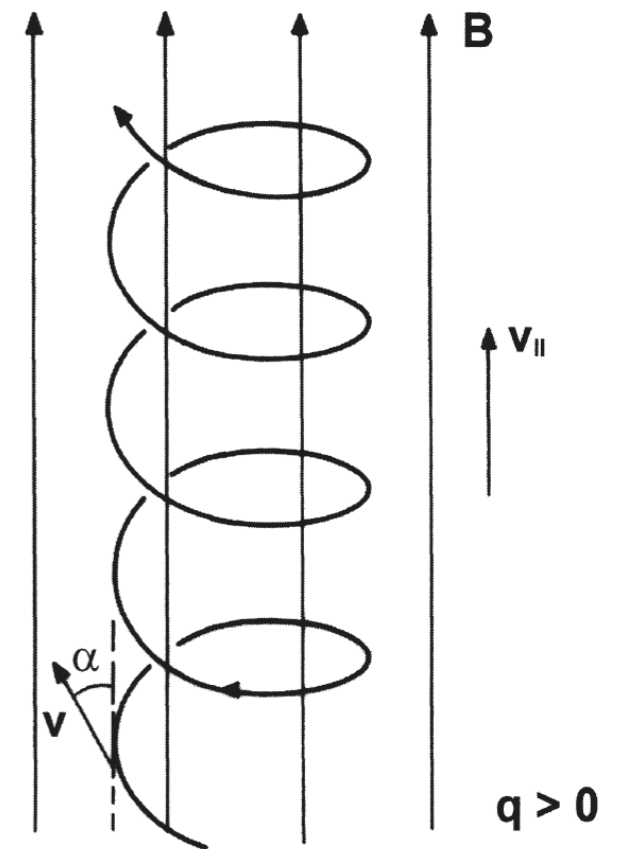
- The angle between B_0 and the direction of the particle motion is called the pitch angle

$$\alpha = \sin^{-1} \left(\frac{v_{\perp}}{v} \right) = \tan^{-1} \left(\frac{v_{\perp}}{v_{\parallel}} \right)$$

- Here, $v = \sqrt{v_{\parallel}^2 + v_{\perp}^2}$ is the total speed of the particle

- When $v_{\parallel} = 0$ and $v_{\perp} \neq 0$, $\alpha = \pi/2$ (Circular/Cyclotron Motion)

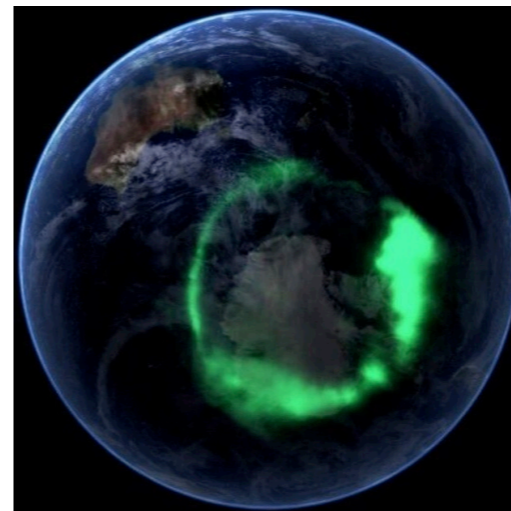
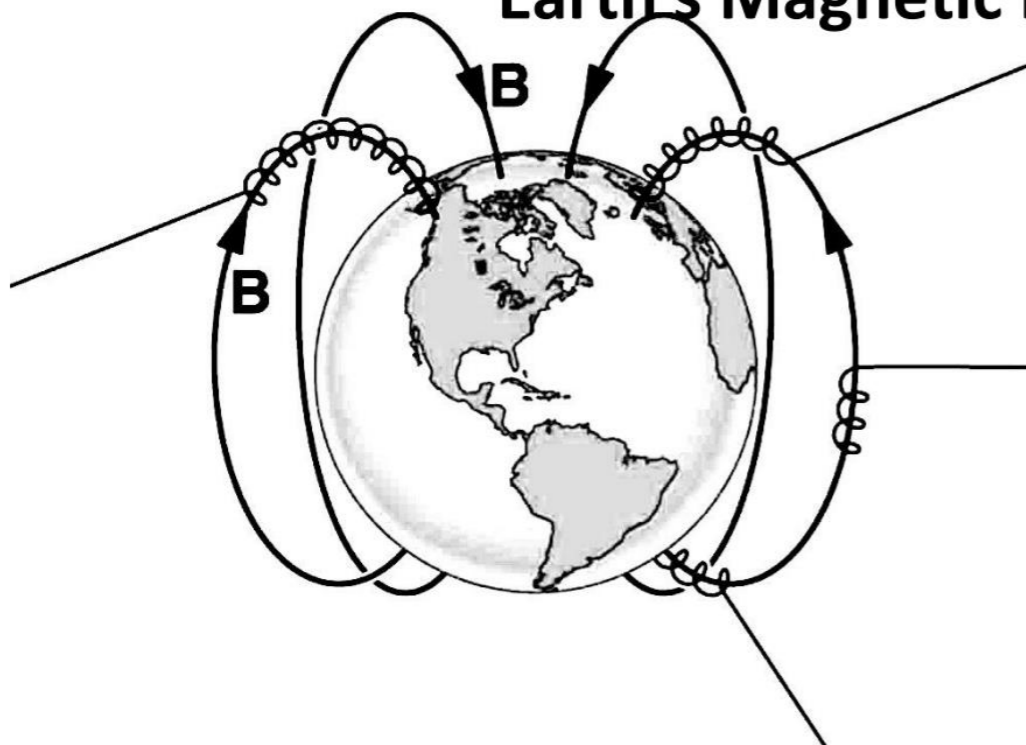
- When $v_{\parallel} \neq 0$ and $v_{\perp} = 0$, $\alpha = 0$ (Uniform Rectilinear Motion along the field line)



The particle orbit theory is well suited to study the trajectory of charged particles entering Earth's atmosphere

- **Charged particles arriving to Earth's atmosphere are deflected towards the poles by the terrestrial magnetic field**
 - The collision of these charged particles with the air molecules in the atmosphere gives rise to the so-called auroras

Charged Particle Trajectories in Earth's Magnetic Field



Magnetic moment associated with the cyclotron motion

- The magnetic moment due to a circulating current I is normal to the area A

$$|\mathbf{m}| = I A$$

- The current due to cyclotron motion is

$$I = \frac{|q|}{T_c} = \frac{|q|\Omega_c}{2\pi}$$

- The area of the current loop is

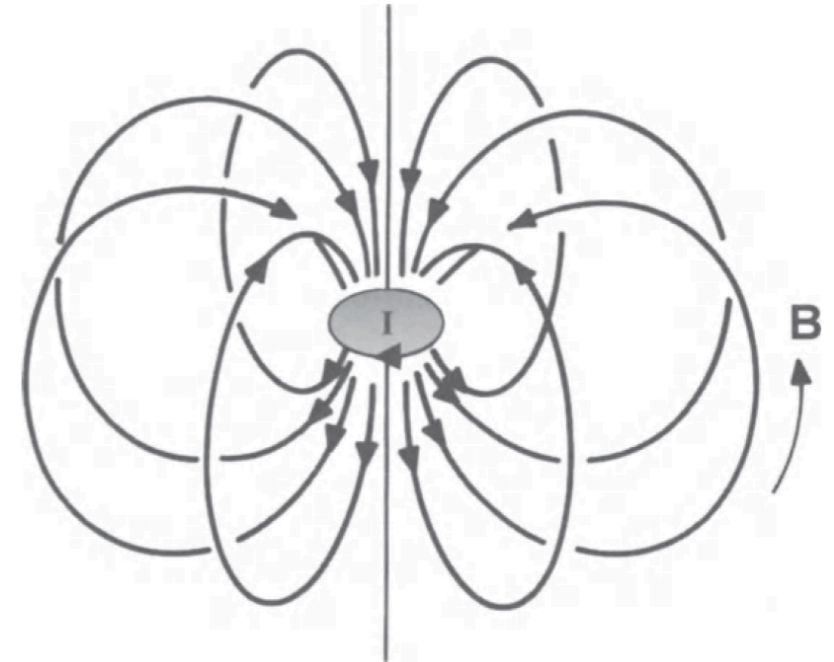
$$A = \pi r_c^2$$

- Therefore, the magnetic moment becomes

$$|\mathbf{m}| = \frac{|q|\Omega_c}{2\pi} \pi r_c^2 = \frac{1}{2} |q| \Omega_c r_c^2 = \frac{\frac{1}{2} m v_{\perp}^2}{B_0} = \frac{W_{\perp}}{B_0}$$

- In vector form, the magnetic moment associated with the cyclotron motion is

$$\mathbf{m} = -\frac{W_{\perp}}{B_0^2} \mathbf{B}_0$$



Magnetization current associated with the cyclotron motion

- The magnetization \mathbf{M} due to the cyclotron motion of several various particles is

$$\mathbf{M} = \frac{1}{V} \sum_{j=1}^N \mathbf{m}_j = \frac{N \mathbf{m}}{V} = n \mathbf{m} \quad \rightarrow \quad \mathbf{M} = -\frac{nW_{\perp}}{B_0^2} \mathbf{B}_0$$

- From classical electrodynamics, the magnetization current is $\mathbf{J}_M = \nabla \times \mathbf{M}$. Writing the total current density as $\mathbf{J}_{\text{total}} = \mathbf{J} + \mathbf{J}_M$, where \mathbf{J} is the current due to free charges, the Ampère-Maxwell equation becomes

$$\nabla \times \mathbf{B} = \mu_0 \left(\mathbf{J} + \mathbf{J}_M + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right) = \mu_0 \left(\mathbf{J} + \nabla \times \mathbf{M} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right)$$

$$\nabla \times \left(\frac{\mathbf{B}}{\mu_0} - \mathbf{M} \right) = \mathbf{J} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \quad \rightarrow \quad \nabla \times \mathbf{H} = \mathbf{J} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}, \text{ where } \mathbf{B} = \mu_0 (\mathbf{H} + \mathbf{M})$$

- A simple linear relation between \mathbf{B} and \mathbf{H} exists when \mathbf{M} is proportional to \mathbf{B} or \mathbf{H}
 - E.g. $\mathbf{M} = \chi_m \mathbf{H}$, where χ_m is the magnetic susceptibility of the medium
- In a plasma, however, $M \propto 1/B$ (non-linear). Therefore, it is NOT convenient to treat plasmas as magnetic media

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Single particle orbits: trajectories of charged particles in uniform and static electric and magnetic fields

- For the case in which $\mathbf{E}_{\text{ext}} = \mathbf{E}_0$ and $\mathbf{B}_{\text{ext}} = \mathbf{B}_0$ are uniform and static, one has

$$m \frac{d\mathbf{v}}{dt} = q (\mathbf{E}_0 + \mathbf{v} \times \mathbf{B}_0)$$

- Decompose \mathbf{v} and \mathbf{E}_0 in their parallel and perpendicular (to \mathbf{B}_0) components

$$m \frac{d\mathbf{v}_{\parallel}}{dt} + m \frac{d\mathbf{v}_{\perp}}{dt} = q (\mathbf{E}_{0,\parallel} + \mathbf{E}_{0,\perp} + \mathbf{v}_{\perp} \times \mathbf{B}_0)$$

- Parallel direction

$$m \frac{d\mathbf{v}_{\parallel}}{dt} = q \mathbf{E}_{0,\parallel} \quad \rightarrow \quad \mathbf{v}(t) = \mathbf{v}_{0\parallel} + \frac{q}{m} \mathbf{E}_{0,\parallel} t \quad \rightarrow \quad \mathbf{r}_{\parallel}(t) = \mathbf{r}_{0\parallel} + \mathbf{v}_{0\parallel} t + \frac{q}{2m} \mathbf{E}_{0,\parallel} t^2$$

- Perpendicular direction

$$m \frac{d\mathbf{v}_{\perp}}{dt} = q (\mathbf{E}_{0,\perp} + \mathbf{v}_{\perp} \times \mathbf{B}_0)$$

Single particle orbits: trajectories of charged particles in uniform and static electric and magnetic fields

- To solve the perpendicular equation, let's change referencial: $\mathbf{v}_\perp(t) = \mathbf{v}_c(t) + \mathbf{v}_{\mathbf{E} \times \mathbf{B}}$

$$m \frac{d\mathbf{v}_c}{dt} = q (\mathbf{E}_{0,\perp} + \mathbf{v}_c \times \mathbf{B}_0 + \mathbf{v}_{\mathbf{E} \times \mathbf{B}} \times \mathbf{B}_0)$$

- Choose the constant velocity $\mathbf{v}_{\mathbf{E} \times \mathbf{B}}$ as

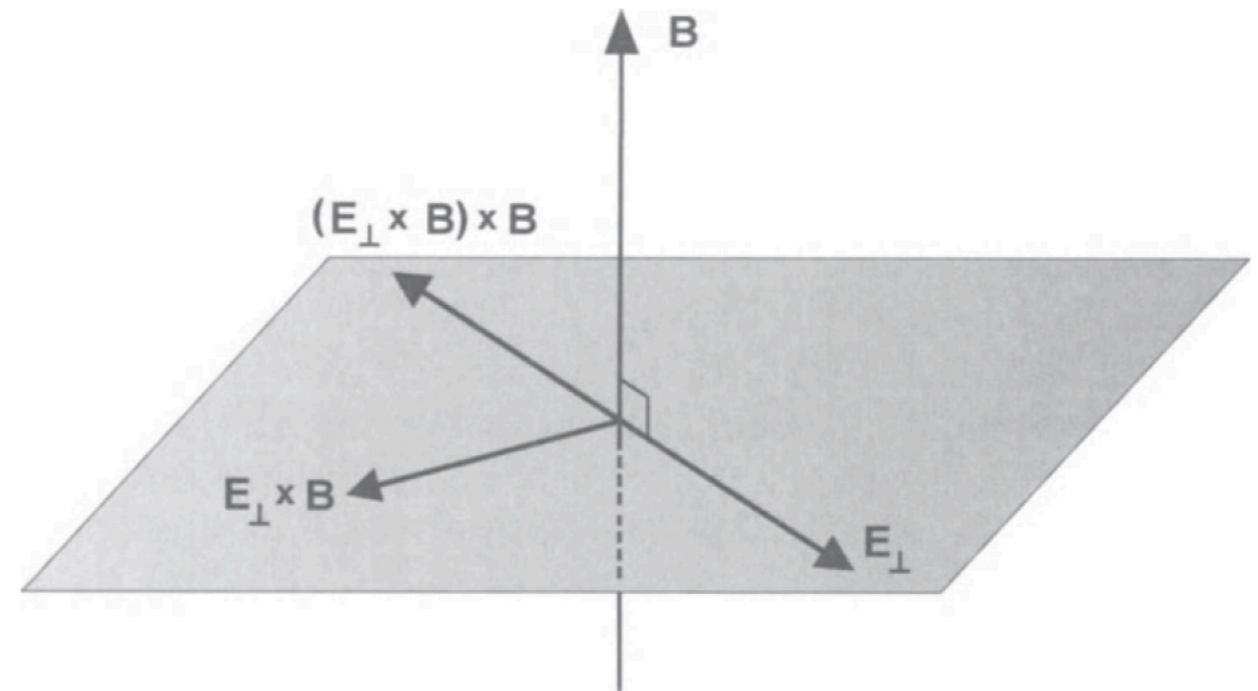
$$\mathbf{v}_{\mathbf{E} \times \mathbf{B}} = \frac{\mathbf{E}_{0,\perp} \times \mathbf{B}_0}{B_0^2}$$

- The equation of motion becomes

$$\frac{d\mathbf{v}_c}{dt} = \frac{q}{m} \mathbf{v}_c \times \mathbf{B}_0$$

- The solution of this equation is the cyclotron motion

$$\mathbf{v}_c(t) = \boldsymbol{\Omega}_c \times \mathbf{r}_c(t)$$

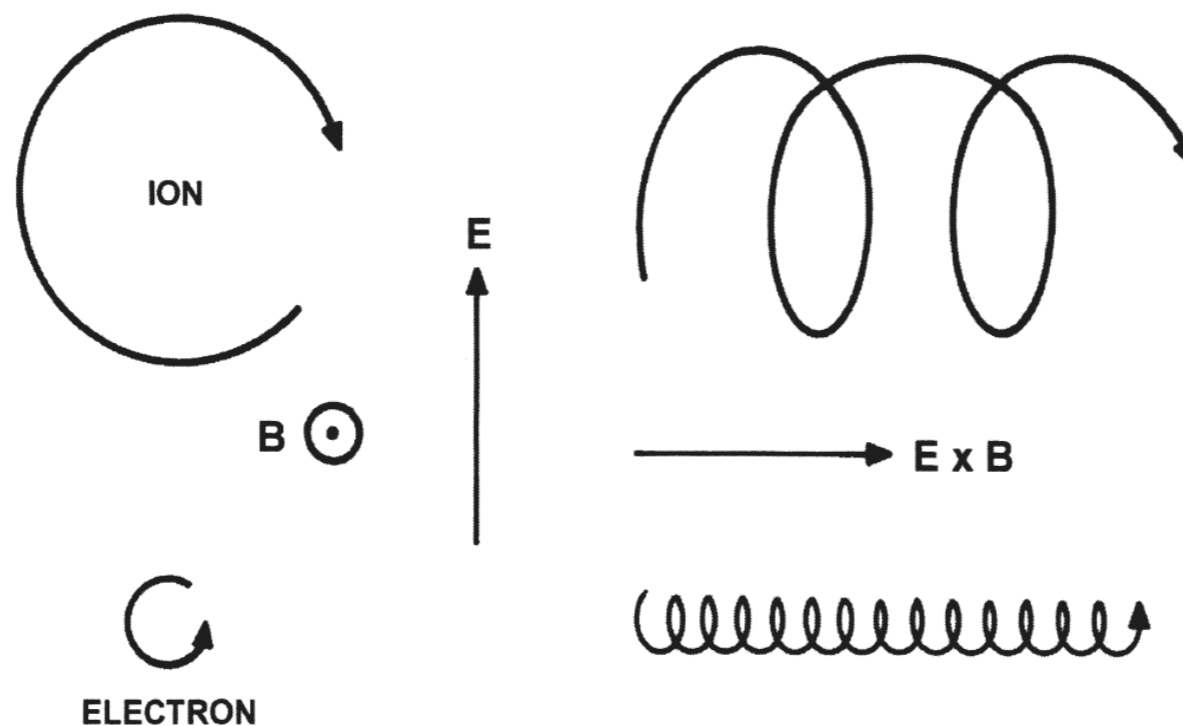


Single particle orbits: trajectories of charged particles in uniform and static electric and magnetic fields

- Therefore, the solution of this problem is

$$\mathbf{v}(t) = \boldsymbol{\Omega}_c \times \mathbf{r}_c(t) + \mathbf{v}_{\text{ExB}} + \mathbf{v}_{0\parallel} + \frac{q \mathbf{E}_{0,\parallel}}{m} t$$

- The constant velocity $\mathbf{v}_{\text{ExB}} = \mathbf{E}_{0,\perp} \times \mathbf{B}_0 / B_0^2$ is termed the ExB drift velocity
 - Note that \mathbf{v}_{ExB} is independent of the particle mass and charge
 - Since $\mathbf{E}_{0,\parallel} \times \mathbf{B}_0 = 0$, one can also write $\mathbf{v}_{\text{ExB}} = \mathbf{E}_0 \times \mathbf{B}_0 / B_0^2$



Drift due to an external force

- For the case in which, in addition to the EM fields, there is a force \mathbf{F} acting on the particle, the equation of motion becomes

$$m \frac{d\mathbf{v}}{dt} = q (\mathbf{E}_0 + \mathbf{v} \times \mathbf{B}_0) + \mathbf{F}$$

- The effect of the force is, in a formal sense, analogous to the effect of \mathbf{E}_0

$$\mathbf{v}_F = \frac{\mathbf{F} \times \mathbf{B}_0}{qB_0^2}$$

- In the case of a uniform gravitational field ($\mathbf{F} = m\mathbf{g}$), the drift velocity is

$$\mathbf{v}_F = \frac{m}{q} \frac{\mathbf{g} \times \mathbf{B}_0}{B_0^2}$$

- Associated to the gravitational drift, there is an electric current density

$$\mathbf{J}_g = \frac{1}{\delta V} \sum_j q_j \mathbf{v}_j = \frac{1}{\delta V} \left(\sum_j m_j \right) \frac{\mathbf{g} \times \mathbf{B}_0}{B_0^2} = \rho_m \frac{\mathbf{g} \times \mathbf{B}_0}{B_0^2}$$

- This current contributes to the so-called equatorial electrojet

Exercise

- **What happens with the \mathbf{ExB} drift velocity when the magnetic field tends to zero while the electric field remains finite? What is the validity of the \mathbf{ExB} drift expression?**

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Drift due to magnetic field gradient (physical insight)

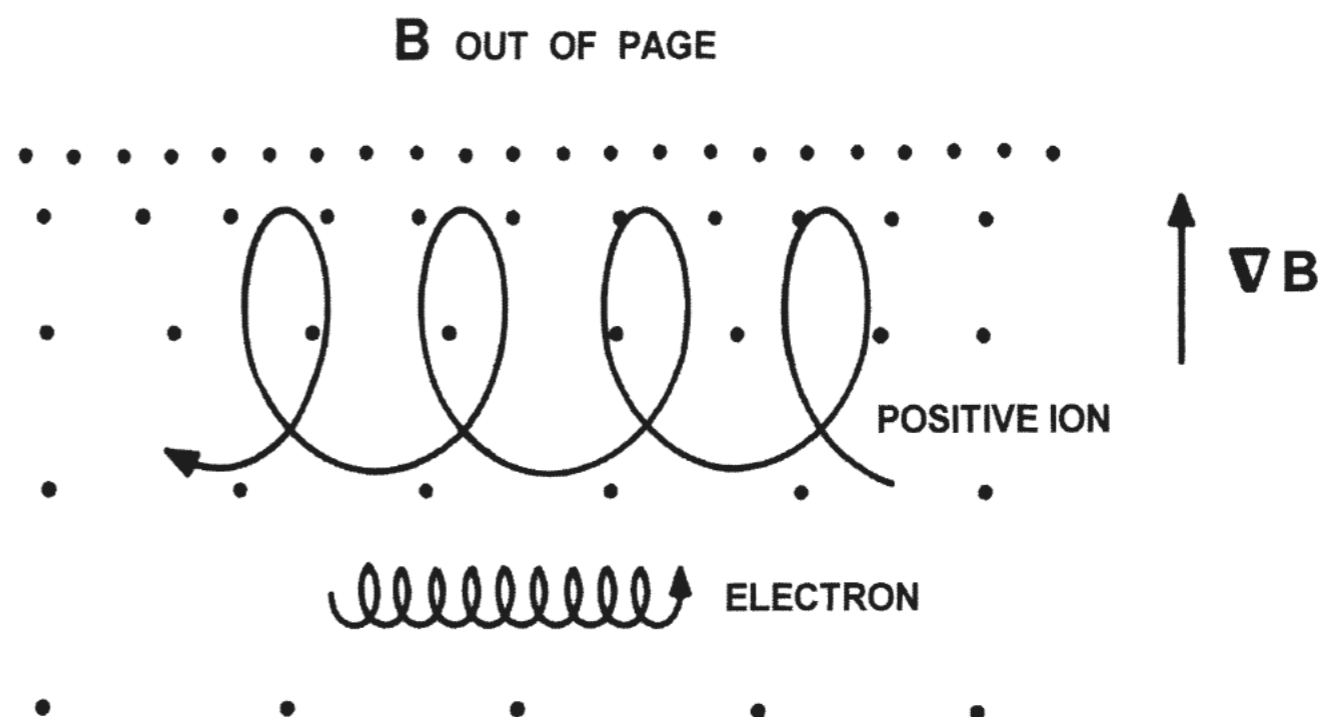
- One can expect that if the magnetic field varies over the Larmor radius, a drift velocity might arise

Larmor radius

$$r_c = \frac{mv_{\perp}}{|q|B_0}$$

Magnetic force due to ∇B_0

$$\langle \mathbf{F} \rangle_L = -|\mathbf{m}| \nabla B_0$$



- The magnetic drift associated to the gradient of the magnetic field (∇B_0) is

$$\mathbf{v}_{\nabla B} = \frac{\langle \mathbf{F} \rangle_L \times \mathbf{B}_0}{qB_0^2} = -\frac{|\mathbf{m}|}{q} \frac{\nabla B_0 \times \mathbf{B}_0}{B_0^2}$$

Drift due to magnetic field curvature (physical insight)

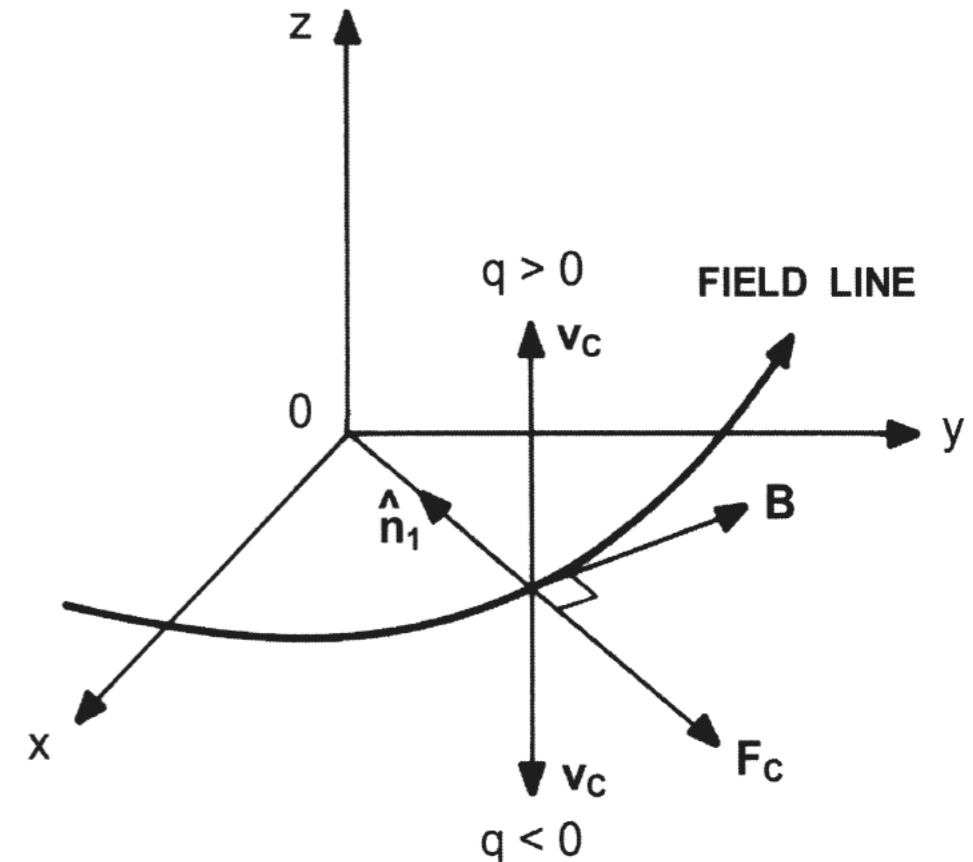
- One can also expect that if the magnetic field direction varies over the Larmor radius, a drift velocity might arise

Magnetic force due to the \mathbf{B}_0 -curvature

$$\langle \mathbf{F} \rangle_L = -\frac{mv_{\parallel}^2}{R} \hat{\mathbf{n}}_1$$

- Magnetic drift due to the \mathbf{B}_0 -curvature

$$\mathbf{v}_{\text{curv}} = \frac{\langle \mathbf{F} \rangle_L \times \mathbf{B}_0}{qB_0^2} = -\frac{mv_{\parallel}^2}{Rq} \frac{\hat{\mathbf{n}}_1 \times \mathbf{B}_0}{B_0^2}$$



Drift due to magnetic field curvature (physical insight)

- One can also expect that if the magnetic field direction varies over the Larmor radius, a drift velocity might arise

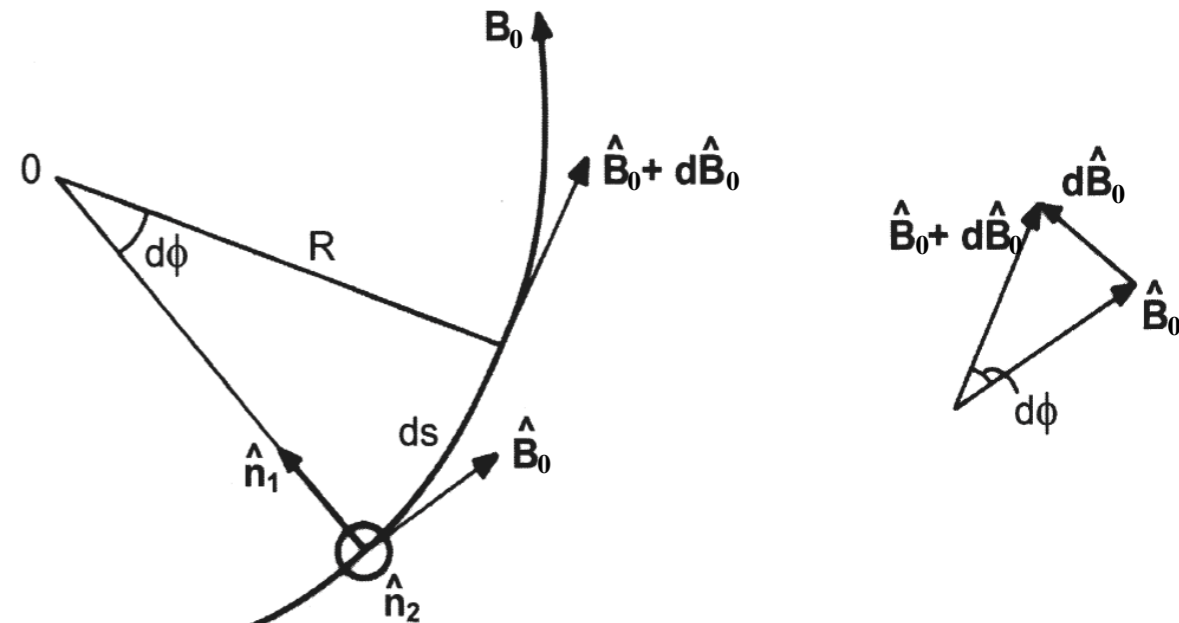
Magnetic force due to the \mathbf{B}_0 -curvature

$$\langle \mathbf{F} \rangle_L = -\frac{mv_{\parallel}^2}{R} \hat{\mathbf{n}}_1$$

- Magnetic drift due to the \mathbf{B}_0 -curvature

$$\mathbf{v}_{\text{curv}} = \frac{\langle \mathbf{F} \rangle_L \times \mathbf{B}_0}{qB_0^2} = -\frac{mv_{\parallel}^2}{Rq} \frac{\hat{\mathbf{n}}_1 \times \mathbf{B}_0}{B_0^2}$$

$$\mathbf{v}_{\text{curv}} = -\frac{mv_{\parallel}^2}{q} \frac{(\mathbf{B}_0 \cdot \nabla) \mathbf{B}_0 \times \mathbf{B}_0}{B_0^4}$$



$$ds = R d\phi \quad \hat{\mathbf{B}}_0 = \frac{\mathbf{B}_0}{B_0}$$

$$|d\hat{\mathbf{B}}_0| = |\hat{\mathbf{B}}_0| d\phi \quad d\hat{\mathbf{B}}_0 = \hat{\mathbf{n}}_1 d\phi$$

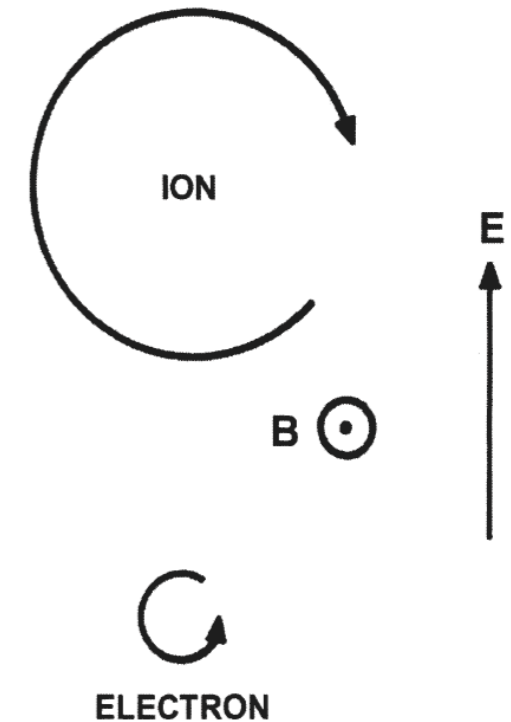
$$\frac{\hat{\mathbf{n}}_1}{R} = \frac{d\hat{\mathbf{B}}_0}{ds}$$

$$\frac{\hat{\mathbf{n}}_1}{R} = (\hat{\mathbf{B}}_0 \cdot \nabla) \hat{\mathbf{B}}_0 = \frac{(\mathbf{B}_0 \cdot \nabla) \mathbf{B}_0}{B_0^2}$$

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Drift due to electric field non-uniformities (physical insight)

- One can also expect that if the electric field varies over the Larmor radius, a drift velocity might arise
- At first order, in which the electric field varies linearly, a charged particle executing a cyclotron motion pass by a region with stronger E_0 -field and pass by a region with weaker E_0 -field
 - On average, the first order correction cancels out
 - Therefore, E_0 -field non-uniformities are important only as 2nd order corrections



References

- **The single particle orbit theory**
 - *Bittencourt: Ch. 2 and 3*