

$$\frac{d}{dz} \int [E^{(1)} \times H^{(2)} - E^{(2)} \times H^{(1)}] \cdot \hat{z} \, dx \, dy =$$

$$i\omega \int [E^{(2)} - E^{(1)}] E^{(1)} \cdot E^{(2)} \, dx \, dy \quad (1)$$

Passo 1:

Estrutura Perturbada:

$$\left\{ \begin{array}{l} E^{(2)} = E(x, y) \quad (2) \\ E^{(2)} = (E_t^T - \hat{z} E_z^T) e^{-i\beta_T z} \quad (3) \\ H^{(2)} = (-H_t^T + \hat{z} H_z^T) e^{-i\beta_T z} \quad (4) \end{array} \right.$$

↖ Modos Estrutura (z) propagando em $-\hat{z}$
Eq. (4) obtida de (3) via Eq. de Maxwell

Estrutura Perturbada:

$$\epsilon^{(1)} = \epsilon(x, y) + \epsilon_0 \Delta n^2 f(z) \quad (5)$$

$$E_t^{(1)} = a(z) E_t^T + b(z) E_t^0 \quad (6)$$

$$H_t^{(1)} = a(z) H_t^T + b(z) H_t^0 \quad (7)$$

$$f(z) = \frac{z}{\pi} \cos\left(\frac{2\pi}{\Lambda} z\right)$$

$$\Delta n^2 = n_j^2 - n_{j-1}^2 \quad j = 1, 2, 3, \dots, N$$

Os campos da Estrutura 1 propagam em $+\hat{z}$

Para obtermos as componentes longitudinais, utilizamos as Eq. de Maxwell.

$$\begin{aligned} \nabla \times E &= -i\omega \mu H \\ \nabla \times H &= -i\omega \epsilon E \end{aligned} \quad \left. \vphantom{\begin{aligned} \nabla \times E &= -i\omega \mu H \\ \nabla \times H &= -i\omega \epsilon E \end{aligned}} \right\} \text{simétricas}$$

$$E_z = -\frac{1}{i\omega \epsilon} \nabla_z \times H_t \quad \checkmark$$

$$H_z = \frac{1}{i\omega \mu} \nabla_z \times E_t \quad \checkmark$$

substituindo

$$E_z^{(1)} = -\frac{1}{i\omega \epsilon^{(1)}} \left[a(z) \nabla_z \times H_t^T + b(z) \nabla_z \times H_t^B \right]$$

$$H_z^{(1)} = \frac{1}{i\omega \mu} \left[a(z) \nabla_z \times E_t^T + b(z) \nabla_z \times E_t^B \right]$$

$$E_z^{(1)} = \frac{\epsilon^{(2)}}{\epsilon^{(1)}} \left[a(z) E_z^T + b(z) E_z^B \right]$$

$$H_z^{(1)} = a(z) H_z^T + b(z) H_z^B$$

porque a estrutura n perturba

8 e 9

8 e 9

$\epsilon^{(2)}$ E_z^B

10

11

Os campos perturbados totais tornam-se:

$$E^{(1)} = a(z) \left(E_1^T + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^T \right) + b(z) \left(E_1^0 + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^0 \right) \quad (12)$$

$$H^{(1)} = a(z) \left(H_1^T + \hat{z} H_2^T \right) + b(z) \left(H_1^0 + \hat{z} H_2^0 \right) \quad (13)$$

Considere o lado esquerdo de (9):

$$LE = \frac{d}{dz} \int \left[\underbrace{E^{(1)} \times H^{(2)}}_{1^\circ \text{ termo}} - \underbrace{E^{(2)} \times H^{(1)}}_{2^\circ \text{ termo}} \right] \cdot \hat{z} dz dy \quad (14)$$

Substituindo (3), (4), (12) e (13) \rightarrow (14)

$$1^\circ \text{ termo} = E^{(1)} \times H^{(2)}$$

$$= \left\{ a(z) \left(E_1^T + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^T \right) + b(z) \left(E_1^0 + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^0 \right) \right\} \times$$

$$\left(-H_1^T + \hat{z} H_2^T \right) e^{-i\beta_T z}$$

$$a(z) \left[-E_1^T \times H_2^T + \cancel{E_1^T \times \hat{z} H_2^T} + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^T \times (-H_1^T) + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^T \times \hat{z} H_2^T \right] e^{-i\beta_T z} + \hat{z} \times \hat{z} = 0$$

* por causa do \hat{z} em (14)

$$b(z) \left[-E_1^0 \times H_1^T + \cancel{E_1^0 \times \hat{z} H_2^T} + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^0 \times (-H_1^T) + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_2^0 \times \hat{z} H_2^T \right] e^{-i\beta_T z}$$

$$E^{(1)} \times H^{(2)} = -a(z) e^{-j\beta_T z} E_T^T \times H_T^T - b(z) e^{-j\beta_T z} E_T^B \times H_T^T$$

$$z^{\text{º}} \text{ termo} = E^{(2)} \times H^{(1)}$$

$$= \left(E_T^T - \hat{z} E_z^T \right) e^{-j\beta_T z} \times \left\{ a(z) \left(H_T^T + \hat{z} H_z^T \right) + b(z) \left(H_T^B + \hat{z} H_z^B \right) \right\}$$

$$= a(z) \left[E_T^T \times H_T^T + E_T^T \times \hat{z} H_z^T - \hat{z} E_z^T \times H_T^T - \hat{z} E_z^T \times \hat{z} H_z^T \right] e^{-j\beta_T z} +$$

$$b(z) \left[E_T^T \times H_T^B + E_T^T \times \hat{z} H_z^B - \hat{z} E_z^T \times H_T^B - \hat{z} E_z^T \times \hat{z} H_z^B \right] e^{-j\beta_T z}$$

$$E^{(2)} \times H^{(1)} = a(z) e^{-j\beta_T z} E_T^T \times H_T^T + b(z) e^{-j\beta_T z} E_T^T \times H_T^B$$

Logo,

$$LE = \frac{d}{dz} \int \left[-a(z) e^{-j\beta_T z} E_T^T \times H_T^T - b(z) e^{-j\beta_T z} E_T^B \times H_T^T - a(z) e^{-j\beta_T z} E_T^T \times H_T^T - b(z) e^{-j\beta_T z} E_T^T \times H_T^B \right] \cdot \hat{z} \, d\alpha \, dy$$

$$LE = -z \frac{d}{dz} \left(a(z) e^{-j\beta_T z} \right) \int E_T^T \times H_T^T \cdot \hat{z} \, d\alpha \, dy -$$

$$\frac{d}{dz} \left(b(z) e^{-j\beta_T z} \right) \int \left[E_T^B \times H_T^T + E_T^T \times H_T^B \right] \cdot \hat{z} \, d\alpha \, dy \quad (15)$$

Definindo: Integral de Overlapping

$$C_{mn} = \frac{1}{2} \int E_T^n \times H_T^m \cdot \hat{z} \, d\alpha \, dy \quad (16)$$

$$m, n = T, B$$

Como estamos lidando com modos da estrutura inteira :

$$C_{TT} = C_{OB} = L \quad C_{TB} = C_{OT} = 0$$

Portanto :

$$LE = -4 \frac{d}{dz} \left(a(z) e^{-i\beta_T z} \right)$$

$$LE = -4 \left[\frac{da(z)}{dz} e^{-i\beta_T z} - i\beta_T a(z) e^{-i\beta_T z} \right] \quad (17)$$

Para concluir, divida (17) por 4 :

$$LE = \left[-\frac{da(z)}{dz} + i\beta_T a(z) \right] e^{-i\beta_T z} \quad (18)$$

Considere, agora, o lado direito de (1) :

$$LD = i\omega \int [E^{(2)} - E^{(1)}] E^{(1)} \cdot E^{(2)} dx dy \quad (19)$$

Substituindo (2), (3), (5) e (12) \rightarrow (19)

$$LD = i\omega \int \left[\epsilon^{(2)} - \epsilon^{(1)} \right] \left\{ a(z) \left(E_t^T + \hat{z} \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^T \right) + b(z) \left(E_t^0 + \hat{z} \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^0 \right) \right\} \cdot \left(E_t^T - \hat{z} E_z^T \right) e^{-i\beta_T z} dx dy$$

$$LD = i\omega \int \Delta \epsilon \left\{ a(z) \left(E_t^T \cdot E_t^T - \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^T \cdot E_z^T \right) + b(z) \left(E_t^0 \cdot E_t^0 - \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^0 \cdot E_z^0 \right) \right\} e^{-i\beta_T z} dx dy \quad (20)$$

onde: $\Delta \epsilon = \epsilon^{(2)} - \epsilon^{(1)} = \epsilon(x, y) - \epsilon_0 \Delta n^2 f(z) = -\epsilon_0 \Delta n^2 f(z)$

Como feito p/ (17), divide (20) por 4:

$$LD = \frac{i\omega}{4} \int \Delta \epsilon \left[a(z) \left(E_t^T \cdot E_t^T - \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^T \cdot E_z^T \right) + b(z) \left(E_t^0 \cdot E_t^0 - \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^0 \cdot E_z^0 \right) \right] e^{-i\beta_T z} dx dy \quad (21)$$

Podemos reduzir (21) definindo uma nova quantidade:

$$K_{mn} = \frac{\omega \epsilon_0 \Delta n^2}{4} \int \left[E_t^m \cdot E_t^n - \frac{\epsilon^{(2)}}{\epsilon^{(1)}} E_z^m \cdot E_z^n \right] dx dy \quad (22)$$

$$m, n = T, 0$$

↖
Coeficiente de Acoplamento

$$LD = - \left[i a(z) K_{TT} f(z) + i b(z) K_{\theta T} f(z) \right] e^{-i\beta_T z} \quad (23)$$

$$LE = LD$$

$$\left[-\frac{da(z)}{dz} + i\beta_T a(z) \right] e^{-i\beta_T z} = - \left[i a(z) K_{TT} f(z) + i b(z) K_{\theta T} f(z) \right] e^{-i\beta_T z}$$

$$\frac{da(z)}{dz} = i \left[\beta_T + K_{TT} f(z) \right] a(z) + i K_{\theta T} f(z) b(z) \quad (24)$$

Passo 2

Deixe q os campos $E^{(2)}$ e $H^{(2)}$ e a permissividade $\epsilon^{(2)}$

representem, novamente, a estrutura multimodo não-perturbada; mas suponha q os campos não-perturbados sejam representados pelos modos do guia inferior:

$$e^{(2)} = \epsilon(x, y) \quad (25)$$

$$E^{(2)} = \left(E_t^0 - \hat{z} E_z^0 \right) e^{-i\beta_0 z} \quad (26)$$

$$H^{(2)} = \left(-H_t^0 + \hat{z} H_z^0 \right) e^{-i\beta_0 z} \quad (27)$$

Os modos da estrutura perturbada e como no passo anterior:

$$E^{(1)} = E(x, y) + \epsilon_0 \Delta n^2 f(z) \quad (28)$$

$$E^{(1)} = a(z) \left(E_t^T + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^T \right) + b(z) \left(E_t^B + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^B \right) \quad (29)$$

$$H^{(1)} = a(z) \left(H_t^T + \hat{z} H_z^T \right) + b(z) \left(H_t^B + \hat{z} H_z^B \right) \quad (30)$$

Substituindo (26), (27), (29) e (30) \rightarrow (14) :

$$\begin{aligned} E^{(1)} \times H^{(2)} &= \left\{ a(z) \left(E_t^T + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^T \right) + b(z) \left(E_t^B + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^B \right) \right\} \times \left(-H_t^B + \hat{z} H_z^B \right) e^{-j\beta_0 z} \\ &= a(z) e^{-j\beta_0 z} \left[-E_t^T \times H_t^B + E_t^T \times \hat{z} H_z^B + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^T \times (-H_t^B) + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^T \times \hat{z} H_z^B \right] + \\ &\quad b(z) e^{-j\beta_0 z} \left[-E_t^B \times H_t^B + E_t^B \times \hat{z} H_z^B + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^B \times (-H_t^B) + \hat{z} \frac{E^{(2)}}{E^{(1)}} E_z^B \times \hat{z} H_z^B \right] \end{aligned}$$

$$E^{(1)} \times H^{(2)} = -a(z) e^{-j\beta_0 z} E_t^T \times H_t^B - b(z) e^{-j\beta_0 z} E_t^B \times H_t^B$$

$$\begin{aligned} E^{(2)} \times H^{(1)} &= \left(E_t^B - \hat{z} E_z^B \right) e^{-j\beta_0 z} \times \left\{ a(z) \left(H_t^T + \hat{z} H_z^T \right) + b(z) \left(H_t^B + \hat{z} H_z^B \right) \right\} \\ &= a(z) e^{-j\beta_0 z} \left[E_t^B \times H_t^T + E_t^B \times \hat{z} H_z^T - \hat{z} E_z^B \times H_t^T - \hat{z} E_z^B \times \hat{z} H_z^T \right] + \\ &\quad b(z) e^{-j\beta_0 z} \left[E_t^B \times H_t^B + E_t^B \times \hat{z} H_z^B - \hat{z} E_z^B \times H_t^B - \hat{z} E_z^B \times \hat{z} H_z^B \right] \end{aligned}$$

$$E^{(2)} \times H^{(1)} = a(z) e^{-j\beta_0 z} E_t^B \times H_t^T + b(z) e^{-j\beta_0 z} E_t^B \times H_t^B$$

$$LE = \frac{d}{dz} \left\{ \left[-a(z) e^{-j\beta_0 z} E_t^T \times H_t^B - b(z) e^{-j\beta_0 z} E_t^B \times H_t^B \right] e^{-j\beta_0 z} - \left[a(z) e^{-j\beta_0 z} E_t^B \times H_t^T + b(z) e^{-j\beta_0 z} E_t^B \times H_t^B \right] \right\} \cdot \hat{z} \, dx \, dy$$

$$LE = \frac{d}{dz} \int \left\{ -a(z) \left[E_t^T \times H_t^B + E_t^B \times H_t^T \right] - b(z) \left[E_t^B \times H_t^B + E_t^B \times H_t^B \right] \right\} e^{-j\beta_0 z} \cdot \hat{z} \, dx \, dy$$

$$LE = -\frac{d}{dz} \int \left\{ a(z) \left[E_t^T \times H_t^B + E_t^B \times H_t^T \right] + 2b(z) E_t^B \times H_t^B \right\} e^{-j\beta_0 z} \cdot \hat{z} \, dx \, dy$$

Utilizando (16):

$$LE = -4 \frac{d}{dz} \left(b(z) e^{-j\beta_0 z} \right) \quad (31)$$

Dividindo (31) por 4 e fazendo a derivada:

$$LE = \left[-\frac{db(z)}{dz} + j\beta_0 b(z) \right] e^{-j\beta_0 z} \quad (32)$$

Considere, agora, o lado direito de (1) como dado

em (19):

$$LD = i\omega \int \left[e^{(2)} - e^{(1)} \right] \left(E_t^B - \hat{z} E_z^B \right) e^{-j\beta_0 z} \cdot \left\{ a(z) \left(E_t^T + \hat{z} \frac{E_z^{(2)}}{E_z^{(1)}} E_z^T \right) + b(z) \left(E_t^B + \hat{z} \frac{E_z^{(2)}}{E_z^{(1)}} E_z^B \right) \right\} \, dx \, dy$$

$$E^T \cdot E^{(1)} \cdot E^{(2)} \quad \text{ou} \quad E^T \cdot E^{(2)} \cdot E^{(1)}$$

(19)

$$LD = i\omega \int \Delta E \left\{ a(z) \left(E_t^T \cdot E_t^0 - \frac{E^{(z)}}{E^{(1)}} E_z^T \cdot E_z^0 \right) + b(z) \left(E_t^0 \cdot E_t^0 - \frac{E^{(z)}}{E^{(1)}} E_z^0 \cdot E_z^0 \right) \right\} e^{-i\beta_0 z} dz dy \quad (33)$$

$$\Delta E = E^{(z)} - E^{(1)} = E(x, y) - E(x, y) - \epsilon_0 \Delta n^2 f(z) = -\epsilon_0 \Delta n^2 f(z)$$

Dividindo (33) por 4 e utilizando (32):

$$LD = -i \left[a(z) K_{TO} f(z) + b(z) K_{BO} f(z) \right] e^{-i\beta_0 z} \quad (34)$$

$$LD = LE$$

$$\left[-\frac{db(z)}{dz} + i\beta_0 b(z) \right] e^{-i\beta_0 z} = -i \left[a(z) K_{TO} f(z) + b(z) K_{BO} f(z) \right] e^{-i\beta_0 z}$$

$$\frac{db(z)}{dz} = i K_{TO} f(z) a(z) + i \left[\beta_0 + K_{BO} f(z) \right] b(z) \quad (35)$$

As amplitudes complexas $a(z)$ e $b(z)$ são funções de z com variações rápidas. Estas variações podem ser removidas introduzindo amplitudes $c/$ variações lentas $A(z)$ e $B(z)$

$$a(z) = A(z) e^{i\beta_T z} \quad (36)$$

$$b(z) = B(z) e^{i\beta_B z} \quad (37)$$

Substituindo (36) e (37) em (29) e (35)

$$\left\{ \begin{array}{l} \frac{dA(z)}{dz} = i f(z) K_{TT} A(z) + i f(z) K_{BT} B(z) e^{i(\beta_0 - \beta_T)z} \end{array} \right. \quad (38)$$

$$\left\{ \begin{array}{l} \frac{dB(z)}{dz} = i f(z) K_{TB} A(z) e^{i(\beta_T - \beta_0)z} + i f(z) K_{BB} B(z) \end{array} \right. \quad (39)$$