

## HINGE MOMENT COEFFICIENT DERIVATIVES FOR TRAILING-EDGE CONTROLS ON WINGS AT SUBSONIC SPEEDS

### 1. NOTATION AND UNITS (see Sketch 1.1)

		<i>SI</i>	<i>British</i>
$A$	aspect ratio		
$(a_1)_0$	lift-curve slope with angle of attack for two-dimensional section in incompressible flow	rad <sup>-1</sup>	rad <sup>-1</sup>
$(a_2)_0$	lift-curve slope with control deflection for two-dimensional section in incompressible flow	rad <sup>-1</sup>	rad <sup>-1</sup>
$b_1$	hinge moment coefficient derivative $\partial C_H / \partial \alpha$	rad <sup>-1</sup>	rad <sup>-1</sup>
$(b_1)_0$	value of $b_1$ for two-dimensional section in incompressible flow	rad <sup>-1</sup>	rad <sup>-1</sup>
$\Delta b_1$	contribution from induced camber to $b_1$	rad <sup>-1</sup>	rad <sup>-1</sup>
$b_2$	hinge moment coefficient derivative $\partial C_H / \partial \delta$	rad <sup>-1</sup>	rad <sup>-1</sup>
$b'_2$	hinge moment coefficient derivative $\partial C_H / \partial \delta'$	rad <sup>-1</sup>	rad <sup>-1</sup>
$(b_2)_0$	value of $b_2$ for two-dimensional section in incompressible flow	rad <sup>-1</sup>	rad <sup>-1</sup>
$\Delta b_2$	contribution from induced camber to $b_2$	rad <sup>-1</sup>	rad <sup>-1</sup>
$C_H$	hinge moment coefficient for control, $H / \frac{1}{2} \rho V^2 \bar{c}_f s_f$		
$C_L$	lift coefficient for wing, $L / \frac{1}{2} \rho V^2 S$		
$c$	wing chord	m	ft
$c'$	wing chord measured normal to wing quarter-chord line	m	ft
$c'_b$	control balance chord forward of hinge line measured normal to wing quarter-chord line	m	ft
$c_f$	control chord aft of hinge line	m	ft
$c'_f$	control chord aft of hinge line measured normal to wing quarter-chord line	m	ft
$\bar{c}_f$	geometric mean control chord aft of hinge line, $\int_{\eta_i}^{\eta_o} c_f d\eta$	m	ft

$\bar{c}_f$	aerodynamic mean control chord aft of hinge line, $\int_{\eta_i}^{\eta_o} c_f^2 d\eta / \bar{c}_f$	m	ft
$F_B$	factor on induced camber contributions to allow for control balance		
$G_1, G_2, G_3$	functions used in calculation of induced camber contribution, see Equations (3.5) and (3.6)	rad <sup>-1</sup>	rad <sup>-1</sup>
$H$	hinge moment	N m	lbf ft
$K$	correction factor for rectangular wing, see Section 3.4		
$L$	lift	N	lbf
$M$	Mach number		
$R$	Reynolds number based on wing mean chord		
$S$	wing area	m <sup>2</sup>	ft <sup>2</sup>
$s$	wing semispan	m	ft
$s_f$	control span, $s_f = s(\eta_o - \eta_i)$	m	ft
$t$	maximum thickness of wing section	m	ft
$t_h$	thickness of control at hinge line	m	ft
$V$	free-stream velocity	m/s	ft/s
$\alpha$	angle of attack	rad	rad
$\alpha_i$	mean induced angle of attack	rad	rad
$\beta$	compressibility parameter, $(1 - M^2)^{1/2}$		
$\delta$	control deflection angle measured in plane parallel to plane of symmetry	rad	rad
$\delta'$	control deflection angle measured in plane normal to hinge line	rad	rad
$\Lambda_{1/4}$	sweep angle of wing quarter-chord line	deg	deg
$\Lambda_{1/2}$	sweep angle of wing half-chord line	deg	deg
$\Lambda_h$	sweep angle of control hinge-line	deg	deg
$\eta$	spanwise distance from wing centre-line as fraction of semispan		

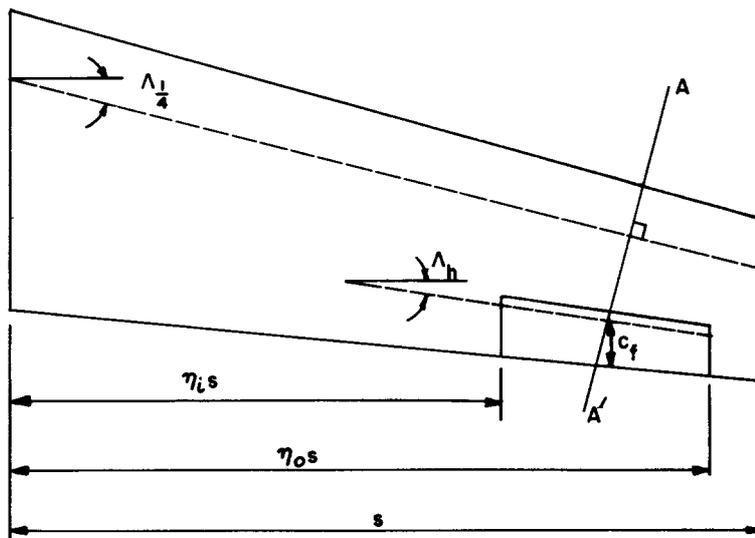
$\eta_i, \eta_o$	values of $\eta$ at inboard and outboard limits of aileron at hinge line		
$\lambda$	ratio of wing tip chord to wing centre-line chord		
$\rho$	density of air	kg/m <sup>3</sup>	slug/ft <sup>3</sup>
$\tau'$	section trailing-edge angle in plane normal to quarter-chord line	deg	deg

**Superscript**

\* as in  $(a_1)_0^*$  denotes properties for a 'standard' section for which  $\tau' = 2 \tan^{-1}(t/c')$

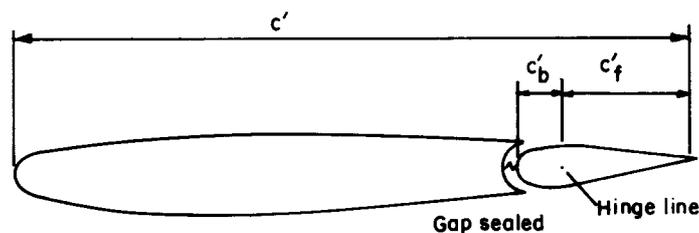
**Subscript**

$T$  as in  $(a_1)_{0T}$  denotes theoretical value



Planform geometry

A A' normal to quarter-chord line and passing through mid-span of control



Section A A'

Sketch 1.1 Planform and section geometry

## 2. INTRODUCTION

This Item provides a method for estimating the hinge moment coefficient derivatives  $b_1$  and  $b_2$  for full-span and part-span sealed controls on swept and unswept wings at speeds where the flow over the wing is wholly subsonic and fully attached. It may also be used for controls on tailplanes or fins. The method follows the traditional approach for a wing of finite aspect ratio, whereby the two-dimensional sectional values are corrected for the effects of induced angle of attack and induced camber, see Derivations 1, 2, 11, 12 and 18 for example. The basic equations for the derivatives are essentially in the form of lifting-line theory, but with lifting-surface theories used to give a better evaluation of the induced angle of attack and to provide the important additional contribution due to induced camber. The equations permit ready substitution of known two-dimensional characteristics.

None of the various published versions of the above technique is sufficiently comprehensive in its theoretical calculations to cover in a satisfactory manner all configurations of interest. The method in this Item follows Derivation 18 in its treatment of wing sweep, which was based on a successful modification to the earlier work of Derivation 11. The presence of control balance is allowed for by adapting the theoretical calculations made in Derivation 2 for unswept wings. Empirical corrections are presented for part-span effects for controls that extend from a general inboard station  $\eta_i$  to near the wing tip,  $\eta_o \geq 0.9$ . The two-dimensional properties of the wing and control are assumed to be effectively constant over the span of the control. Prandtl-Glauert similarity parameters are employed to model the first order effects of compressibility at low to moderate subsonic speeds.

Item No. Aero C.04.01.00 (Derivation 29) provides a general introduction to the treatment of control hinge moment coefficient derivatives within the Aerodynamics Sub-series. It should be consulted for a brief description of the individual Items that are available and their inter-relationship. In addition, it gives qualitative guidance on the effects of control geometry and flow conditions on the range of linearity of hinge moment characteristics.

## 3. METHOD

### 3.1 Basic Equations

In incompressible flow the derivatives for a finite wing are related to the section values normal to the wing quarter-chord line by the equations

$$b_1 = (b_1)_0 \left(1 - \frac{\alpha_i}{\alpha}\right) \cos \Lambda_{1/4} \cos \Lambda_h + \Delta b_1 \quad (3.1)$$

$$\text{and} \quad b_2 = \left( (b_2)_0 - \frac{\alpha_i}{\delta} (b_1)_0 \right) \cos \Lambda_{1/4} \cos \Lambda_h + \Delta b_2, \quad (3.2)$$

where  $\alpha_i$  is the mean induced angle of attack and  $\Delta b_1$  and  $\Delta b_2$  are induced camber contributions. These equations make the assumption that the angle  $\alpha_i$  is constant across the span, which is a reasonable assumption for wings that are approximately elliptically loaded. For loadings that depart significantly from the elliptical a spanwise integration with  $\alpha_i$  as a function of  $\eta$  is necessary, as described in Derivation 2 or Item No. Aero C.04.01.06 (Derivation 28). However, a relatively simple correction is possible for a rectangular wing that has constant two-dimensional properties, see Section 3.4.

### 3.2 Derivative $b_1$

The assumptions made in Derivation 18 lead to the substitution

$$\frac{\alpha_i}{\alpha} = 1 - \frac{dC_L/d\alpha}{(a_1)_0 \cos \Lambda_{1/4}}, \quad (3.3)$$

where  $(a_1)_0$  is the two-dimensional lift-curve slope for the section normal to the wing quarter-chord line and  $dC_L/d\alpha$  is the wing lift-curve slope estimated by lifting-surface theory. Values of  $dC_L/d\alpha$  can be obtained from Item No. 70011 (Derivation 34). As in Derivation 20, for example, Prandtl-Glauert similarity parameters are incorporated to allow for first-order compressibility effects. With  $dC_L/d\alpha$  calculated at the Mach number of interest, the general form for  $b_1$  becomes

$$b_1 = \frac{(b_1)_0}{(a_1)_0} \left( \frac{dC_L}{d\alpha} \right) \cos \Lambda_h + \Delta b_1. \quad (3.4)$$

In this Item the correction  $\Delta b_1$  is evaluated as

$$\Delta b_1 = G_1 + G_2, \quad (3.5)$$

where  $G_1$  is the theoretically based contribution for a full-span control and  $G_2$  is a purely empirical correction for part-span effects.

Lifting-surface values of  $G_1$  for a plain control in incompressible flow were calculated in Derivation 18 for  $2 \leq A \leq 6$  and  $0 \leq \Lambda_{1/4} \leq 45^\circ$ . These are reproduced in Figure 1 where  $2\pi\beta G_1/F_B(a_1)_0 \cos \Lambda_h$  is given in three carpets as a function of  $(1/\beta)\tan \Lambda_{1/4}$  and  $\beta A$  for values of the control chord ratio measured normal to the quarter-chord line,  $c_f'/c' = 0.2, 0.3$  and  $0.4$ . The data for unswept wings have been slightly adjusted to be compatible with the data for unswept wings given in Derivation 2 as these covered higher aspect ratios. For swept wings, extrapolations above  $\beta A = 6$  have been made by taking the form of the variation for the unswept wing as a guide and noting that, for all sweep angles,  $G_1$  must tend to zero as  $\beta A$  becomes very large. The factor  $F_B$  allows for control balance and has been deduced from the data given in Derivation 2 for nose-balanced and internally-balanced controls on unswept wings. For those two types of control  $F_B$  is given in Figure 2 as a function of  $c_f'/c'$  and balance  $c_b'/c_f'$ . For a plain control  $F_B$  is unity.

The empirical part-span correction  $G_2$  is obtained from Figure 3 where  $2\pi\beta G_2/F_B(a_1)_0 \cos \Lambda_h$  is given in a carpet as a function of  $\eta_i$  and  $A \tan \Lambda_{1/2}$ . It is assumed that  $\eta_o \geq 0.9$ .

### 3.3 Derivative $b_2$

In the evaluation of  $b_2$  given in Derivation 18 for full-span controls, some necessary simplifications were made to facilitate the lifting-surface calculations of the control-deflected loadings and only a partial evaluation of the induced effects was performed. Comparisons with experimental data showed that the overall contribution of the neglected components was acceptably small. However, although the resulting equation for  $b_2$  resembles Equation (3.2) the approximation is such that the separate effects of induced angle of attack and induced camber cannot be identified. The equation for  $b_2$  is

$$b_2 = \left( (b_2)_0 - \frac{(a_2)_0}{(a_1)_0} (b_1)_0 \right) \frac{\cos \Lambda_h}{(\beta^2 + \tan^2 \Lambda_{1/4})^{1/2}} + \frac{(a_2)_0}{(a_1)_0} (b_1 + G_3), \quad (3.6)$$

where the function  $G_3$  has been introduced in this Item to provide a further empirical allowance for part-span effects on the induced camber contribution, in addition to that made in  $b_1$ . Similarity parameters are again introduced to allow for the effects of compressibility. The lift-curve slope with control deflection  $(a_2)_0$  is for a section normal to the wing quarter-chord line and a control chord ratio  $c_f'/c'$ . Figure 4 gives  $2\pi\beta G_3/F_B(a_1)_0 \cos\Lambda_h$  as a function of  $\eta_i$ . It is assumed that  $\eta_o \geq 0.9$ .

For control deflection angles measured normal to the hinge line, the hinge moment coefficient derivative is

$$\frac{\partial C_H}{\partial \delta'} = b_2' = b_2 \cos\Lambda_h. \tag{3.7}$$

### 3.4 Rectangular Wing

For a rectangular wing with uniform section and control geometry, Item No. Aero C.04.01.06 contains information on the spanwise variation of the ratio of  $\alpha_i$  for a rectangular wing to the value of  $\alpha_i$  for an elliptically-loaded wing. This information has been used to deduce a simple factor  $K$  to modify Equations (3.4) and (3.6). The factor is simply the mean value of the  $\alpha_i$  ratio over the span of the control. Figure 5 shows  $K$  as a function  $\eta_i$  for  $\eta_o \geq 0.9$ . The equations for  $b_1$  and  $b_2$  become

$$b_1 = (b_1)_0 \left\{ 1 - K \left( 1 - \frac{dC_L/d\alpha}{(a_1)_0} \right) \right\} + \Delta b_1 \tag{3.8}$$

and 
$$b_2 = \frac{1}{\beta} \left( (b_2)_0 - K \frac{(a_2)_0}{(a_1)_0} (b_1)_0 \right) + \frac{(a_2)_0}{(a_1)_0} (b_1 + G_3). \tag{3.9}$$

With  $K$  taken from Figure 5, the equations are otherwise evaluated as in Sections 3.2 and 3.3. If  $K = 1$ , Equations (3.8) and (3.9) reduce to Equations (3.4) and (3.6) with  $\Lambda_{1/4} = \Lambda_h = 0$ . For a full-span control  $K \approx 1$ , but the correction becomes important for part-span controls.

### 3.5 Sectional Properties

The sectional properties that are required in the calculation of  $b_1$  and  $b_2$  can be obtained as indicated in the table, or experimental values may be substituted if they are known.

<i>Parameter</i>	<i>Item No. Aero</i>	<i>Derivation</i>	
$(a_1)_0$	W.01.01.05	30	
$(a_2)_0$	C.01.01.03	31	
$(b_1)_0$	C.04.01.01 (plain control)	32	
$(b_2)_0$	C.04.01.02 (plain control)	33	
Corrections to $(b_1)_0$	}	C.04.01.03 (nose balance)	26
and $(b_2)_0$ for balance		C.04.01.04 (Irving internal balance)	27

The use of these Derivations is illustrated in the Example (see Section 6).

## 4. ACCURACY AND APPLICABILITY

### 4.1 Accuracy

Sketches 4.1 and 4.2 demonstrate the overall accuracy of prediction at low speeds using test data from Derivations 2 and 4 to 25. In general,  $b_1$  is predicted to within about  $\pm 0.05 \text{ rad}^{-1}$  and  $b_2$  to within  $\pm 0.07 \text{ rad}^{-1}$ . As would be expected, estimated values are usually more reliable when known two-dimensional section properties are available. The method gives best estimates for full-span controls on unswept wings, but there is otherwise no general trend within the overall scatter. For comparison, the accuracy quoted in Item Nos Aero C.04.01.01 and 02 for  $(b_1)_0$  and  $(b_2)_0$  is  $\pm 0.05 \text{ rad}^{-1}$ .

The use of similarity parameters provides a reasonable representation of compressibility effects for Mach numbers at which the flow over the wing is wholly subsonic and fully attached, provided the section profiles are thin,  $t/c' \leq 0.14$  say, and the controls have straight-tapered profiles aft of the hinge line. The variations with Mach number are small in these cases. The method is unsatisfactory for predicting the larger variations associated with thicker sections and more complex control profiles. Examples of the different variations with Mach number that are displayed by various controls are contained in Derivation 9.

The empirical part-span correction for  $b_1$  differs considerably from the theoretical one given for unswept wings in Derivations 3 and 12. Indeed, as demonstrated by the form of  $G_2$ , best predictions are achieved with no correction for unswept wings, although one is needed for swept wings. The further empirical part-span correction introduced in the calculation of  $b_2$ , namely  $G_3$ , could only be defined crudely within a considerable scatter of experimental results but it does remove a bias that would otherwise exist in the prediction. For both part-span corrections, particular notice was taken of the results of parametric tests in which the control span was altered systematically.

### 4.2 Applicability

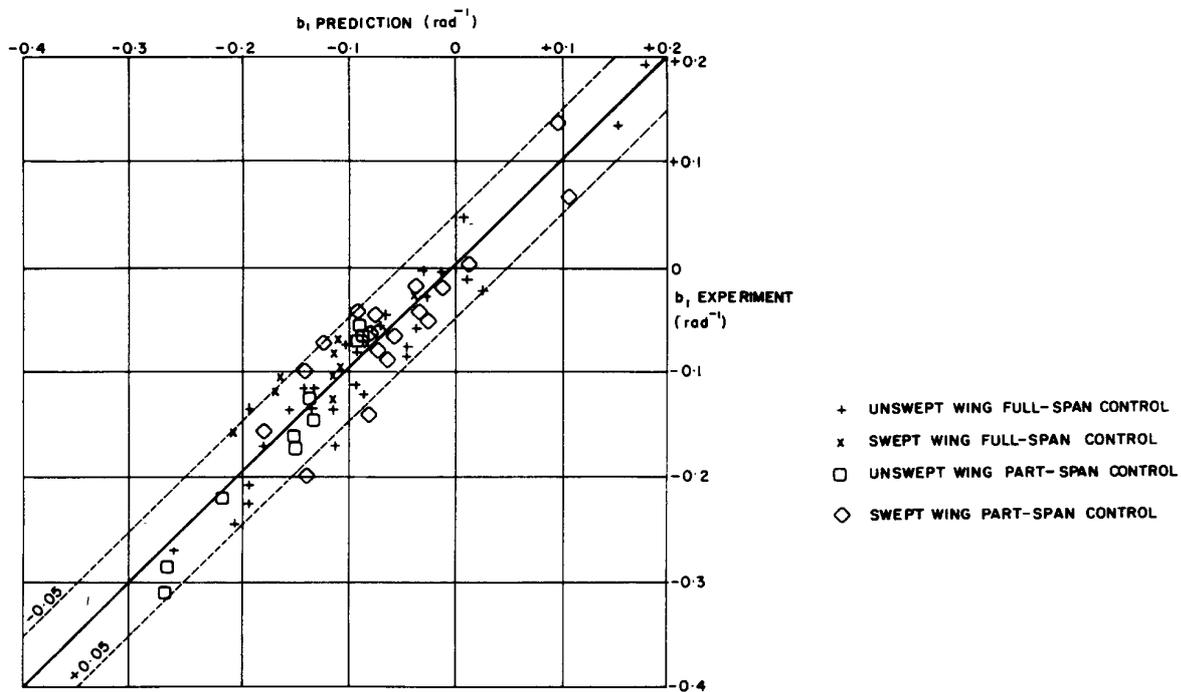
The derivatives  $b_1$  and  $b_2$  are defined over small ranges of angle of attack and control deflection, where the hinge moment coefficient varies linearly. A general discussion on factors affecting the extent of the linear range is given in the introductory Item No. Aero C.04.01.00 (Derivation 29).

An indication of the range of wing and control geometries covered in the development of the method is given in the table. The method applies to sealed controls with streamwise side-edges. If the side edges are normal to a swept hinge-line then there can be a significant change in  $b_2$ , although the effect on  $b_1$  is small (Derivation 10).

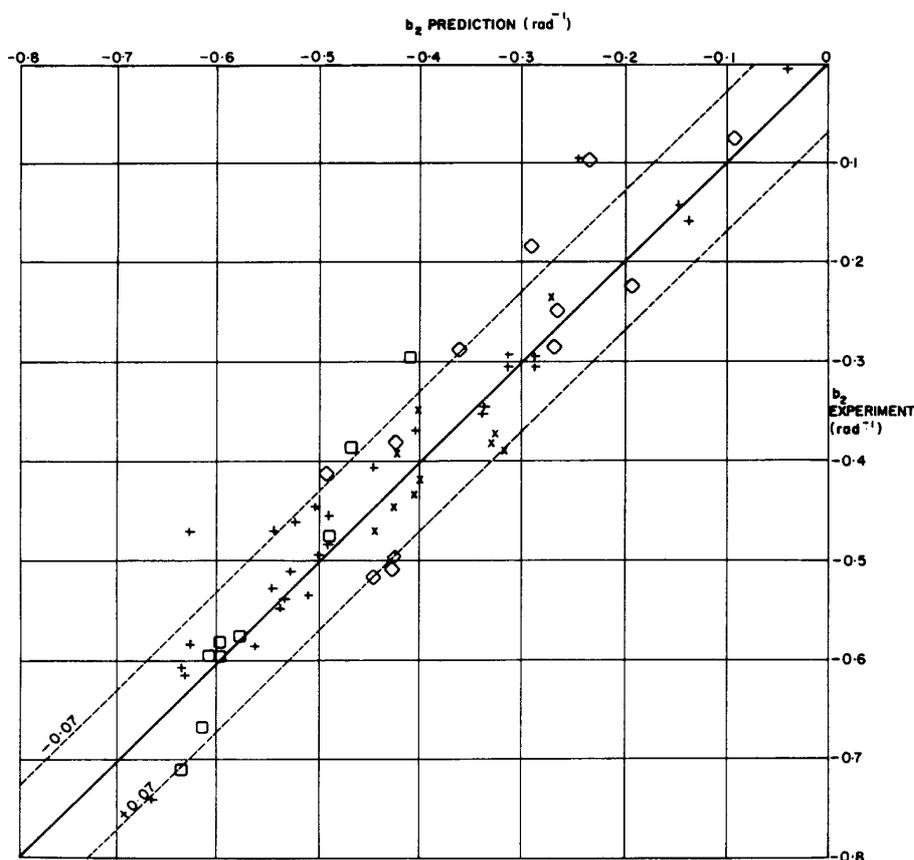
<i>Parameter</i>	<i>Range</i>	<i>Parameter</i>	<i>Range</i>
$A$	2 to 8	$t/c'$	0.06 to 0.14
$\Lambda_{1/2}$	0 to $50^\circ$	$\tau'$	$6^\circ$ to $20^\circ$
$c_f'/c'$	0.2 to 0.4	$\eta_i$	0 to 0.8

The main method in Sections 3.2 and 3.3 assumes that the wing loading is elliptical and that the wing and control two-dimensional properties are essentially constant over the span of the control. The modest departures from these conditions that are normally expected will not lead to undue errors in prediction. A simple modification is given in Section 3.4 for rectangular wings with constant two-dimensional properties. For the special case of full-span controls on unswept wings with a large spanwise variation in sectional properties the method of Item No. Aero C.04.01.06 (Derivation 28) may be used.

The method is based on an analysis of data on controls that extend to near the wing tip; therefore, it must be used with caution if applied to controls that are well inboard with  $\eta_o$  substantially less than 0.9.



Sketch 4.1 Comparison of predicted and experimental values of  $b_1$



Sketch 4.2 Comparison of predicted and experimental values of  $b_2$

## 5. DERIVATION

The Derivation lists selected sources that have assisted in the preparation of this Item.

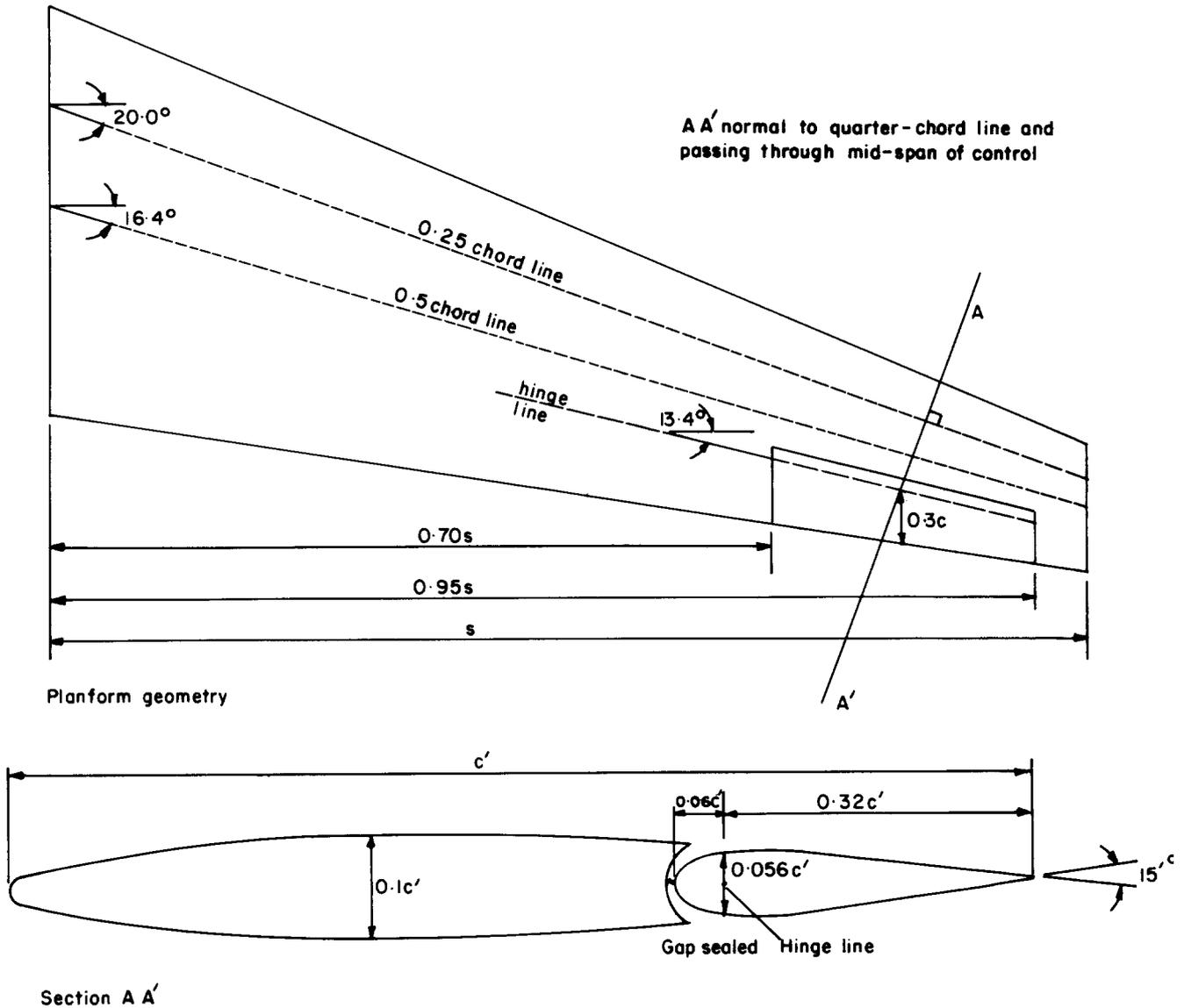
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26. ESDU Effect of nose balance on two-dimensional control hinge-moment coefficients. Item No. Aero C.04.01.03, ESDU International, 1949.
  27. ESDU Effect of Irving internal balance on hinge-moment coefficient in two-dimensional flow. Item No. Aero C.04.01.04, ESDU International, 1949.

28. ESDU Full-span control hinge moment coefficient derivatives in incompressible flow for unswept wings with allowance for spanwise variation of sectional properties. Item No. Aero C.04.01.06, ESDU International, 1949.
29. ESDU Introduction to Data Items on control hinge moments. Item No. C.04.01.00, ESDU International, 1950.
30. ESDU Slope of lift curve for two-dimensional flow. Item No. Aero W.01.01.05, ESDU International, 1955.
31. ESDU Rate of change of lift coefficient with control deflection in incompressible two-dimensional flow,  $(a_2)_0$ . Item No. Aero C.01.01.03, ESDU International, 1956.
32. ESDU Rate of change of hinge moment coefficient with incidence for a plain control in incompressible two-dimensional flow,  $(b_1)_0$ . Item No. Aero C.04.01.01, ESDU International, 1956.
33. ESDU Rate of change of hinge moment coefficient with control deflection for a plain control in incompressible two-dimensional flow,  $(b_2)_0$ . Item No. Aero C.04.01.02, ESDU International, 1956.
34. ESDU Lift-curve slope and aerodynamic centre position of wings in inviscid subsonic flow. Item No. 70011, ESDU International, 1970.

## 6. EXAMPLE

Calculate the hinge moment derivatives for the control shown in Sketch 6.1. Assume a Mach number of 0.4 and a Reynolds number of  $3.5 \times 10^7$  based on wing mean chord. The required geometric parameters are summarised in the table. The control has a nose balance with a rounded forward profile.



Sketch 6.1

From planform geometry		From section geometry			
$A$	$= 7.7$	$\lambda$	$= 0.3$	$\tau'$	$= 15.0^\circ$
$\Lambda_{1/4}$	$= 20.0^\circ$	$\eta_i$	$= 0.70$	$t/c'$	$= 0.1$
$\Lambda_{1/2}$	$= 16.4^\circ$	$\eta_o$	$= 0.95$	$c_b'/c'$	$= 0.06$
$\Lambda_h$	$= 13.4^\circ$	$c_f/c$	$= 0.30$	$c_f'/c'$	$= 0.32$

(i) Determine section properties  $(a_1)_0$ ,  $(a_2)_0$ ,  $(b_1)_0$ ,  $(b_2)_0$ .

If experimental data are available for section properties, go to Step (ii).

From Item No. Aero W.01.01.05, with a Reynolds number  $R = 3.5 \times 10^7$ , a trailing-edge angle  $\tau' = 15^\circ$ , a thickness chord ratio  $t/c' = 0.10$ , and an assumed boundary-layer transition point of  $0.3 c'$ ,

$$\frac{(a_1)_0}{(a_1)_{0T}} = 0.883$$

and  $(a_1)_{0T} = 6.788 \text{ rad}^{-1}$ ,

so  $(a_1)_0 = 5.994 \text{ rad}^{-1}$ .

From Item No. Aero C.01.01.03, with  $c_f'/c' = 0.32$ ,  $t/c' = 0.10$  and  $(a_1)_0/(a_1)_{0T} = 0.883$ ,

$$(a_2)_{0T} = 4.600 \text{ rad}^{-1},$$

and  $\frac{(a_2)_0}{(a_2)_{0T}} = 0.835$ ,

so  $(a_2)_0 = 3.841 \text{ rad}^{-1}$ .

The calculation of  $(b_1)_0$  and  $(b_2)_0$  requires corresponding values, denoted  $(a_1)_0^*$  etc, for a 'standard' aerofoil section with  $\tau' = 2 \tan^{-1}(t/c')$ . By again using Item Nos Aero W.01.01.05 and Aero C.01.01.03, with  $\tau' = 2 \tan^{-1}(0.10) = 11.4^\circ$ ,

$$\frac{(a_1)_0^*}{(a_1)_{0T}^*} = 0.906$$

and  $(a_1)_{0T}^* = 6.781 \text{ rad}^{-1}$ ,

so  $(a_1)_0^* = 6.143 \text{ rad}^{-1}$ .

Also  $(a_2)_{0T}^* = 4.600 \text{ rad}^{-1}$

and  $\frac{(a_2)_0^*}{(a_2)_{0T}^*} = 0.862$ ,

so  $(a_2)_0^* = 3.965 \text{ rad}^{-1}$ .

Then from Item No. Aero C.04.01.01 with  $t/c' = 0.10$ ,  $c_f'/c' = 0.32$  and  $(a_1)_0^*/(a_1)_{0T}^* = 0.906$ , for a plain unbalanced control in incompressible flow

$$-(b_1)_{0T}^* = 0.580 \text{ rad}^{-1}$$

and 
$$\frac{(b_1)_0^*}{(b_1)_{0T}^*} = 0.795,$$

so 
$$\begin{aligned} (b_1)_0 &= (b_1)_0^* + 2[(a_1)_{0T}^* - (a_1)_0^*](\tan \frac{1}{2}\tau' - t/c') \\ &= -0.580 \times 0.795 + 2[6.781 - 6.143](\tan 7.5^\circ - 0.10) \\ &= -0.421 \text{ rad}^{-1}. \end{aligned}$$

Similarly, from Item No. Aero C.04.01.02 with  $t/c' = 0.10$ ,  $c_f'/c' = 0.32$  and  $(a_2)_0^*/(a_2)_{0T}^* = 0.862$ ,

and 
$$\frac{(b_2)_0^*}{(b_2)_{0T}^*} = 0.880,$$

so 
$$\begin{aligned} (b_2)_0 &= (b_2)_0^* + 2[(a_2)_{0T}^* - (a_2)_0^*](\tan \frac{1}{2}\tau' - t/c') \\ &= -0.902 \times 0.880 + 2[4.600 - 3.965](\tan 7.5^\circ - 0.10) \\ &= -0.754 \text{ rad}^{-1}. \end{aligned}$$

The control has a nose balance so Item No. Aero C.04.01.03 must be used to correct the values of  $b_1$  and  $b_2$  that have been calculated above for plain controls. (For an Irving internal balance see Item No. Aero C.04.01.04.)

With  $t/c' = 0.10$ , a round nose profile, and a balance ratio given by

$$[(c_b'/c_f')^2 - (\frac{1}{2} t_h/c_f')^2]^{1/2} = [(0.06/0.32)^2 - (\frac{1}{2} \times 0.056/0.32)^2]^{1/2} = 0.166,$$

Item No. Aero C.04.01.03 gives the ratio of balanced to plain control derivatives

and 
$$\frac{(b_1)_{0Bal}}{(b_1)_{0Plain}} = 0.89$$

and 
$$\frac{(b_2)_{0Bal}}{(b_2)_{0Plain}} = 0.78.$$

Therefore, for the balanced control,

$$(b_1)_0 = -0.421 \times 0.89 = -0.375 \text{ rad}^{-1}$$

and 
$$(b_2)_0 = -0.754 \times 0.78 = -0.588 \text{ rad}^{-1}.$$

(ii) Calculate  $dC_L/d\alpha$  for the wing

For  $M = 0.4$ ,  $\beta = (1 - M^2)^{1/2} = 0.917.$

From Item No. 70011, with  $\lambda = 0.3$ ,  $A \tan \Lambda_{1/2} = 7.7 \tan 16.4^\circ = 2.27$ , and  $\beta A = 0.917 \times 7.7 = 7.06$ ,

$$\frac{1}{A} \frac{dC_L}{d\alpha} = 0.630 \text{ rad}^{-1}$$

so 
$$\frac{dC_L}{d\alpha} = 7.7 \times 0.630 = 4.851 \text{ rad}^{-1}.$$

(iii) Calculate  $b_1$

From Figure 1 with  $\beta A = 7.06$ ,  $c_f'/c' = 0.32$  and  $(1/\beta) \tan \Lambda_{1/4} = (1/0.917) \tan 20^\circ = 0.397$ , the full-span induced camber contribution parameter is

$$\frac{2\pi\beta G_1}{F_B(a_1)_0 \cos \Lambda_h} = -0.008 \text{ rad}^{-1}.$$

For a nose-balanced control with  $c_b'/c_f' = 0.06/0.32 = 0.188$  and  $c_f'/c' = 0.32$ , the balance factor  $F_B$  from Figure 2a is 0.935,

so 
$$G_1 = -0.008 \times \frac{0.935 \times 5.994 \times \cos 13.4^\circ}{2 \times \pi \times 0.917} = -0.008 \times 0.946 = -0.008 \text{ rad}^{-1}.$$

(Note that the bracketed numerical expression involved in evaluating  $G_1$ ,  $G_2$  and  $G_3$  is dimensionless because  $(a_1)_0$  is normalised by the presence of the theoretical thin aerofoil value  $2\pi \text{ rad}^{-1}$ .)

From Figure 3 with  $A \tan \Lambda_{1/2} = 2.27$  and  $\eta_i = 0.70$  the part-span induced camber contribution parameter is

$$\frac{2\pi\beta G_2}{F_B(a_1)_0 \cos \Lambda_h} = 0.070 \text{ rad}^{-1},$$

so 
$$G_2 = 0.070 \times \frac{0.935 \times 5.994 \times \cos 13.4^\circ}{2 \times \pi \times 0.917} = 0.070 \times 0.946 = 0.066 \text{ rad}^{-1}.$$

Therefore, from Equation (3.5)

$$\begin{aligned} \Delta b_1 &= G_1 + G_2 \\ &= -0.008 + 0.066 = 0.058 \text{ rad}^{-1}, \end{aligned}$$

and from Equation (3.4)

$$\begin{aligned} b_1 &= \frac{(b_1)_0}{(a_1)_0} \left( \frac{dC_L}{d\alpha} \right) \cos \Lambda_h + \Delta b_1 \\ &= \frac{-0.375}{5.994} \times 4.851 \times \cos 13.4^\circ + 0.058 = -0.295 + 0.058 = -0.237 \text{ rad}^{-1}. \end{aligned}$$

(iv) Calculate  $b_2$

From Figure 4 with  $\eta_i = 0.7$  the parameter giving the further part-span induced camber contribution is

$$\frac{2\pi\beta G_3}{F_B(a_1)_0 \cos\Lambda_h} = 0.105 \text{ rad}^{-1},$$

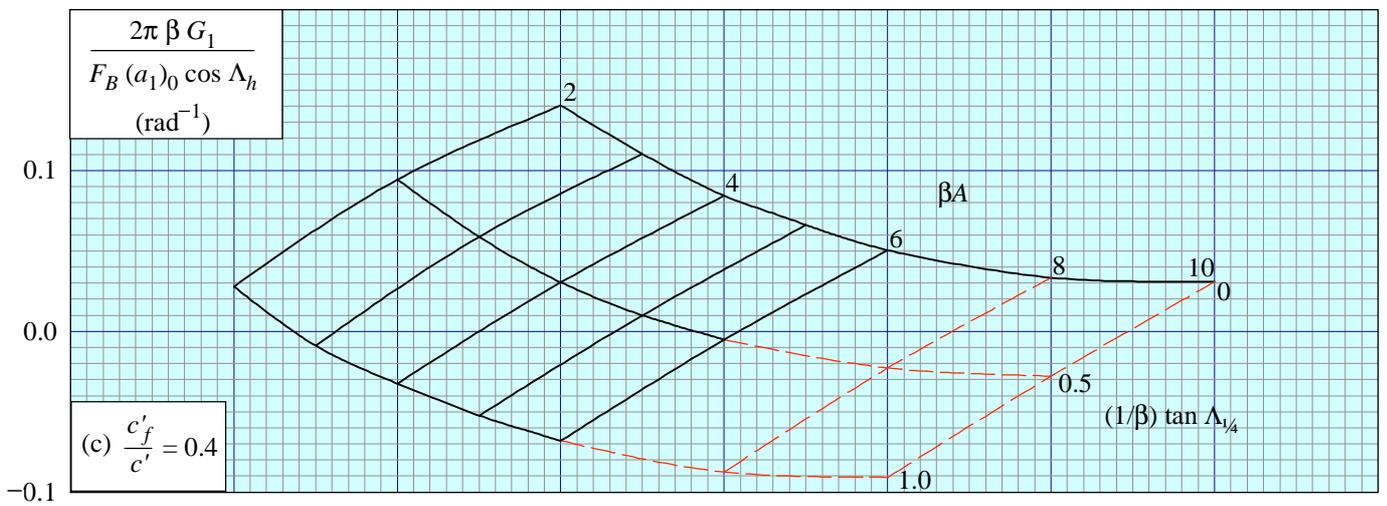
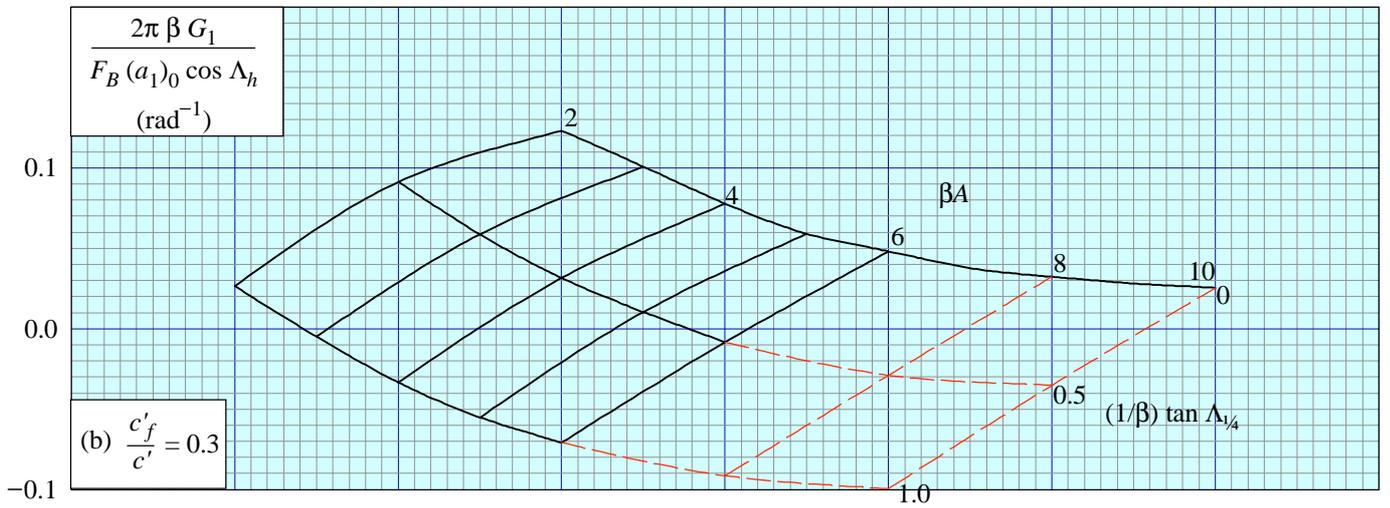
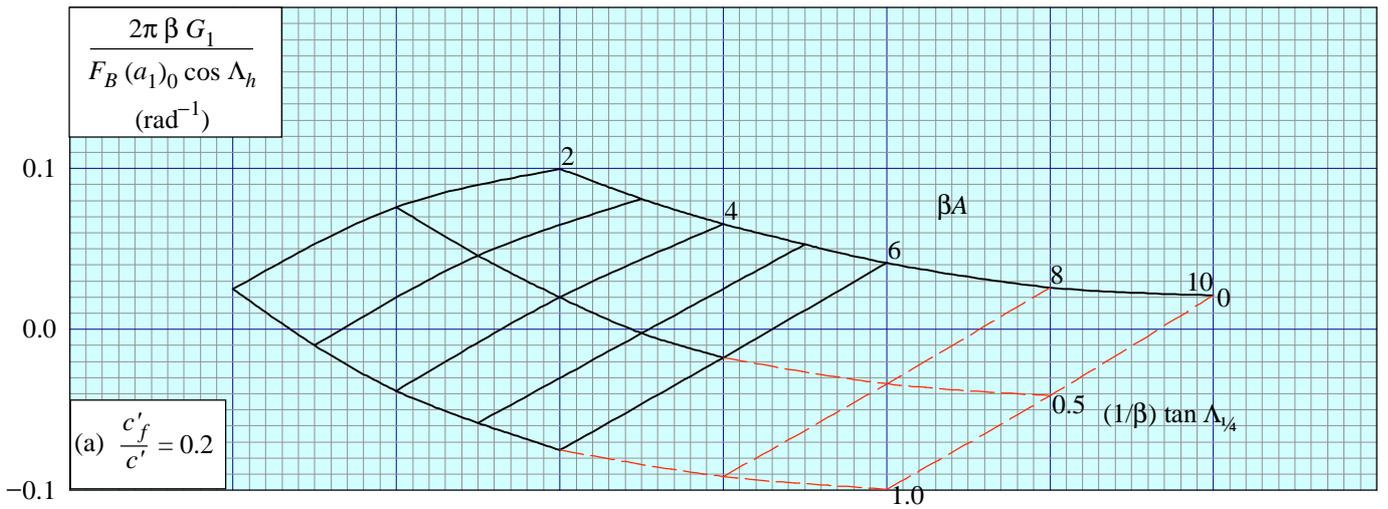
so 
$$G_3 = 0.105 \times \frac{0.935 \times 5.994 \times \cos 13.4^\circ}{2 \times \pi \times 0.917} = 0.105 \times 0.946 = 0.099 \text{ rad}^{-1}.$$

Therefore, from Equation (3.6)

$$\begin{aligned} b_2 &= \left( (b_2)_0 - \frac{(a_2)_0}{(a_1)_0} (b_1)_0 \right) \frac{\cos\Lambda_h}{(\beta^2 + \tan^2\Lambda_{1/4})^{1/2}} + \frac{(a_2)_0}{(a_1)_0} (b_1 + G_3) \\ &= \left( -0.588 - \frac{3.841}{5.994} (-0.375) \right) \frac{\cos 13.4^\circ}{(0.917^2 + \tan^2 20^\circ)^{1/2}} + \frac{3.841}{5.994} (-0.237 + 0.099) \\ &= -0.343 - 0.088 = -0.431 \text{ rad}^{-1}. \end{aligned}$$

For control angles measured normal to the hinge line,

$$\begin{aligned} b'_2 &= b_2 \cos\Lambda_h \\ &= -0.431 \cos 13.4^\circ \\ &= -0.419 \text{ rad}^{-1}. \end{aligned}$$



**FIGURE 1 INDUCED CAMBER CONTRIBUTION FOR FULL-SPAN CONTROLS**

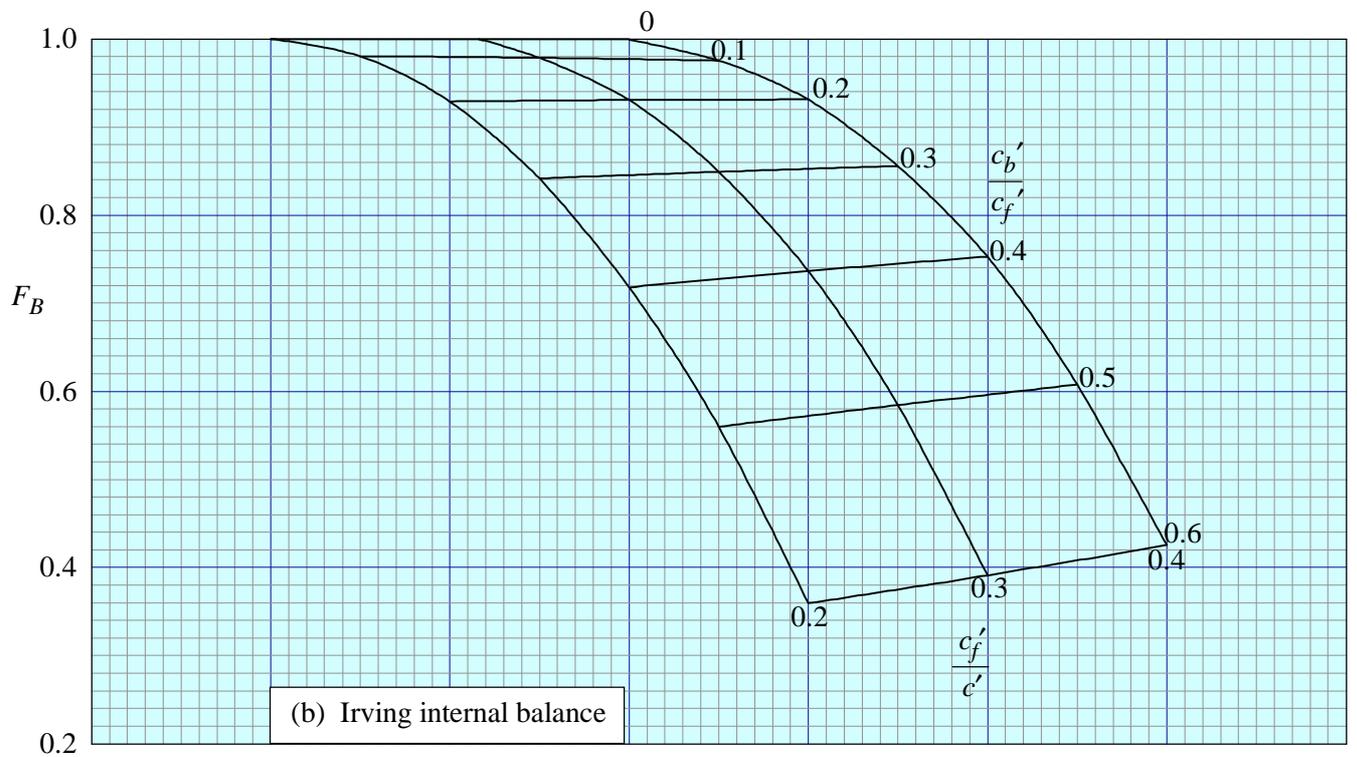
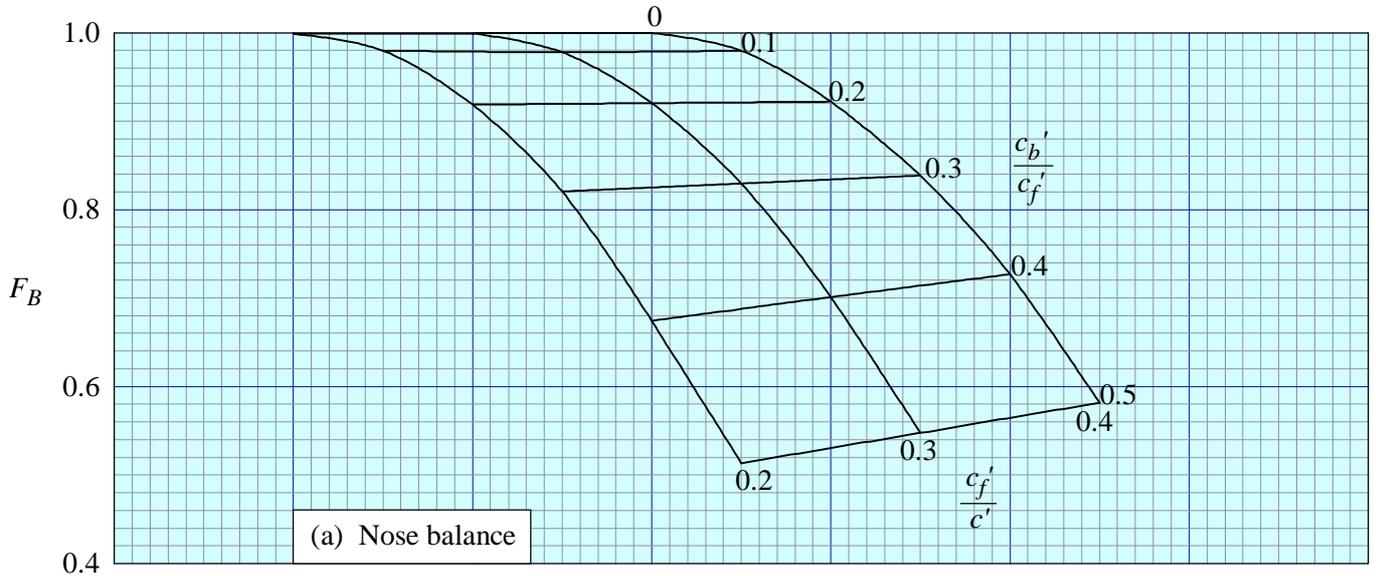


FIGURE 2 FACTOR ON INDUCED CAMBER CONTRIBUTION TO ALLOW FOR CONTROL BALANCE

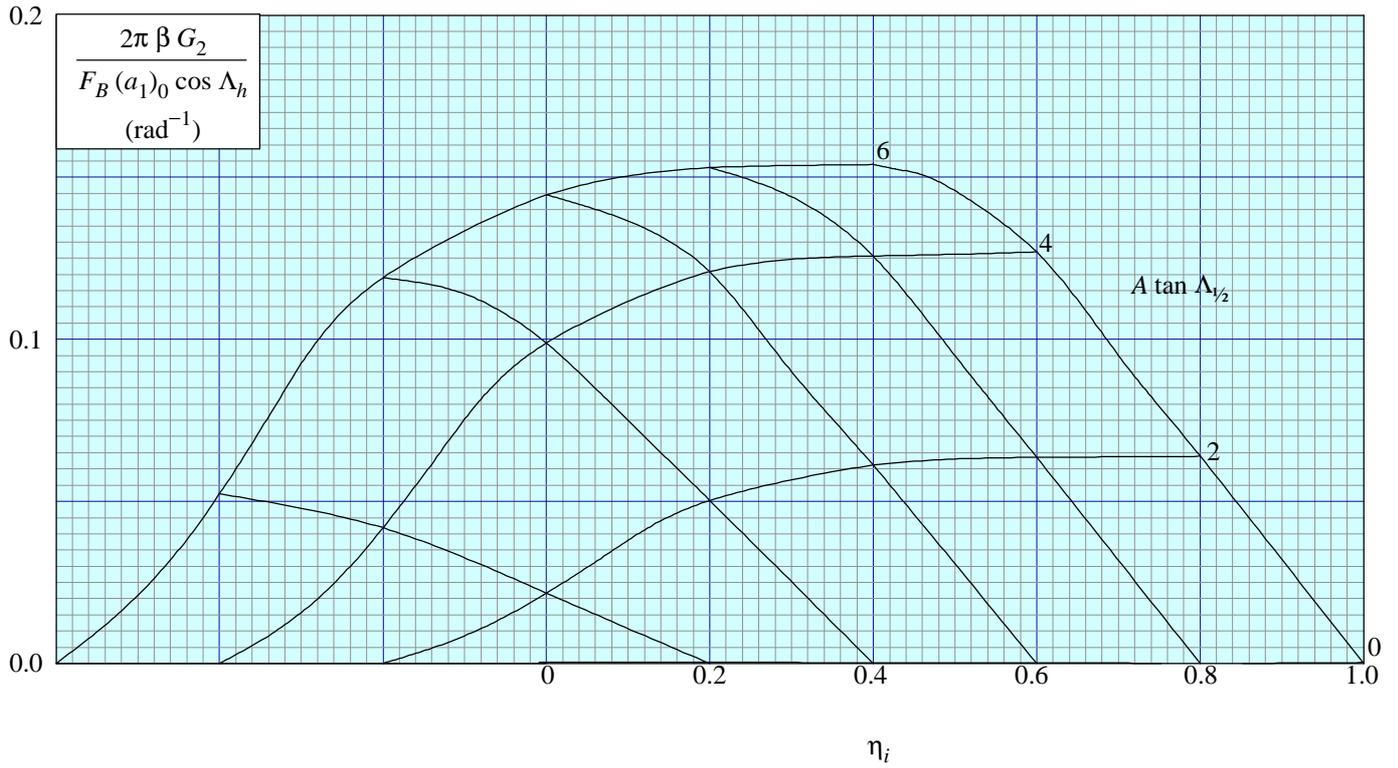


FIGURE 3 INDUCED CAMBER PART-SPAN CORRECTION FOR  $b_1$

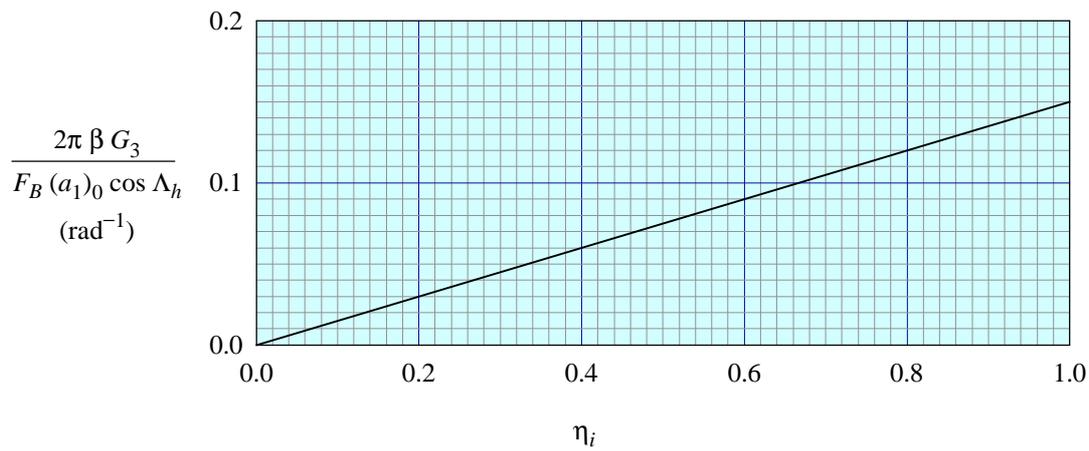
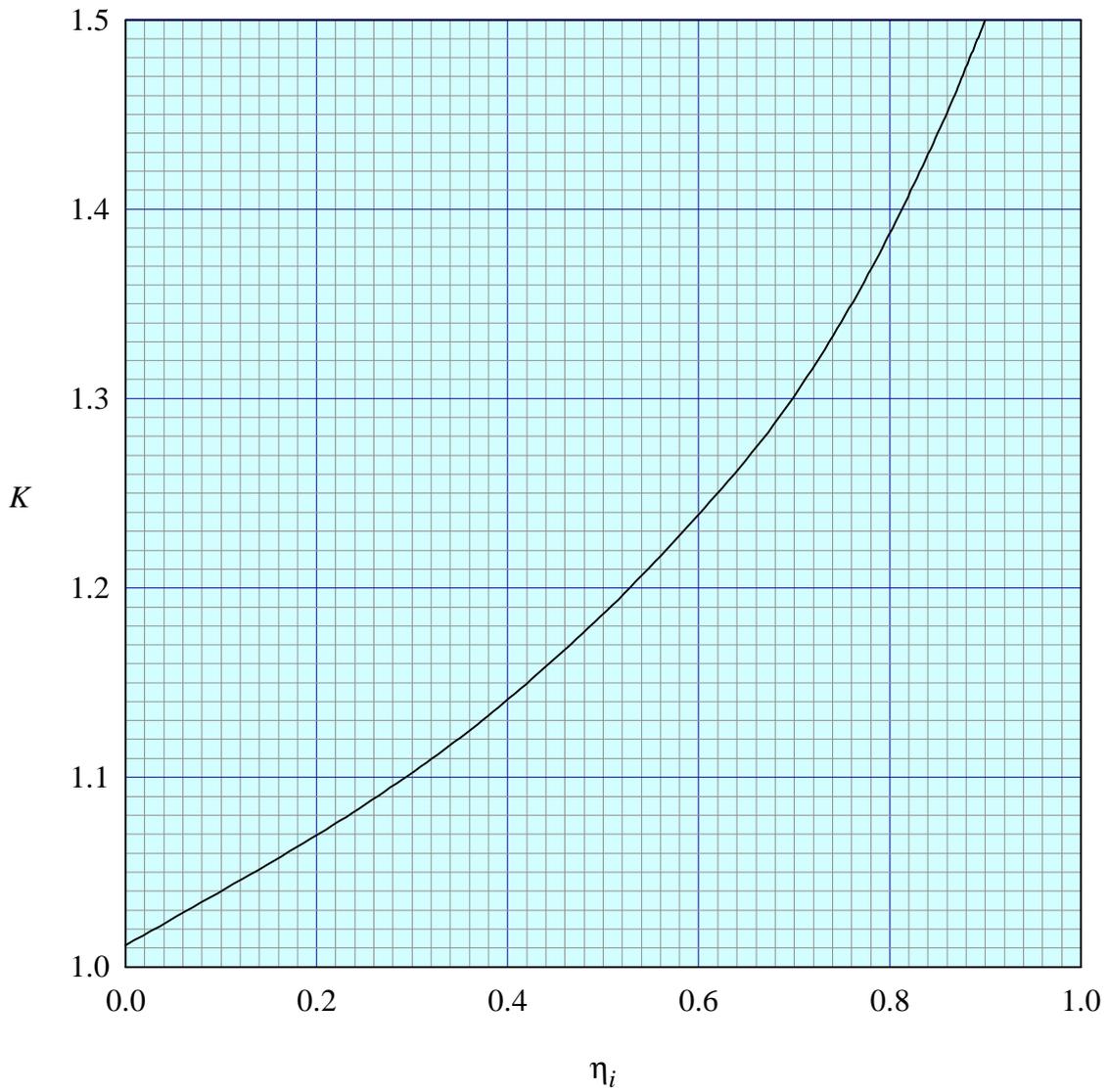


FIGURE 4 ADDITIONAL PART-SPAN CORRECTION FOR  $b_2$



**FIGURE 5 INDUCED ANGLE OF ATTACK CORRECTION FACTOR FOR RECTANGULAR WINGS**

## THE PREPARATION OF THIS DATA ITEM

The work on this particular Item, which supersedes Item No. Aero C.04.01.05 and in part Item No. Aero C.04.01.00, was monitored and guided by the Aerodynamics Committee which first met in 1942 and now has the following membership:

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Mr M.R. Smith – British Aerospace (Commercial Aircraft) Ltd, Bristol

Miss J. Willaume – Aérospatiale, Toulouse, France.

\* Corresponding Member

The technical work in the assessment of the available information and the construction and subsequent development of the Data Item was carried out by

Mr R.W. Gilbey – Senior Engineer.