

# A literature survey on fatigue analysis approaches for rubber

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## Abstract

Rubber components subjected to fluctuating loads often fail due to the nucleation and growth of defects or cracks. The prevention of such failures depends upon an understanding of the mechanics underlying the failure process. This paper reviews analysis approaches that are currently available for predicting fatigue life in rubber. Both crack nucleation and crack growth approaches are considered. A discussion of each approach's strengths and limitations, and examples of how these approaches have been applied in engineering analysis are presented. © 2002 Elsevier Science Ltd. All rights reserved.

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## 1. Introduction

Rubber's ability to withstand very large strains without permanent deformation or fracture makes it ideal for many applications. Applications include tires, vibration isolators, seals, hoses, belts, structural bearings, impact bumpers, medical devices, and footwear, to name a few. These applications impose large static and time-varying strains over a long time. Long-term durability is therefore a critical issue. While many factors contribute to long-term durability, mechanical fatigue, the nucleation and growth of cracks in the rubber, is often the primary consideration. To address the issue effectively and economically, engineers need to model and design for mechanical fatigue early in the product development process. This need has partially been addressed by the development of simulation software capable of predicting stress and strain histories [1–4]. The question remains, however, of how to use these histories to estimate component life.

The objective of this paper is to review analytical approaches that are currently available for predicting fatigue life in rubber. Typically, the fatigue failure process involves two distinct phases. The first phase is a

period during which cracks nucleate in regions that were initially free of observable cracks. The second phase is a period during which nucleated cracks grow to the point of failure. It will be seen that nucleation, growth, and final failure may all be rationalized in terms of the fracture mechanical behavior of rubber. There are, however, issues unique to the crack nucleation phase, which deserve careful study.

Models for predicting fatigue life in rubber follow two overall approaches. One approach focuses on predicting crack nucleation life, given the history of quantities that are defined at a material point, in the sense of continuum mechanics. Stress and strain are examples of such quantities. The other approach, based on ideas from fracture mechanics, focuses on predicting the growth of a particular crack, given the initial geometry and energy release rate history of the crack. For each approach, existing theories are presented. A discussion of each approach's strengths and limitations, and examples of how these approaches have been applied in engineering analysis are also included.

Some of the information presented in this paper has been reviewed previously [5–13]. This literature survey updates these existing reviews to reflect recent and previously unnoticed developments. This survey also offers new interpretations of existing studies and theories, and identifies areas where additional research is needed. Another paper reviews factors that affect the fatigue life of rubber [14]. These include the effects of mechanical

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loading history, environmental effects, effects of rubber formulation, and effects due to dissipative aspects of the constitutive response of rubber.

## 2. Crack nucleation approaches

The crack nucleation approach considers that a material has an intrinsic life determined by the history of stresses or strains at a point. This approach is convenient because it is formulated in terms of stresses and strains, which are familiar to designers. The fatigue crack nucleation life may be defined as the number of cycles required to cause the appearance of a crack of a certain size. The earliest known study of this type was August Wöhler's work with railroad axles in the 1860s [15]. A similar analysis approach was applied to rubber as early as the 1940s [16,17], and remains in use today [18,19]. The approach is particularly appropriate in applications where the initial flaws that eventually determine component life are several orders of magnitude smaller than component features, and where it is desired to analyze the spatial distribution of fatigue life.

The two widely used fatigue life parameters for crack nucleation prediction in rubber are maximum principal strain (or stretch), and strain energy density. The octahedral shear strain has also been used, but less commonly. Strain is a natural choice because it can be directly determined from displacements, which can be readily measured in rubber. When strain energy density is applied to fatigue analysis in rubber, it is often estimated from a hyperelastic strain energy density function, which is defined entirely in terms of strains. Stress, apparently, has rarely been used as a fatigue life parameter in rubber [20]. This seems to be related to the fact that fatigue testing in rubber has traditionally been conducted in displacement control, and that accurate stress determination in rubber components can be formidably difficult.

### 2.1. Maximum principal strain

A common hypothesis [21], implicit in many studies, is that the alternating and mean values of the maximum principal strain uniquely predict nucleation life. Note that it is commonly observed in rubber that cracks initiate on a plane normal to the maximum tensile strain. The earliest fatigue studies in rubber focused on developing an empirical description of the number of cycles to failure as a function of alternating strain and minimum strain. In 1940, Cadwell and co-workers [16] studied unfilled, vulcanized natural rubber. They investigated minimum engineering strains in the range  $-40\%$  to greater than  $+500\%$ , and strain amplitudes in the range  $12.5\%$  to  $350\%$ . They found that, for constant strain amplitude, the fatigue life of natural rubber improves

with increasing minimum strain, up to a moderately high minimum strain level ( $200\%$ ), beyond which additional minimum strain decreased the life. Similar effects were observed in both axial and shear fatigue tests. Several years later, Fielding [17] applied the same approach (based simply on axial engineering strain) in studying the effect of minimum strain on two newly developed synthetic rubbers. In general, for rubbers that strain crystallize, increasing the minimum strain (i.e. increasing  $R$ -ratio) of the strain cycle can significantly lengthen the fatigue life. A detailed discussion of the minimum strain effects on fatigue crack nucleation as well as growth for both crystallized and non-crystallized rubber is presented in [14].

While both uniaxial and shear experiments had been performed in the study of Cadwell et al. [16], it was not attempted to quantitatively reconcile the results from different strain states, or to explicitly develop a theory of how to relate relatively simple lab tests to more complicated strain histories.

Roberts and Benzies [22] and Roach [23] investigated fatigue life under conditions of simple and equibiaxial tension. When plotted against the maximum principal stretch (or strain), fatigue life is longer in simple tension than in equibiaxial tension. The difference was pronounced for natural rubber (NR), and much less pronounced in styrene butadiene rubber (SBR). Ro [24] reanalyzed the data from these studies, using other strain-based parameters, including octahedral shear strain, and maximum shear strain. Ro concluded that none of these parameters were generally optimal for unifying simple and equibiaxial tension data.

### 2.2. Strain energy density

In the late 1950s and early 1960s, success with crack growth models [25–37] had a significant impact on the subsequent development of the nucleation life approach in rubber. Prior to this work, the independent variables in fatigue studies were usually taken as alternating and minimum tensile strain, or stretch. After the development of fracture mechanics for rubber, strain energy density came into use as a parameter to predict fatigue crack initiation [12].

Under certain conditions, the energy release rate is proportional to the product of strain energy density (far from the crack), and the crack size [36,38,39]. When this is the case, it may be considered that the strain energy density is a measure of the energy release rate of naturally occurring flaws. The conditions under which the strain energy density may be uniquely related to the energy release rate are limited. For the relationship to hold, it is assumed that crack growth is self-similar, that the far-field strain gradient across the crack is negligible, and that the stress state is one of simple tension. Several

researchers have investigated strain energy density as a fatigue life parameter in rubber [22–24,40–42].

Roberts and Benzies [22], and Roach [23], found that for NR, equibiaxial tension fatigue life was longer than simple tension fatigue life by a factor of approximately four, when compared based on equal strain energy density. For SBR, equibiaxial tension fatigue life was longer than simple tension fatigue life by a factor of approximately 16, when compared based on equal strain energy density. Note that this ranking is opposite of that found when compared on the basis of maximum principal strain. Roach proposed that these differences could be explained by considering only that portion of the strain energy density that was actually available for flaw growth. For simple tension, Roach proposed that all of the strain energy density is available for flaw growth, while for equibiaxial tension, only one half of the strain energy density is available for flaw growth. This hypothesis gave the best correlation between simple and equibiaxial tension fatigue data.

Ro [24] re-analyzed the data of Roach, and Roberts and Benzies and concluded that strain energy density is a better correlation parameter for high-cycle fatigue of rubber than other strain-based parameters. It should be noted, however, that Ro's analysis is entirely dependent upon assumptions of a contrived dependence of Poisson's ratio on strain, and of linear elastic stress–strain behavior. Curiously, Ro did not further investigate Roach's idea of an available energy density, despite the fact that it appeared to give the most consistent explanation of Roach's results. To this point, no distinction has been made between total strain energy density, and distortional strain energy density. Ro correctly pointed out that, because of rubber's near incompressibility, the distinction is usually unnecessary.

Strain energy density was proposed and studied as a fatigue parameter in metals [43], but the correlation was not satisfactory, and theoretical objections have been raised [44,45]. Findley et al. [43] devised an experiment in which the stresses were cycled, but the strain energy density remained constant. Specimen failure was still observed to occur. Applied as a scalar criterion, strain energy density does not predict the fact that cracking appears in a specific orientation. In addition, the strain energy density cannot be a general measure of energy release rate, since the energy released depends on how the flaw is oriented with respect to the strains.

A number of other approaches have been proposed and evaluated for multiaxial fatigue nucleation life in metals [46–57]. It is particularly worthwhile mentioning the critical plane approaches, which have enjoyed a great deal of success. In critical plane approaches, the history of parameters associated with specific material planes are used to predict fatigue life. For rubber, however, multiaxial loading effects are not yet well understood.

Rubber components are quite commonly subjected to

compressive loading, and this fact must be considered carefully in any analysis of fatigue life. Compressive loading along one direction is almost always associated with simultaneous shear and/or tensile loading in other directions. The only exception would be the case of pure hydrostatic compression. Although planes perpendicular to a compression axis experience closure, planes in other orientations experience shear and/or tension. Cracks will tend to nucleate and grow on these planes. Fatigue crack nucleation criteria (maximum principal strain, strain energy density) which do not consider crack closure may be particularly unreliable for cases involving compressive loading [58].

### 2.3. Applications of crack nucleation approaches

Many engineers and researchers have used strain energy density to correlate analysis results to experimental component life data. Often, such studies refer to the original work of Gent, Lindley, and Thomas [25,36], or the follow-up studies of Lake and Lindley [59,60], reflecting the argument presented in the previous section, that the strain energy density is a measure of the energy release rate of naturally occurring flaws.

Grosch [61] developed a simple analytical model to predict the endurance mileage of tires under various operating conditions. This model uses an analytical estimate of the strain energy density in a tire, along with a semi-empirical relationship between strain energy density and fatigue life. His model predicts relative differences in failure mileage, given differences in operating conditions. The model does not attempt to account for differences due to tire design.

Building on Ro's results [24], DeEskinazi et al. [62] used the Finite Element Method to compute the strain energy density in three tires with design differences. They correlated observed differences in fatigue life to computed strain energy density levels. Oh [63] and Yamashita [64] used strain energy density to predict the fatigue life of bushings and vibration damping devices, respectively.

The presumption of a unique relationship between strain energy density and crack nucleation life is implicit in these studies. While many in the rubber industry have used strain energy density as a predictive parameter, few have tried to ascertain its range of validity under the general conditions experienced by components in service.

Taken together, the studies cited here show that a nucleation life approach to design analysis with rubber is often needed, and that the approach must be general enough to handle situations involving multiaxial loading. Existing approaches have not been very successful in this regard, based on the limited data that have been reported for multiaxial conditions.

### 3. Crack growth approach

The crack growth approach explicitly considers pre-existing cracks or flaws. The idea of focusing attention on individual flaws was introduced by Inglis [65] in 1913, and Griffith [66] in 1920. Griffith proposed a fracture criterion based on an energy balance including both the mechanical energy of a cracked body, and the energy associated with the crack surfaces. Griffith's approach was further developed for rubber by Thomas, Greensmith, Lake, Lindley, Mullins, and Rivlin in the 1950s and 1960s [25–37]. Irwin [67–69], Rice [70], and others developed the approach in metals. While the original application of this approach to rubber was to predict static strength [25,71–78], in the late 1950s Thomas [29] extended the approach to analyze the growth of cracks under cyclic loads in natural rubber. He discovered a square-law relationship between peak energy release rate and crack growth rate for unfilled natural rubber. Paris et al. [79] independently found a similar power-law relationship in metals.

Two important developments in the fracture mechanics of rubber predated the analogous developments in metals. In his work on the J-integral, Rice [70] credits Thomas [26] for first showing the connection between the energy release rate and the strain concentration at the crack tip. Also, Thomas' proposal that the relationship between the cyclic energy release rate and the rate of fatigue crack growth follows a square-law [29], predated Paris' work with a power-law [79] by 3 years.

#### 3.1. The energy release rate

Griffith's [66] hypothesis was that crack growth is due to the conversion of a structure's stored potential energy to surface energy associated with new crack surfaces. He was able to show that the surface energy associated with the crack faces of a broken glass filament was equal to the elastic energy released by the fracture. In rubber, the potential energy released from surrounding material is spent on both reversible and irreversible changes to create the new surfaces [5,25,80]. In any case, the energy release rate is simply the change in the stored mechanical energy  $dU$ , per unit change in crack surface area  $dA$ . In the rubber literature, this quantity is often called the tearing energy  $T$ , regardless of whether the applied loading results in fatigue crack growth or sudden fracture ("tearing").

$$T = -\frac{dU}{dA} \quad (1)$$

The energy release rate was first applied to the analysis of rubber specimens under static loading [25]. It was quickly realized, however, that the concept also applied to crack growth under cyclic loading. In this case, it was found that the maximum energy release rate achieved

during a cycle determined the crack growth rate, for  $R = 0$  cycles [29].

#### 3.2. Fracture mechanics test specimens for rubber

Rivlin and Thomas [25] showed that static crack growth occurs above a critical value of the energy release rate, independent of the type of test specimen they used, suggesting that the critical energy release rate may be considered a true material property. Their initial study used three specimen types: a center-cracked sheet, an edge-cracked sheet, and a "trouser" test piece. A subsequent study by Thomas with three additional test specimen types confirmed the independence of the critical energy release rate from specimen geometry [30]. Other studies [29,36,37,81,82], using the same specimen types, showed that the fatigue crack growth rate is also uniquely determined by the energy release rate.

The aforementioned studies did not investigate the effect of specimen thickness on the fatigue or fracture properties of the material. Thickness effects have been reported by Kadir and Thomas [83], and by Mazich et al. [84]. Kadir and Thomas showed that thickness dependence is related to the development of crack tip roughness. They hypothesized that under hydrostatic tension, rubber can cavitate, and that this may be the cause of the roughness developed around the tip. When the fracture surface remained smooth during growth, little dependence of the growth rate on specimen thickness was observed. Thicknesses ranging from 0.1 mm to 10 mm were investigated. When so called rough, or stick-slip crack growth occurs, however, a thickness effect can arise. On a plot of crack growth rate, at constant energy release rate, as a function of specimen thickness, two plateaus were observed. Below 0.5 mm, and above 5 mm, thickness had little effect. Between these thicknesses, however, the crack growth rate changes by more than an order of magnitude. The crack growth rate for thin specimens was larger than for thick specimens, by an order of magnitude, at equal energy release rate. Their explanation is that the transition in crack growth rate with thickness is related to the scale of the observed crack tip roughness, which is of the same order of magnitude as the transition thickness.

Mazich et al. [84] looked at the dependence of the critical energy release rate for fracture on specimen thickness. They found that, in the range from 1 mm to 3 mm, the critical energy release rate for crack growth increased by a factor of two. Note that these results are broadly consistent with those of Kadir and Thomas, since higher crack growth rates are associated with lower critical energy release rates. Both of these studies were conducted with gum SBR.

The two most commonly used specimens in rubber fatigue crack growth studies are the single edge cut planar tension specimen and single edge cut simple ten-

sion specimen. The planar tension specimen geometry (also known as the pure shear specimen in the rubber literature) has proven to be especially useful in fracture mechanics tests for rubber [81,85]. This specimen is short and wide such that lateral contraction is prevented by the grips, while an axial strain is introduced in the direction of the short dimension. The pure shear designation arises because, for small, volume-conserving strains, this strain state is one of pure shear, in a coordinate system suitably rotated with respect to the specimen.

In the single edge cut, planar tension specimen, the energy release rate  $T$  has an especially simple form, provided that the cut is sufficiently deep and that growth of the cut results only in translation of the crack tip fields. The expression depends only on the strain energy density  $W$  remote from the crack and from specimen edges, and the specimen height  $h$ . Note that, in displacement control, the energy release rate for this specimen is independent of the crack size [25].

$$T = Wh \tag{2}$$

In the single edge cut, simple tension specimen, the energy release rate  $T$  depends on the gauge section strain energy density  $W$ , the size of the crack  $a$ , and a strain-dependent parameter  $k$ .

$$T = 2kWa \tag{3}$$

The dependence of the parameter  $k$  on strain was studied by Greensmith [34], and Lindley [72], for the case in which the crack is much smaller than the specimen width. Nait-Abdelaziz et al. [82] have extended the parameter for cracks of finite size, relative to specimen width. A plot of Greensmith's data is shown in Fig. 1. A practical approximation of the data is given in terms of the engineering strain  $\epsilon$  by [72]:

$$k = \frac{2.95 - 0.08\epsilon}{(1 + \epsilon)^{1/2}} \tag{4}$$

The trouser specimen, shown in Fig. 2, was used in early studies of fatigue crack growth in rubber [29] to demonstrate geometry independence of the relationship between the energy release rate and the fatigue crack growth rate. The energy release rate  $T$  depends on the applied force  $F$ , the extension ratio  $\lambda$  and strain energy density  $W$  in the "legs" of the specimen, the specimen thickness  $t$ , and the "leg" width  $b$ .

$$T = \frac{2F\lambda}{t} - bW \tag{5}$$

Energy release rate calculations have been developed for other specimen types that have occasionally been used in fracture mechanics studies of rubber. These include the so-called angled, and split specimens. Since these have been reviewed previously [5–

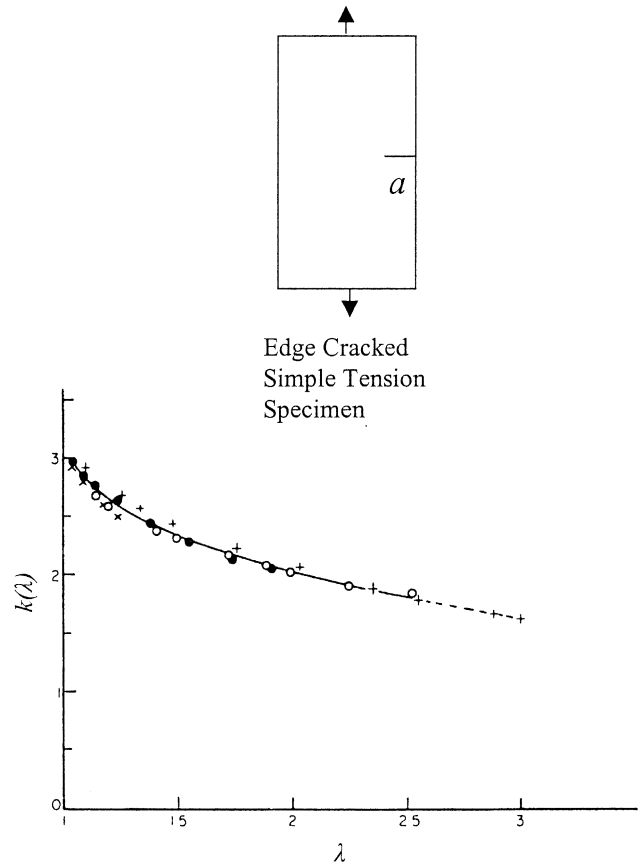


Fig. 1. Greensmith's [34] data for variation of  $k$  with maximum principal stretch  $\lambda$ , in the energy release rate of a crack in the simple tension specimen,  $T = 2kWa$ . Different data point types are for different rubber compounds.

11,25,30,59,81,85,86], and since their application has been limited primarily to static strength measurements, the corresponding energy release rate calculations are not given here.

### 3.3. Regimes of fatigue crack growth

Lake and Lindley [59] identified four distinct regimes of fatigue crack growth behavior, based on the maximum energy release rate per cycle,  $T$ , for  $R = 0$  cycles in rubber. The full range of behavior is shown in Fig. 3, for unfilled NR and SBR.

So long as the peak energy release rate  $T$  remains below a threshold  $T_0$ , crack growth proceeds at a constant rate  $r$ , due solely to environmental attack. The crack growth rate  $da/dN$  below  $T_0$  is independent of the mechanical loading, and is denoted Regime 1.

$$\frac{da}{dN} = r \quad T \leq T_0 \tag{6}$$

There is then a range of  $T$ , between  $T_0$  and  $T_i$ , over which there is a transition. The transition is described

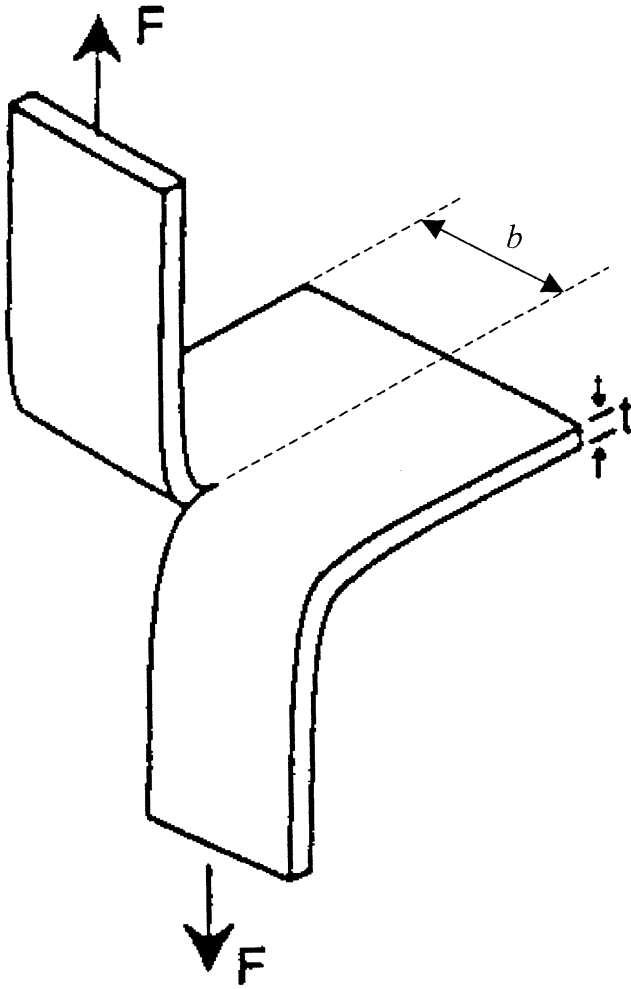


Fig. 2. Trouser test specimen used in early fatigue crack growth studies.

by the following relationship, in which  $A$  is a material property. This is denoted Regime 2.

$$\frac{da}{dN} = A(T - T_0) + r \quad T_0 \leq T \leq T_t \quad (7)$$

After the transition, there is a range between  $T_t$  and  $T_c$ , over which the relationship between the fatigue crack growth rate and the energy release rate obeys a power-law. The associated material properties are  $B$  and  $F$ . This is denoted Regime 3.

$$\frac{da}{dN} = BT^F \quad T_t \leq T \leq T_c \quad (8)$$

Finally, beyond  $T_c$ , unstable crack growth ensues. In this regime, the crack growth rate is essentially infinite. This is denoted Regime 4.

$$\frac{da}{dN} = \infty \quad T = T_c \quad (9)$$

Aglan and Moet [87] developed a single relationship that predicts Regimes 2, 3, and 4, for  $R = 0$  loading.

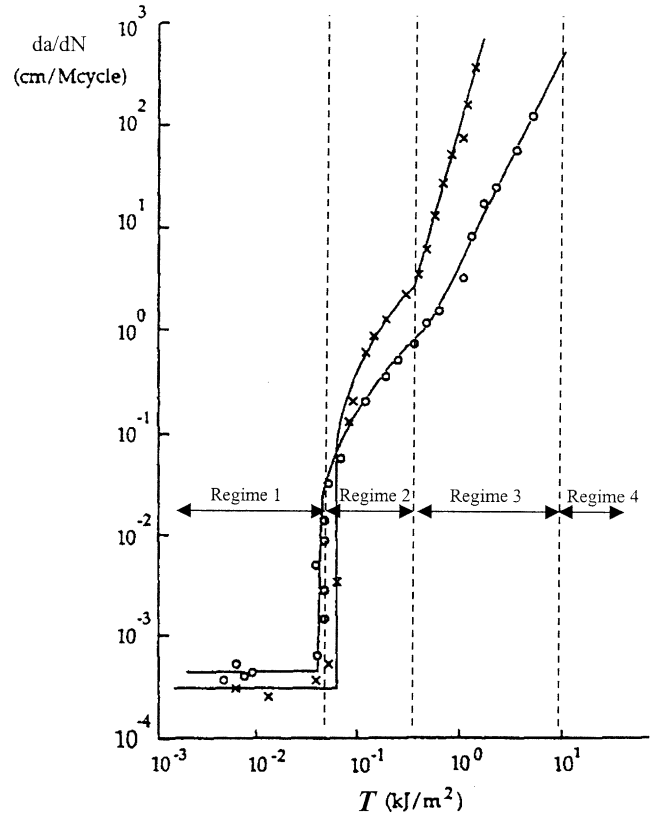


Fig. 3. Regimes of fatigue crack growth behavior in unfilled rubber under  $R = 0$  loading ( $\times$ , SBR;  $\circ$ , NR) [59].

Their model, the Crack–Layer theory for rubber, is based on the irreversible thermodynamics of an “active zone”, preceding the crack tip [87–89]. The model takes the thermodynamic flux to be the crack growth rate, and the conjugate driving force to be the energy release rate range  $\Delta T$ . The model assumes that the energy dissipation associated with crack growth is proportional to the square of the energy-release rate range, in analogy to the Dugdale strip-yield model [90–92]. In addition to the critical energy release rate  $T_c$ , the model uses material parameters  $\beta$  and  $\mu$ .

$$\frac{da}{dN} = \frac{\beta \Delta T^2}{\mu T_c - \Delta T} \quad (10)$$

Chow and Lu also developed a multi-regime model, based on thermodynamic arguments [93,94]. Their model is claimed to be valid for a wide range of materials. In addition to the critical energy release rate  $T_c$ , the model uses material parameters  $\beta$  and  $m$ .

$$\frac{da}{dN} = \frac{\beta \Delta T^m}{T_c - T_{\max}} \quad (11)$$

These multi-regime models resemble models proposed for fatigue crack growth in metals by Foreman [95], and by Weertman [96]. Foreman’s model depends on the  $R$

ratio, the range of the stress intensity factor  $\Delta K$ , the fracture toughness  $K_c$ , and material parameters  $\beta$  and  $n$ .

$$\frac{da}{dN} = \frac{\beta \Delta K^n}{(1-R)K_c - \Delta K} \quad (12)$$

Weertman's model depends also on the peak stress intensity factor,  $K_{\max}$ .

$$\frac{da}{dN} = \frac{\beta \Delta K^4}{K_c^2 - K_{\max}^2} \quad (13)$$

In Foreman's model, the resemblance can be seen by taking  $n = 4$ , and comparing to Aglan and Moet mode with  $\mu = 1$ . Note that while the Foreman equation is in terms of the stress intensity factor, the Aglan–Moet and Chow–Lu equations are in terms of energy release rate  $T$ . The energy release rate is proportional to the square of the stress intensity factor. All models also contain a proportionality constant, herein denoted by the common symbol  $\beta$  to emphasize algebraic similarity.

In some elastomers, the occurrence of compressive loading during a fatigue cycle can also have a dramatic effect on crack growth rate through the mechanism of strain crystallization at the crack tip. For cases in which the minimum tensile stress of the crack tip during the load cycle is sufficient to induce crystallization, the material remains crystalline at all times, retarding crack growth. For cases where the loading conditions permit the crack tip crystalline region to “melt” at some point during the cycle, the fatigue crack growth rate increases [97]. The effect of  $R$ -ratio on strain-crystallizing rubbers can influence crack growth rate and fatigue life by several orders of magnitude.

A major shortcoming of the aforementioned models, when applied to strain-crystallizing rubbers, is that they predict increased crack growth rates for  $R > 0$  conditions. In contrast, for strain-crystallizing rubbers, the crack growth rate is significantly retarded by  $R > 0$  conditions, as shown by Lindley [98]. Environmental effects can interact significantly with  $R$  ratio effects [97]. It appears that no current multi-regime fatigue crack growth model predicts the observed  $R > 0$  effects for an important class of rubbers.

### 3.4. Relationship of energy release rate to crack tip conditions

The success of the energy release rate as a parameter for predicting fatigue crack growth and fracture has been attributed to its unique relationship to the local conditions at the crack tip. Thomas was the first to demonstrate this relationship experimentally [26]. He studied the strain distribution around the tip of a model crack in a sheet of rubber, in several specimen geometries. He found that the average strain energy density of material surrounding an idealized crack tip was uniquely related

to the energy release rate, independent of the specimen type.

The relationship between the energy release rate and local crack tip conditions in rubber has also been confirmed in independent studies by Andrews [99,100], Knauss [101], Lee and Donovan [102], and Morman et al. [103]. Andrews used a microscopic, photoelastic technique to quantify the strain field around the crack tip. He showed that a combination of hysteresis and large displacements result in blunting of the crack tip in highly deformable materials. Knauss used a printed-grid technique. Lee and Donovan used Thomas' original approach. Morman et al. developed an analytical expression relating the energy release rate to the crack tip radius. The studies of Andrews and Knauss produced a more detailed map of the crack tip strain distribution than Thomas's; and all studies confirmed Thomas' conclusion.

Rice's development of the J-integral [70] provided a mathematical argument to explain the relationship between the energy release rate and the local crack tip conditions. The J-integral expresses an energy balance on a volume surrounding a crack tip. Rice proved that the value of the integral is independent of the choice of integration path. Since the integral is independent of the integration path, the path can be chosen close to the crack tip. The integral is therefore a measure of local crack tip conditions. The integration path may also be chosen to follow the boundaries of the specimen. In this case, the integral turns out to be equivalent to the energy release rate [84,102]. The energy release rate is therefore a measure of the intensity of local crack tip fields, for a given material and crack tip geometry. Rice's original formulation was valid for nonlinear elastic materials and infinitesimal strains. Chang generalized the J-integral for nonlinear elastic materials at finite strains [104].

A practical consequence of the J-integral is that the details of the processes occurring at the crack tip often do not need to be quantified in any way other than the J-integral in order to model crack growth. Instead, the details are accounted for by treating them as intrinsic to the material/crack-tip system. In this manner, nonlinearities due to finite network extensibility [105], strain crystallization [16,17,106–108], frictional losses due to filler interactions [109–111], and the Mullins effect [112–114] may be rolled into the fracture properties of the material [5,6]. A general theory to account for the effects of general nonlinear, dissipative constitutive behavior on crack tip fields has been proposed by Andrews [115]. Of course, this approach remains subject to the assumption that the crack propagates through an isotropic, homogeneous continuum. It seems that the energy release rate approach avoids the necessity of modeling crack tip dissipative processes by focusing on where the energy for driving the crack comes from (from strain energy stored beyond the J-integral boundary), instead

of where that energy is spent (on dissipative processes near the crack tip).

In situations where the stress–strain behavior is strongly time-dependent, the J-integral does not uniquely characterize local crack tip conditions. Lindley presented an approach for addressing time-dependent crack growth in SBR [116]. A time-dependent path-integral approach has also been developed [117–121], but it does not appear that this approach has been applied to elastomers.

Compressive loading must be considered carefully in any analysis of crack growth [122]. One consideration is that crack growth can only occur due to loads which result in tensile stresses at the crack tip. Purely compressive loading of a crack results in closure of crack faces until contact is achieved. Further compressive loading is then transmitted across the crack without causing crack growth. The strain energy release rate in this case is zero, despite the fact that the compressive loading may be large. When compressive and shear loads are both present, the crack tip can experience tensile loading. Determination of the strain energy release rate in this case can be complex, depending not only on the crack geometry and loading conditions, but also on the frictional properties of the crack faces.

Of course, the details of crack-tip processes are of interest when a deeper understanding of the failure process is sought. Theories for the relationship between the failure properties of rubber and rubber's molecular characteristics have been proposed and studied by several researchers [123–125].

### 3.5. Applications of the crack growth approach

An early design application of the fracture mechanics approach in rubber came from Lake and Clapson [126]. They developed an estimate of the energy release rate cycle at the base of tread pattern grooves in tires, based on the crack mouth opening displacement. They successfully predicted the rate of growth of cracks in tires on a road test from fatigue crack growth data generated in the lab.

Southern and Thomas [127] applied a fracture mechanics approach to develop a model for abrasive wear of rubber. Their model is based on fatigue crack growth at the base of an incipient wear particle. The model relates observed wear rates to measured fatigue crack growth properties, for the particular wear mechanism studied.

Huang and Yeoh [128] developed a fatigue crack growth model to rationalize the nucleation phase of the belt edge separation process in tires. The model was based on an estimate of the energy release rate of an array of penny-shaped cracks, each located at the end of a cord. Their model was validated via fatigue testing of a model cord–rubber composite. Choi, Roland, and Bissonette [129] analyzed failure in a large, elastomeric torpedo launcher. Medri and Strozzi [130] analyzed

crack growth in elastomeric seals. Lindley [131] analyzed the life of metal–rubber bonds using a fracture mechanics approach. Stevenson et al. [132], and Gundersen et al. [133] have used FEA to evaluate energy release rate cycles experienced by cracks in rubber supports on off-shore structures. Stevenson and Malek [134] developed a model to predict the puncture behavior of thick rubber components penetrated by a sharp-cornered cylindrical indenter.

Ebbott [135] and Wei et al. [136] have used FEA to evaluate the energy release rate cycle experienced by a crack at the edge of cord–rubber tire components. Ebbott used a global–local analysis procedure in which a coarse mesh of the whole tire was first analyzed, followed by a refined-mesh analysis involving only the region of interest. Wei et al. performed the analysis with a single mesh. In both cases, the assumed crack geometry was built into the FE model. Both studies reported reasonable estimates of the crack growth rate for the tires analyzed. A limitation of this approach is that each potential failure mode requires its own mesh and analysis. In addition, each mesh applies only to a given crack size. A full life analysis with this approach requires large effort and expense in creating and analyzing multiple models. Automated mesh adaptation is required to make these approaches suitable for general use.

A common difficulty of using the crack growth approach in rubber is that it requires up-front knowledge of the initial location and state of the crack that causes the final failure. Often, this information is not available, or it is the very information the designer needs to predict. In addition, the changing geometry of the problem must be considered. Numerical implementations of a direct fracture mechanics approach remain labor intensive and computationally expensive. There is a great need for robust, general-purpose algorithms for crack growth analyses in rubber.

## 4. A flaw growth model for crack nucleation

A crack growth approach has been successfully used to predict uniaxial fatigue crack nucleation life from fatigue crack growth measurements [36,59]. This unified analysis is based on integration of the fatigue crack growth rate. The energy release rate of an assumed, pre-existing flaw is estimated from the flaw size and the strain energy density. The analysis only applies to small cracks, when the energy release rate may be factored into a product proportional to both strain energy density and crack size. Nevertheless, the small crack assumption often covers the most important portion of a component's life, since the presence of a large crack usually means the component has already failed.



#### 4.1. Integrated power-law model

Fatigue life is ultimately determined by a pre-existing flaw that is the first to grow to a critical size. The life is obtained by integrating the growth rate of the fastest growing flaw, from its initial size to its critical or final size. For the purposes of integration, it is assumed that flaw growth is planar and self-similar.

Piece-wise integration over the four regimes of fatigue crack growth behavior is possible, and yields the most accurate fatigue life predictions [9,137]. A practical shortcut assumes power-law behavior over the entire life of the flaw [5,85,138]. This results in a closed-form relationship between the fatigue life and the fatigue crack growth properties. Combining Eqs. (8) and (3),

$$\frac{da}{dN} = f[T(a,W)] = BT^F = B(2kW a)^F \quad (14)$$

Then integrating,

$$N_f = \int_0^{N_f} dN = \int_{a_0}^{a_f} \frac{1}{f[T(a,W)]} da = \int_{a_0}^{a_f} \frac{1}{B(2kW)^F a^{-F}} da \quad (15)$$

$$N_f = \frac{1}{F-1} \frac{1}{B(2kW)^F} \left[ \frac{1}{a_0^{F-1}} - \frac{1}{a_f^{F-1}} \right] \quad (16)$$

From Eq. (16), we see that if the initial flaw size  $a_0$  is much smaller than the critical flaw size,  $a_f$ , then the life becomes independent of the critical flaw size.

$$N_f = \frac{1}{F-1} \frac{1}{B(2kW)^F} \frac{1}{a_0^{F-1}} \quad (17)$$

If the initial flaw size is regarded as a property intrinsic to a given virgin material, and if variation of  $k$  with strain is neglected, the constants of the equation may be combined into a single material constant  $D$ .

$$N_f = DW^{-F} \quad (18)$$

This derivation applies only to uniaxial loading, where the energy release rate can be factored into the strain energy density and the crack size. For multiaxial situations, not all of the elastically stored energy is available to be released [23]. A general-purpose calculation of the available strain energy density has not been published.

#### 4.2. Intrinsic flaws in rubber

The preceding theory suggests a way to estimate the size of naturally occurring flaws in rubber, by using the initial flaw size as a curve fit parameter to obtain agreement between crack nucleation and crack growth experiments [34,59,139]. The resulting flaw size is actually an effective flaw size, reflecting both the size and shape of the flaws [36]. Effective flaw sizes in the range of  $20 \times 10^{-6}$  m to  $50 \times 10^{-6}$  m were observed in a study

by Lake and Lindley [59], which covered eight different polymer types, and various fillers, curatives, and other compounding variables. It has been confirmed that the measured flaw size is independent of temperature [140]. Flaw size has some dependence on crosslink density [59,141], carbon black type [59,142], and degree of dispersion of compound ingredients [138]. The flaw size can also be deduced from static strength measurements [10,140], and optical microscopy techniques [142], independent of fatigue measurements. Agreement of these methods to within a factor of 2 has been reported [139,142]. Damage characterization of elastomeric composites using X-ray attenuation has been reported by Bathias et al. [143].

A basic assumption of fracture mechanics is continuity and homogeneity of the material. In materials where the initial flaws are small enough that this assumption is not true, additional considerations are necessary, as is the case in metals [144–146]. In rubber, the intrinsic flaws are larger than features of the molecular network structure by a factor of more than 10,000, and larger than individual filler particles by a factor of more than 100. Agglomerations of carbon black particles can exhibit dimensions of the same scale of intrinsic flaws [142]. The actual scale of such features seems likely to depend on manufacturing processes such as mixing. Table 1 summarizes the size scales of importance in filled rubber.

The independent agreement of nucleation and growth approaches, using initial flaw size as the sole fitting parameter suggests that it is appropriate to assume that the initial flaws are embedded in continuous, homogeneous material. The precise nature of such flaws remains obscure because it appears that there are multiple sources for flaws of the observed effective size. These sources may include naturally occurring contaminants or voids in the base polymer, imperfectly dispersed compounding

Table 1  
Geometric features of filled rubber

Feature	Size
Large carbon black agglomerate [142]	$200 \times 10^{-6}$ m
Smallest flaw visible to naked eye	$100 \times 10^{-6}$ m
Typical size of intrinsic defects [139]	$40 \times 10^{-6}$ m
Small carbon black agglomerate [142]	$20 \times 10^{-6}$ m
Coarse carbon black particle [147]	$500 \times 10^{-9}$ m
Fine carbon black particle [147]	$10 \times 10^{-9}$ m
Distance along polymer chain between crosslinks (assuming 300% macro-stretch = 100% chain extension, crosslink density of $5.8 \times 10^{19}/\text{cm}^3$ ) [148]	$1 \times 10^{-9}$ m
Length of a single monomer unit (isoprene) [123,149]	$500 \times 10^{-12}$ m
Spatial distance between crosslinks (based on crosslink density of $5.8 \times 10^{19}/\text{cm}^3$ ) [148]	$300 \times 10^{-12}$ m
Length of a main-chain, polysulfidic bond [150]	$100 \times 10^{-12}$ m

ingredients, filler agglomerates, mold lubricants, and imperfections in mold surfaces.

## 5. Summary

Two approaches have developed for analyzing fatigue life in rubber components, the crack nucleation approach, and the crack growth approach. In rubber, the crack growth approach has been studied and used extensively. The nucleation approach has received less attention in the literature, although many engineers still use this approach for its simplicity and familiarity.

The nucleation approach is advantageous for analyzing the spatial distribution of fatigue life, since it is based on quantities that are defined at a material point, in the sense of continuum mechanics. In rubber, uniaxial fatigue life results are most commonly correlated based on maximum principal strain (or stretch), and strain energy density. Neither of these parameters has been robustly successful in correlating results from different strain states, particularly simple tension and equibiaxial tension.

Some success has been achieved in developing and applying the crack growth approach in rubber. A major practical challenge is computation of the energy release rate associated with the crack of interest, and predicting the location and path of the fastest growing crack, especially when the geometry and loading are complicated. Robust numerical procedures are inevitably required, but are not widely available. When the crack of interest is small, another problem is determining the initial size and shape of the crack. Small flaws are often of particular importance, since most of a component's life may be spent on the growth of small flaws.

For uniaxial situations in which failure initiates from a small flaw, the strain energy density can be used to estimate the energy release rate of the flaw, from which fatigue life can be computed, given the fatigue crack growth curve. For multiaxial situations, the strain energy density is not generally appropriate because not all of the energy is available to be released by the growth of a flaw. An adequate multiaxial nucleation life approach is needed to accurately predict fatigue life in rubber components.

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