



# 7600026 – Laboratório de Óptica

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**Índice de Refração**

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## CLASSICAL ELETROMAGNETISM

In uniform isotropic linear media, the wave equation is:

$$\nabla^2 \mathbf{E} - \mu\epsilon \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0$$

$$\nabla^2 \mathbf{H} - \mu\epsilon \frac{\partial^2 \mathbf{H}}{\partial t^2} = 0$$

They are satisfied by plane wave

$$\psi = A e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

$$k = |\mathbf{k}| = \omega \sqrt{\mu\epsilon}$$

$\psi$  can be any Cartesian components of  $\mathbf{E}$  and  $\mathbf{H}$

The phase velocity of plane wave travels in the direction of  $\mathbf{k}$  is

$$v = \frac{\omega}{k} = \frac{1}{\sqrt{\mu\epsilon}}$$



We can define the index of refraction as

$$n = \frac{v}{c} = \sqrt{\frac{\mu\epsilon}{\mu_0\epsilon_0}}$$

Most media are nonmagnetic and have a magnetic permeability  $\mu=\mu_0$ , in this case

$$n = \sqrt{\frac{\epsilon}{\epsilon_0}}$$

In most media,  $n$  is a function of frequency.



## Classical Electron Model ( Lorentz Model)

Let the electric field of optical wave in an atom be

$$\mathbf{E} = \mathbf{E}_0 e^{-i\omega t}$$

the electron obeys the following equation of motion

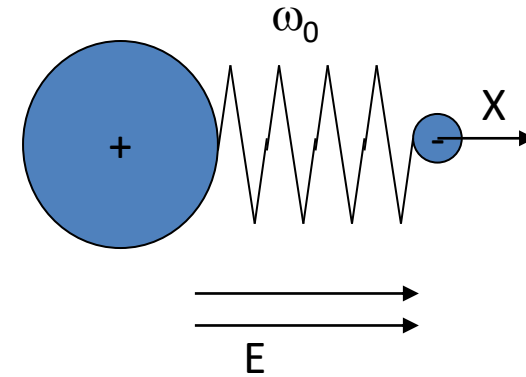
$$m \frac{d^2}{dt^2} \mathbf{X} + m\gamma \frac{d}{dt} \mathbf{X} + m\omega_0^2 \mathbf{X} = -e\mathbf{E}$$

$\mathbf{X}$  is the position of the electron relative to the atom

$m$  is the mass of the electron

$\omega_0$  is the resonant frequency of the electron motion

$\gamma$  is the damping coefficient





The solution is 
$$\mathbf{X} = \frac{-e\mathbf{E}_0}{m(\omega_0^2 - \omega^2 - i\omega\gamma)} e^{-i\omega t}$$

The induced dipole moment is

$$\mathbf{p} = -e\mathbf{X} = \frac{e^2}{m(\omega_0^2 - \omega^2 - i\omega\gamma)} \mathbf{E} = \alpha\mathbf{E}$$

$\alpha$  is atomic polarizability

$$\alpha = \frac{e^2}{m(\omega_0^2 - \omega^2 - i\omega\gamma)}$$

The dielectric constant of a medium depends on the manner in which the atoms are assembled. Let  $N$  be the number of atoms per unit volume. Then the polarization can be written approximately as

$$\mathbf{P} = N \mathbf{p} = N \alpha \mathbf{E} = \epsilon_0 \chi \mathbf{E}$$



The dielectric constant of the medium is given by

$$\varepsilon = \varepsilon_0 (1 + \chi) = \varepsilon_0 (1 + N\alpha / \varepsilon_0)$$

If the medium is nonmagnetic, the index of refraction is

$$n = (\varepsilon / \varepsilon_0)^{1/2} = (1 + N\alpha / \varepsilon_0)^{1/2}$$

$$n^2 = \frac{\varepsilon}{\varepsilon_0} = 1 + \frac{Ne^2}{\varepsilon_0 m(\omega_0^2 - \omega^2 - i\omega\gamma)}$$

If the second term is small enough then

$$n = 1 + \frac{Ne^2}{2\varepsilon_0 m(\omega_0^2 - \omega^2 - i\omega\gamma)}$$

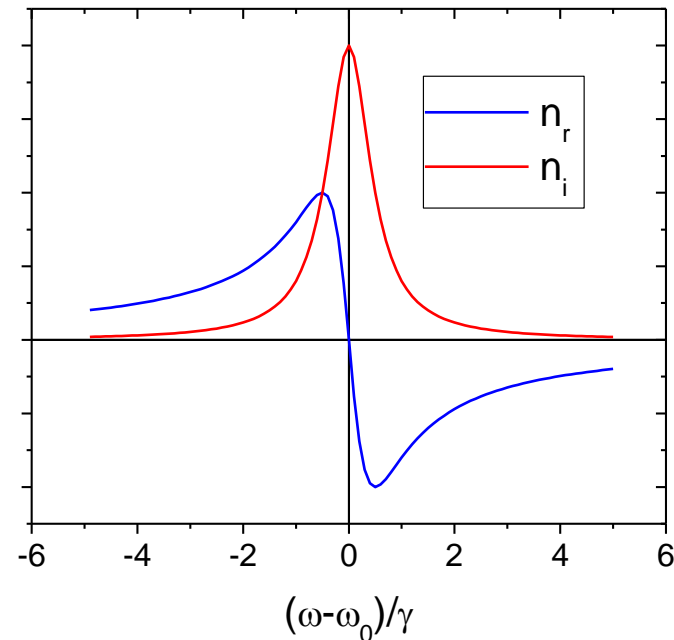


The complex refractive index is

$$n \rightarrow n_r + in_i = 1 + \frac{Ne^2(\omega_0^2 - \omega^2)}{2\varepsilon_0 m [(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2]} + i \frac{Ne^2 \gamma \omega}{2\varepsilon_0 m [(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2]}$$

at  $\omega \sim \omega_0$ ,

$$n_r + in_i = 1 + \frac{Ne^2(\omega_0 - \omega)}{4\varepsilon_0 m \omega_0 [(\omega_0 - \omega)^2 + (\gamma/2)^2]} + i \frac{Ne^2 \gamma}{8\varepsilon_0 m \omega_0 [(\omega_0 - \omega)^2 + (\gamma/2)^2]}$$



Normalized plot of  $n-1$  and  $k$  versus  $\omega-\omega_0$



For more than one resonance frequencies for each atom,

$$n^2 = 1 + \frac{Ne^2}{\epsilon_0 m} \sum_j \frac{f_j}{(\omega_j^2 - \omega^2 - i\omega\gamma_j)} \quad \sum_j f_j = Z$$

## Classical Electron Model ( Drude model)

If we set  $\omega_0=0$ , the Lorentz model become Drude model. This model can be used in free electron metals

$$n^2 = 1 - \frac{Ne^2}{\epsilon_0 m(\omega^2 + i\omega\gamma)}$$





## Relation Between Dielectric Constant and Refractive Index

By definition,

$$n^2 = \frac{\varepsilon}{\varepsilon_0}$$

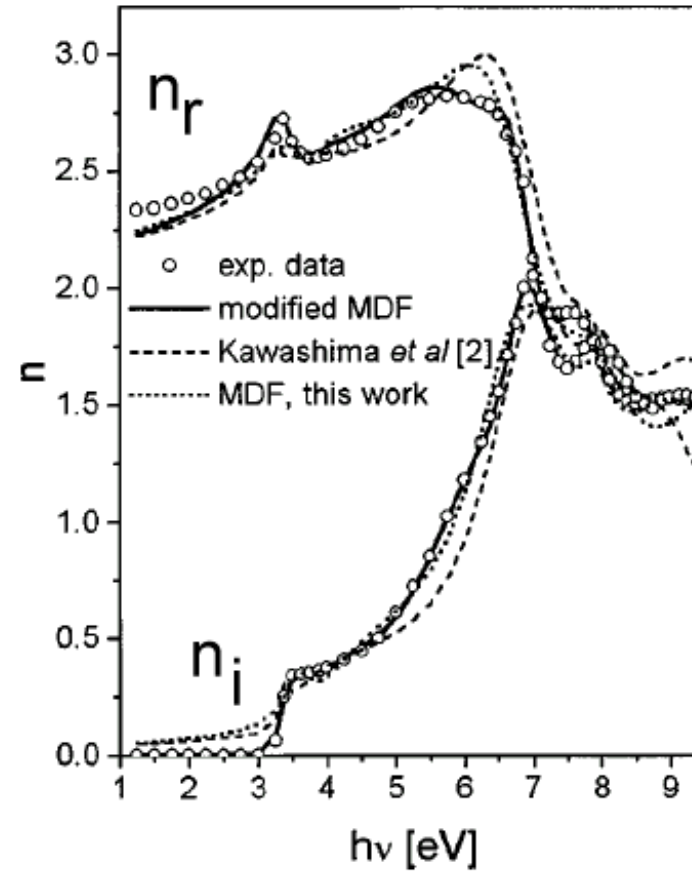
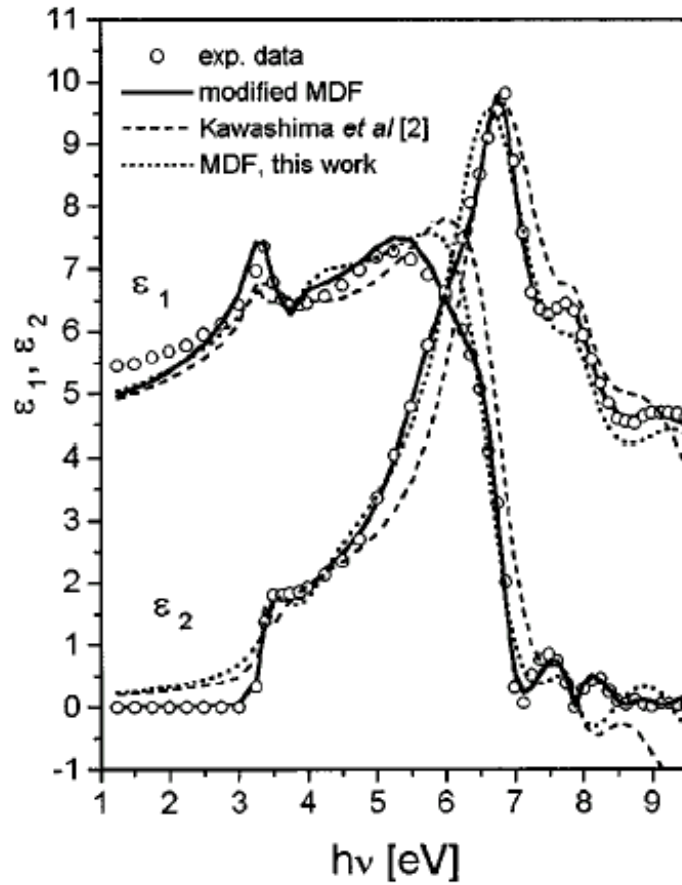
$$n = n_r + in_i$$

$$\varepsilon = \varepsilon_1 + i\varepsilon_2$$

We can easily get:

$$n_r = \left\{ \frac{1}{2} [(\varepsilon_1^2 + \varepsilon_2^2)^{1/2} + \varepsilon_1] \right\}^{1/2} / \varepsilon_0$$

$$n_i = \left\{ \frac{1}{2} [(\varepsilon_1^2 + \varepsilon_2^2)^{1/2} - \varepsilon_1] \right\}^{1/2} / \varepsilon_0$$



$$\epsilon(E) = \sum_{\alpha=A,B,C} \left( \sum_{n=1}^{\infty} \frac{A_{0\alpha}^{\text{ex}}}{n^3} \frac{1}{E_{0\alpha} - (G_{0\alpha}^{3D}/n^2) - E - i\Gamma} \right)$$

Real and imaginary part of the index of refraction of GaN vs. energy;



The real part and imaginary part of the complex dielectric function  $\varepsilon(\omega)$  are not independent. they can be connected by Kramers-Kronig relations:

$$\varepsilon_1(\omega) = \varepsilon_0 + \frac{2}{\pi} P \int_0^{\infty} \frac{\varepsilon_2(\omega') \omega'}{\omega'^2 - \omega^2} d\omega'$$

$$\varepsilon_2(\omega) = \frac{2\omega}{\pi} P \int_0^{\infty} \frac{\varepsilon_1(\omega') - \varepsilon_0}{\omega'^2 - \omega^2} d\omega'$$

P indicates that the integral is a principal value integral.

K-K relation can also be written in other form, like

$$n(\lambda) = \frac{1}{\pi} P \int_0^{\infty} \frac{\alpha(\lambda')}{1 - (\lambda'/\lambda)^2} d\lambda'$$



## Model:

Light: electromagnetic wave

Atom and molecule: classical dipole oscillator

$$n(\omega), \alpha(\omega)$$

$$E(z, t) = E_0 e^{i(k \cdot z - \omega t)}, \quad k = (n + i\kappa)\omega / c$$
$$= E_0 e^{-\alpha z} e^{i(\omega n z / c - \omega t)}.$$

Two propagation parameters:

$$n, \alpha$$

## Propagation of light in a dense optical medium

Three types of oscillators:

1. bound electron (atomic) oscillator
2. vibrational oscillator;
3. free electron oscillators

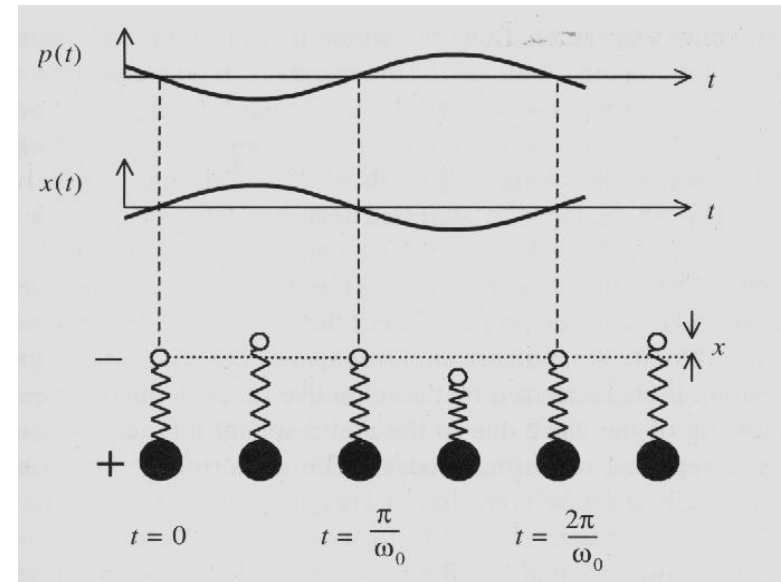
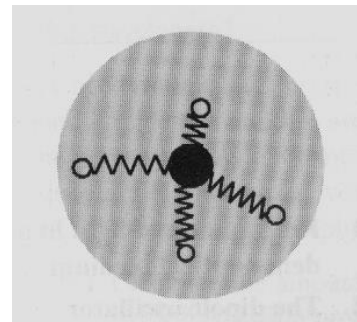
### Atomic oscillators

$$\frac{1}{\mu} = \frac{1}{m_0} + \frac{1}{m_N} \approx \frac{1}{m_0},$$

$$\omega_0 = \sqrt{\frac{K_S}{\mu}} = 10^{14} \sim 10^{15},$$

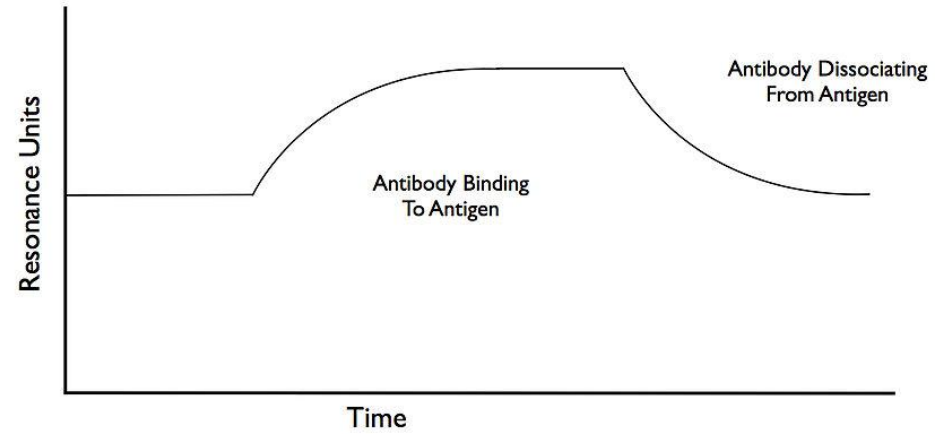
$$\vec{p} = q(\vec{r}_+ - \vec{r}_-)$$

$$p(t) = -ex(t)$$





Graph of Reference Channel 2-1



# A

