



# Nonlinear dynamics of structures and mechanical systems

Prof. Carlos Eduardo Nigro Mazzilli  
Universidade de São Paulo

# Lesson 1.2



# Lagrangian formulation for discrete systems or discretized continua

Holonomic constraints

$$\mathbf{R} = \mathbf{R}(q_1, q_2, \dots, q_n, t) \quad \Rightarrow \quad \dot{\mathbf{R}} = \frac{\partial \mathbf{R}}{\partial q_i} \dot{q}_i + \frac{\partial \mathbf{R}}{\partial t}$$

Kinetic energy

$$T = \frac{1}{2} \int_{\Omega} \rho (\dot{\mathbf{R}} \cdot \dot{\mathbf{R}}) d\Omega = \frac{1}{2} A^{ij} \dot{q}_i \dot{q}_j + B^i \dot{q}_i + \frac{1}{2} C$$

$$A^{ij} (q_1, q_2, \dots, q_n, t) = \int_{\Omega} \rho \left( \frac{\partial \mathbf{R}}{\partial q_i} \cdot \frac{\partial \mathbf{R}}{\partial q_j} \right) d\Omega$$

$$B^i (q_1, q_2, \dots, q_n, t) = \int_{\Omega} \rho \left( \frac{\partial \mathbf{R}}{\partial q_i} \cdot \frac{\partial \mathbf{R}}{\partial t} \right) d\Omega$$

$$C (q_1, q_2, \dots, q_n, t) = \int_{\Omega} \rho \left( \frac{\partial \mathbf{R}}{\partial t} \cdot \frac{\partial \mathbf{R}}{\partial t} \right) d\Omega$$

Total potential energy

$$V(q_1, q_2, \dots, q_n, t)$$

# Lagrangian formulation...

Lagrange's equation

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = N_i$$

$$L = T(q_1, q_2, \dots, q_n, \dot{q}_1, \dot{q}_2, \dots, \dot{q}_n, t) - V(q_1, q_2, \dots, q_n, t)$$



$$\begin{aligned} & A^{ij} \ddot{q}_j + \left( \frac{\partial A^{ij}}{\partial q_k} - \frac{1}{2} \frac{\partial A^{jk}}{\partial q_i} \right) \dot{q}_j \dot{q}_k + \frac{\partial A^{ij}}{\partial t} \dot{q}_j \\ & + \left( \frac{\partial B^i}{\partial q_j} - \frac{\partial B^j}{\partial q_i} \right) \dot{q}_j + \frac{\partial B^i}{\partial t} - \frac{1}{2} \frac{\partial C}{\partial q_i} + \frac{\partial V}{\partial q_i} = N_i \end{aligned}$$

# Lagrangian formulation...

Gyroscopic force (Coriolis)

$$\left( \frac{\partial B^i}{\partial q_j} - \frac{\partial B^j}{\partial q_i} \right) \dot{q}_j$$


Gyroscopic “damping” matrix is anti-symmetric

Scleronomic systems

$$\mathbf{R} = \mathbf{R}(q_1, q_2, \dots, q_n)$$

$$\frac{\partial A^{ij}}{\partial t} = 0; \quad B^i = 0; \quad C = 0 \quad \Rightarrow \quad A^{ij} \ddot{q}_j + \left( \frac{\partial A^{ij}}{\partial q_k} - \frac{1}{2} \frac{\partial A^{jk}}{\partial q_i} \right) \dot{q}_j \dot{q}_k + \frac{\partial V}{\partial q_i} = N_i$$

If the mass matrix is constant...

$$\frac{\partial A^{ij}}{\partial q_k} = 0 \quad \Rightarrow \quad A^{ij} \ddot{q}_j + \frac{\partial V}{\partial q_i} = N_i$$

# Hamiltonian formulation...

Legendre's transformation

$$f(x, y) \Rightarrow g(u, y) = ux - f(x, y)$$

$$\text{with } u = \frac{\partial f}{\partial x}$$

Observe that  $df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy = u dx + \frac{\partial f}{\partial y} dy$

$$dg = \frac{\partial g}{\partial u} du + \frac{\partial g}{\partial y} dy$$

$$= \underbrace{x du + u dx}_{0} - \frac{\partial f}{\partial x} dx - \frac{\partial f}{\partial y} dy$$

$$= x du - \frac{\partial f}{\partial y} dy \Rightarrow x = \frac{\partial g}{\partial u} \quad \text{e} \quad \frac{\partial g}{\partial y} = -\frac{\partial f}{\partial y}$$

# Hamiltonian formulation...

Duality

“Old”

“New”

$$f(x, y) = ux - g(u, y) \quad g(u, y) = ux - f(x, y)$$

$$u = \frac{\partial f}{\partial x} \quad x = \frac{\partial g}{\partial u}$$

Application to the Lagrangian

$$L(\mathbf{q}, \dot{\mathbf{q}}, t) \quad \rightarrow \quad H(\mathbf{q}, \mathbf{p}, t) = \mathbf{p}^T \dot{\mathbf{q}} - L(\mathbf{q}, \dot{\mathbf{q}}, t)$$

$$\dot{q}_i = \frac{\partial H}{\partial p_i} \quad \leftarrow \quad \text{with} \quad p_i = \frac{\partial L}{\partial \dot{q}_i}$$

# Hamiltonian formulation...

$$H(\mathbf{q}, \mathbf{p}, t) = \mathbf{p}^T \dot{\mathbf{q}} - L(\mathbf{q}, \dot{\mathbf{q}}, t)$$

$$\begin{aligned} dH &= \frac{\partial H}{\partial q_i} dq_i + \frac{\partial H}{\partial p_i} dp_i + \frac{\partial H}{\partial t} dt = \\ &= \dot{q}_i dp_i + p_i d\dot{q}_i - \underbrace{\frac{\partial L}{\partial \dot{q}_i} d\dot{q}_i}_{0} - \frac{\partial L}{\partial q_i} dq_i - \frac{\partial L}{\partial t} dt \end{aligned}$$



$$\left\{ \begin{array}{l} \dot{q}_i = \frac{\partial H}{\partial p_i} \\ \frac{\partial H}{\partial q_i} = -\frac{\partial L}{\partial q_i} \\ \frac{\partial H}{\partial t} = -\frac{\partial L}{\partial t} \end{array} \right.$$

# Hamiltonian formulation...

Lagrange's equation with

$$N_i = 0$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$$



$$\dot{p}_i = \frac{\partial L}{\partial q_i} = - \frac{\partial H}{\partial q_i}$$

$$\dot{q}_i = \frac{\partial H}{\partial p_i}$$

First-order system

Hamilton's canonical equations

$$\dot{H} = \frac{dH}{dt} = \frac{\partial H}{\partial q_i} \dot{q}_i + \frac{\partial H}{\partial p_i} \dot{p}_i + \frac{\partial H}{\partial t} = \frac{\partial H}{\partial q_i} \frac{\partial H}{\partial p_i} - \frac{\partial H}{\partial p_i} \frac{\partial H}{\partial q_i} + \frac{\partial H}{\partial t} = \frac{\partial H}{\partial t}$$



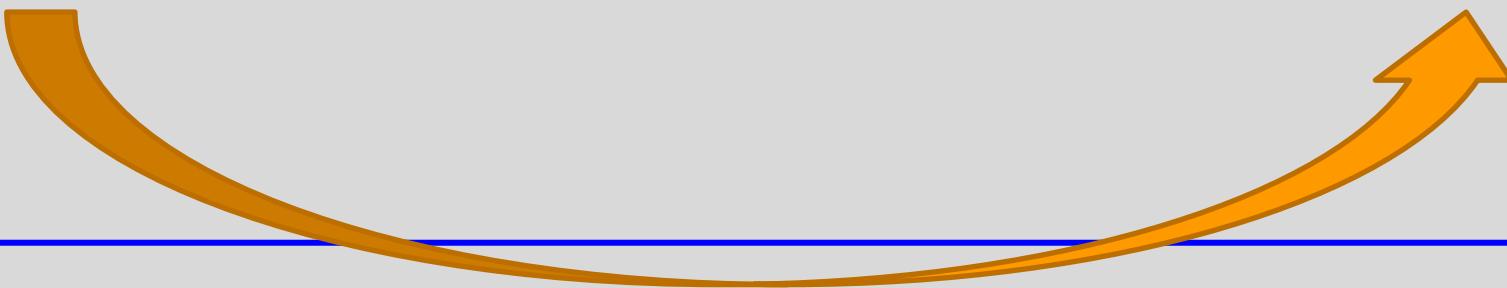
# Hamiltonian formulation...

Lagrange's equation with  $N_i \neq 0$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = N_i \quad \Rightarrow \quad \dot{p}_i = N_i + \frac{\partial L}{\partial q_i} = N_i - \frac{\partial H}{\partial q_i}$$
$$\dot{q}_i = \frac{\partial H}{\partial p_i}$$

Hamilton's canonical equations

$$\dot{H} = \frac{dH}{dt} = \frac{\partial H}{\partial q_i} \dot{q}_i + \frac{\partial H}{\partial p_i} \dot{p}_i + \frac{\partial H}{\partial t} = \frac{\partial H}{\partial q_i} \frac{\partial H}{\partial p_i} + \frac{\partial H}{\partial p_i} \left( N_i - \frac{\partial H}{\partial q_i} \right) + \frac{\partial H}{\partial t} = N_i \frac{\partial H}{\partial p_i} + \frac{\partial H}{\partial t}$$



# Hamiltonian formulation...

Scleronomous conservative system

$$\mathbf{R} = \mathbf{R}(\mathbf{q}) \quad \text{e} \quad N_i = 0$$

$$H = H(\mathbf{p}, \mathbf{q}) \therefore \frac{\partial H}{\partial t} = 0 \Rightarrow \dot{H} = 0 \Rightarrow H = \text{const.}$$

$$L = T - V = \frac{1}{2} A^{ij} \dot{q}_i \dot{q}_j - V(\mathbf{q}) \Rightarrow \frac{\partial L}{\partial \dot{q}_j} = A^{ij} \dot{q}_i$$


$$H(\mathbf{q}, \mathbf{p}) = \mathbf{p}^T \dot{\mathbf{q}} - L(\mathbf{q}, \dot{\mathbf{q}}) = p_j \dot{q}_j - L = \frac{\partial L}{\partial \dot{q}_j} \dot{q}_j - L$$

$$= A^{ij} \dot{q}_i \dot{q}_j - \frac{1}{2} A^{ij} \dot{q}_i \dot{q}_j + V = \frac{1}{2} A^{ij} \dot{q}_i \dot{q}_j + V$$
$$= T + V$$

# Hamiltonian formulation...

Rheonomic system  $\mathbf{R} = \mathbf{R}(\mathbf{q}, t)$  with  $N_i = 0$  or  $N_i \neq 0$

$$H = H(\mathbf{p}, \mathbf{q}, t) \text{ and } \dot{H} = N_i \frac{\partial H}{\partial p_i} + \frac{\partial H}{\partial t} \Rightarrow H \neq \text{const.}$$

$$L = T - V = \frac{1}{2} A^{ij} \dot{q}_i \dot{q}_j + B^i \dot{q}_i + \frac{1}{2} C - V(\mathbf{q}) \Rightarrow \frac{\partial L}{\partial \dot{q}_j} = A^{ij} \dot{q}_i + B^j$$


$$\begin{aligned} H(\mathbf{q}, \mathbf{p}) &= \mathbf{p}^T \dot{\mathbf{q}} - L(\mathbf{q}, \dot{\mathbf{q}}) = p_j \dot{q}_j - L = \frac{\partial L}{\partial \dot{q}_j} \dot{q}_j - T + V \\ &= (A^{ij} \dot{q}_i \dot{q}_j + B^j \dot{q}_j) - T + V = (2T - C - B^j \dot{q}_j) - T + V \\ &= T + V - C - B^j \dot{q}_j \end{aligned}$$

mechanic energy      pseudo-potential energy      Coriolis 'energy'