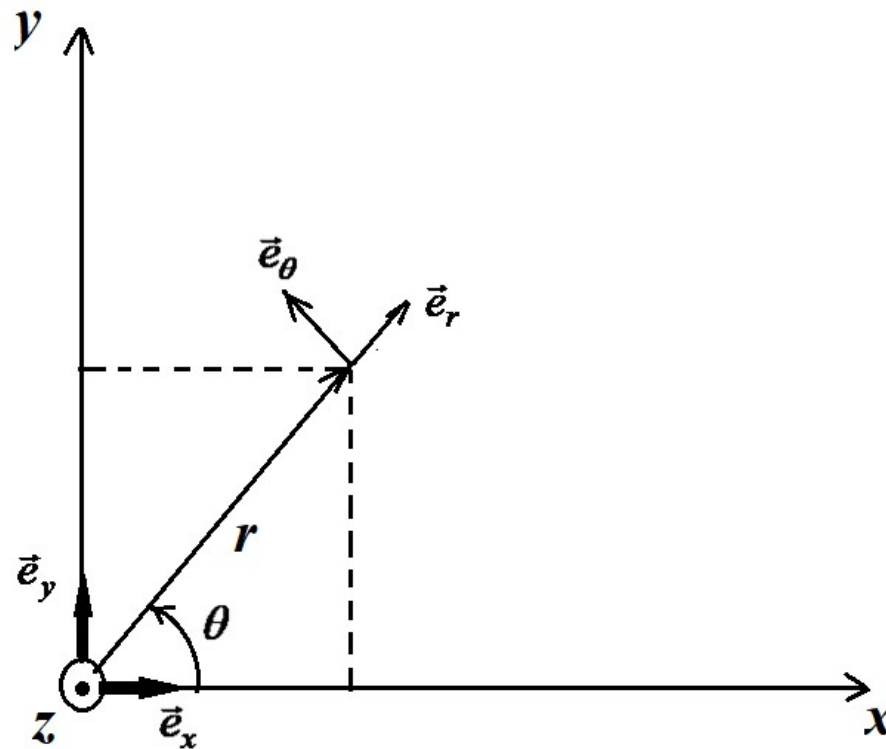


Coordenadas Cilíndricas

Aceleração, Divergente, Rotacional em coordenadas
cilíndricas

1) Relação entre os sistemas de coordenadas cartesiano e cilíndrico



$$r \vec{e}_r = x \vec{e}_x + y \vec{e}_y \Rightarrow \vec{e}_r = \frac{x}{r} \vec{e}_x + \frac{y}{r} \vec{e}_y \Rightarrow \boxed{\vec{e}_r = \cos \theta \vec{e}_x + \sin \theta \vec{e}_y}$$

$$\vec{e}_\theta = \vec{e}_z \times \vec{e}_r \Rightarrow \boxed{\vec{e}_\theta = -\sin \theta \vec{e}_x + \cos \theta \vec{e}_y}$$

Das expressões dos versores $\vec{e}_r$ e $\vec{e}_\theta$:

$$\frac{d\vec{e}_r}{d\theta} = -\sin\theta \vec{e}_x + \cos\theta \vec{e}_y = \vec{e}_\theta$$

$$\frac{d\vec{e}_\theta}{d\theta} = -\cos\theta \vec{e}_x - \sin\theta \vec{e}_y = -\vec{e}_r$$

O sistema de coordenadas cilíndrico é um sistema de coordenadas não-inercial, e devido à dependência dos versores $\vec{e}_r$ e $\vec{e}_\theta$ com a coordenada θ temos o aparecimento de termos adicionais nas equações da aceleração, do divergente de um vetor e do rotacional de um vetor.

2) Aceleração em coordenadas cilíndricas

A expressão de aceleração é:

$$\vec{a} = \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v}$$

Num sistema cilíndrico de coordenadas, o vetor da velocidade é dado por:

$$\vec{v} = v_r \vec{e}_r + v_\theta \vec{e}_\theta + v_z \vec{e}_z$$

E o vetor nabla (operador gradiente):

$$\nabla = \vec{e}_r \frac{\partial}{\partial r} + \vec{e}_\theta \frac{1}{r} \frac{\partial}{\partial \theta} + \vec{e}_z \frac{\partial}{\partial z}$$

A expressão da aceleração resulta:

$$\vec{a} = \frac{\partial}{\partial t} (\nu_r \vec{e}_r + \nu_\theta \vec{e}_\theta + \nu_z \vec{e}_z) + \left(\nu_r \frac{\partial}{\partial r} + \nu_\theta \frac{\partial}{r \partial \theta} + \nu_z \frac{\partial}{\partial z} \right) (\nu_r \vec{e}_r + \nu_\theta \vec{e}_\theta + \nu_z \vec{e}_z)$$

Ou seja:

$$\begin{aligned} \vec{a} = & \frac{\partial(\nu_r \vec{e}_r)}{\partial t} + \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial t} + \frac{\partial(\nu_z \vec{e}_z)}{\partial t} + \nu_r \frac{\partial(\nu_r \vec{e}_r)}{\partial r} + \nu_\theta \frac{\partial(\nu_r \vec{e}_r)}{r \partial \theta} + \nu_z \frac{\partial(\nu_r \vec{e}_r)}{\partial z} \\ & + \nu_r \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial r} + \nu_\theta \frac{\partial(\nu_\theta \vec{e}_\theta)}{r \partial \theta} + \nu_z \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial z} + \nu_r \frac{\partial(\nu_z \vec{e}_z)}{\partial r} + \nu_\theta \frac{\partial(\nu_z \vec{e}_z)}{r \partial \theta} + \nu_z \frac{\partial(\nu_z \vec{e}_z)}{\partial z} \end{aligned}$$

Continuando:

$$\begin{aligned}
 \vec{a} = & \frac{\partial v_r}{\partial t} \vec{e}_r + \frac{\partial v_\theta}{\partial t} \vec{e}_\theta + \frac{\partial v_z}{\partial t} \vec{e}_z + v_r \frac{\partial v_r}{\partial r} \vec{e}_r + v_\theta \frac{\partial v_r}{r \partial \theta} \vec{e}_r + \frac{v_\theta v_r}{r} \underbrace{\frac{\partial \vec{e}_r}{\partial \theta}}_{\vec{e}_\theta} + v_z \frac{\partial v_r}{\partial z} \vec{e}_r \\
 & + v_r \frac{\partial v_\theta}{\partial r} \vec{e}_\theta + v_\theta \frac{\partial v_\theta}{r \partial \theta} \vec{e}_\theta + \frac{v_\theta^2}{r} \underbrace{\frac{\partial \vec{e}_\theta}{\partial \theta}}_{-\vec{e}_r} + v_z \frac{\partial v_\theta}{\partial z} \vec{e}_\theta \\
 & + v_r \frac{\partial v_z}{\partial r} \vec{e}_z + v_\theta \frac{\partial v_z}{r \partial \theta} \vec{e}_z + v_z \frac{\partial v_z}{\partial z} \vec{e}_z
 \end{aligned}$$

Resultam as componentes:

$$a_r = \frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + v_\theta \frac{\partial v_r}{r \partial \theta} + v_z \frac{\partial v_r}{\partial z} - \frac{v_\theta^2}{r}$$

$$a_\theta = \frac{\partial v_\theta}{\partial t} + v_r \frac{\partial v_\theta}{\partial r} + v_\theta \frac{\partial v_\theta}{r \partial \theta} + v_z \frac{\partial v_\theta}{\partial z} + \frac{v_r v_\theta}{r}$$

$$a_z = \frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + v_\theta \frac{\partial v_z}{r \partial \theta} + v_z \frac{\partial v_z}{\partial z}$$

3) Divergente de um vetor

O divergente da velocidade é dado por:

$$\nabla \cdot \vec{v} = \left(\vec{e}_r \frac{\partial}{\partial r} + \vec{e}_\theta \frac{1}{r} \frac{\partial}{\partial \theta} + \vec{e}_z \frac{\partial}{\partial z} \right) \cdot (\nu_r \vec{e}_r + \nu_\theta \vec{e}_\theta + \nu_z \vec{e}_z)$$

Isso resulta:

$$\begin{aligned} \nabla \cdot \vec{v} = & \vec{e}_r \cdot \frac{\partial(\nu_r \vec{e}_r)}{\partial r} + \vec{e}_\theta \cdot \frac{\partial(\nu_r \vec{e}_r)}{r \partial \theta} + \vec{e}_z \cdot \frac{\partial(\nu_r \vec{e}_r)}{\partial z} \\ & + \vec{e}_r \cdot \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial r} + \vec{e}_\theta \cdot \frac{\partial(\nu_\theta \vec{e}_\theta)}{r \partial \theta} + \vec{e}_z \cdot \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial z} \\ & + \vec{e}_r \cdot \frac{\partial(\nu_z \vec{e}_z)}{\partial r} + \vec{e}_\theta \cdot \frac{\partial(\nu_z \vec{e}_z)}{r \partial \theta} + \vec{e}_z \cdot \frac{\partial(\nu_z \vec{e}_z)}{\partial z} \end{aligned}$$

Continuando:

$$\begin{aligned}
 \nabla \cdot \vec{v} = & \vec{e}_r \cdot \frac{\partial v_r}{\partial r} \vec{e}_r + \vec{e}_\theta \cdot \frac{\partial v_r}{r \partial \theta} \vec{e}_r + \underbrace{\vec{e}_\theta \cdot \frac{v_r}{r} \frac{\partial \vec{e}_r}{\partial \theta}}_{\vec{e}_\theta} + \vec{e}_z \cdot \frac{\partial v_r}{\partial z} \vec{e}_r \\
 & + \vec{e}_r \cdot \frac{\partial v_\theta}{\partial r} \vec{e}_\theta + \vec{e}_\theta \cdot \frac{\partial v_\theta}{r \partial \theta} \vec{e}_\theta + \underbrace{\vec{e}_\theta \cdot \frac{v_\theta}{r} \frac{\partial \vec{e}_\theta}{\partial \theta}}_{-\vec{e}_r} + \vec{e}_z \cdot \frac{\partial v_\theta}{\partial z} \vec{e}_\theta \\
 & + \vec{e}_r \cdot \frac{\partial v_z}{\partial r} \vec{e}_z + \vec{e}_\theta \cdot \frac{\partial v_z}{r \partial \theta} \vec{e}_z + \vec{e}_z \cdot \frac{\partial v_z}{\partial z} \vec{e}_z
 \end{aligned}$$

Finalmente:

$$\nabla \cdot \vec{v} = \frac{\partial v_r}{\partial r} + \frac{v_r}{r} + \frac{\partial v_\theta}{r \partial \theta} + \frac{\partial v_z}{\partial z}$$

Em geral, os dois primeiros termos são combinados usando a regra da cadeia:

$$\nabla \cdot \vec{v} = \frac{1}{r} \frac{\partial (r v_r)}{\partial r} + \frac{\partial v_\theta}{r \partial \theta} + \frac{\partial v_z}{\partial z}$$

Assim, para um escoamento incompressível:

$$\boxed{\nabla \cdot \vec{v} = \frac{1}{r} \frac{\partial (r v_r)}{\partial r} + \frac{\partial v_\theta}{r \partial \theta} + \frac{\partial v_z}{\partial z} = 0}$$

4) Rotacional de um vetor

O rotacional do vetor da velocidade é:

$$\nabla \times \vec{v} = \left(\vec{e}_r \frac{\partial}{\partial r} + \vec{e}_\theta \frac{1}{r} \frac{\partial}{\partial \theta} + \vec{e}_z \frac{\partial}{\partial z} \right) \times (\nu_r \vec{e}_r + \nu_\theta \vec{e}_\theta + \nu_z \vec{e}_z)$$

E podemos escrever:

$$\begin{aligned} \nabla \times \vec{v} = & \vec{e}_r \times \frac{\partial(\nu_r \vec{e}_r)}{\partial r} + \vec{e}_\theta \times \frac{\partial(\nu_r \vec{e}_r)}{r \partial \theta} + \vec{e}_z \times \frac{\partial(\nu_r \vec{e}_r)}{\partial z} \\ & + \vec{e}_r \times \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial r} + \vec{e}_\theta \times \frac{\partial(\nu_\theta \vec{e}_\theta)}{r \partial \theta} + \vec{e}_z \times \frac{\partial(\nu_\theta \vec{e}_\theta)}{\partial z} \\ & + \vec{e}_r \times \frac{\partial(\nu_z \vec{e}_z)}{\partial r} + \vec{e}_\theta \times \frac{\partial(\nu_z \vec{e}_z)}{r \partial \theta} + \vec{e}_z \times \frac{\partial(\nu_z \vec{e}_z)}{\partial z} \end{aligned}$$

Continuando:

$$\begin{aligned}
 \nabla \times \vec{v} = & \vec{e}_r \times \frac{\partial v_r}{\partial r} \vec{e}_r + \vec{e}_\theta \times \frac{\partial v_r}{r \partial \theta} \vec{e}_r + \underbrace{\vec{e}_\theta \times \frac{v_r}{r} \frac{\partial \vec{e}_r}{\partial \theta}}_{\vec{e}_\theta} + \vec{e}_z \times \frac{\partial v_r}{\partial z} \vec{e}_r \\
 & + \vec{e}_r \times \frac{\partial v_\theta}{\partial r} \vec{e}_\theta + \vec{e}_\theta \times \frac{\partial v_\theta}{r \partial \theta} \vec{e}_\theta + \underbrace{\vec{e}_\theta \times \frac{v_\theta}{r} \frac{\partial \vec{e}_\theta}{\partial \theta}}_{-\vec{e}_r} + \vec{e}_z \times \frac{\partial v_\theta}{\partial z} \vec{e}_\theta \\
 & + \vec{e}_r \times \frac{\partial v_z}{\partial r} \vec{e}_z + \vec{e}_\theta \times \frac{\partial v_z}{r \partial \theta} \vec{e}_z + \vec{e}_z \times \frac{\partial v_z}{\partial z} \vec{e}_z
 \end{aligned}$$

Isso resulta:

$$\nabla \times \vec{v} = -\frac{\partial v_r}{r \partial \theta} \vec{e}_z + \frac{\partial v_r}{\partial z} \vec{e}_\theta + \frac{\partial v_\theta}{\partial r} \vec{e}_z + \frac{v_\theta}{r} \vec{e}_z - \frac{\partial v_\theta}{\partial z} \vec{e}_r - \frac{\partial v_z}{\partial r} \vec{e}_\theta + \frac{\partial v_z}{r \partial \theta} \vec{e}_r$$

Reorganizando os termos:

$$\nabla \times \vec{v} = \left(\frac{\partial v_z}{r \partial \theta} - \frac{\partial v_\theta}{\partial z} \right) \vec{e}_r + \left(\frac{\partial v_r}{\partial z} - \frac{\partial v_z}{\partial r} \right) \vec{e}_\theta + \left(\frac{\partial v_\theta}{\partial r} - \frac{\partial v_r}{r \partial \theta} + \frac{v_\theta}{r} \right) \vec{e}_z$$

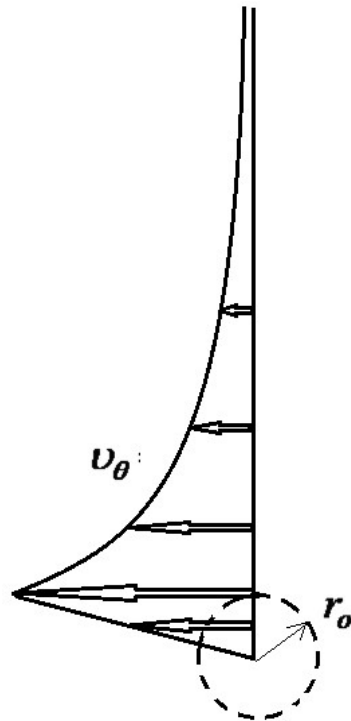
O rotacional da velocidade é chamado de vorticidade $\vec{\omega}$. Na forma de determinante:

$$\vec{\omega} = \nabla \times \vec{v} = \begin{vmatrix} \vec{e}_r & \vec{e}_\theta & \vec{e}_z \\ \frac{\partial}{\partial r} & \frac{\partial}{r \partial \theta} & \frac{\partial}{\partial z} \\ v_r & v_\theta & v_z \end{vmatrix} + \frac{v_\theta}{r} \vec{e}_z$$

Exercício: O vórtice de Rankine é um modelo usado para representar furacões ou tornados. Ele é dado pelo campo de velocidades:

$$v_r \cong 0 \quad v_z \cong 0 \quad v_\theta = \begin{cases} \Omega r & \text{para } r \leq r_o \\ \frac{\Omega r_o^2}{r} & \text{para } r > r_o \end{cases}$$

Calcule as acelerações e verifique se o escoamento é incompressível e irrotacional.



Solução:

$$a_r = \frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + v_\theta \frac{\partial v_r}{r \partial \theta} + v_z \frac{\partial v_r}{\partial z} - \frac{v_\theta^2}{r}$$

$$a_r = -\frac{v_\theta^2}{r}$$

Logo:

$$a_r = -\Omega^2 r \quad \text{para } r \leq r_o \quad \text{e} \quad a_r = -\frac{\Omega^2 r_o^4}{r^3} \quad \text{para } r > r_o$$

$$a_{\theta} = \frac{\partial v_{\theta}}{\partial t} + v_r \frac{\partial v_{\theta}}{\partial r} + v_{\theta} \frac{\partial v_{\theta}}{r \partial \theta} + v_z \frac{\partial v_{\theta}}{\partial z} + \frac{v_r v_{\theta}}{r} = 0$$

$$a_z = \frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + v_{\theta} \frac{\partial v_z}{r \partial \theta} + v_z \frac{\partial v_z}{\partial z} = 0$$

$$\nabla \cdot \vec{v} = \frac{1}{r} \frac{\partial (r v_r)}{\partial r} + \frac{\partial v_{\theta}}{r \partial \theta} + \frac{\partial v_z}{\partial z} = 0 \quad \text{para qualquer } r.$$

Assim, o escoamento é incompressível.

Para $r \leq r_o$:

$$\vec{\omega} = \nabla \times \vec{v} = \left(\underbrace{\frac{\partial v_z}{r \partial \theta}}_0 - \underbrace{\frac{\partial v_\theta}{\partial z}}_0 \right) \vec{e}_r + \left(\underbrace{\frac{\partial v_r}{\partial z}}_0 - \underbrace{\frac{\partial v_z}{\partial r}}_0 \right) \vec{e}_\theta + \left(\underbrace{\frac{\partial v_\theta}{\partial r}}_\Omega - \underbrace{\frac{\partial v_r}{r \partial \theta}}_0 + \underbrace{\frac{v_\theta}{r}}_\Omega \right) \vec{e}_z = 2\Omega \vec{e}_z$$

Assim, para $r \leq r_o$ o escoamento é rotacional.

Para $r > r_o$:

$$\vec{\omega} = \nabla \times \vec{v} = \left(\underbrace{\frac{\partial v_z}{r \partial \theta}}_0 - \underbrace{\frac{\partial v_\theta}{\partial z}}_0 \right) \vec{e}_r + \left(\underbrace{\frac{\partial v_r}{\partial z}}_0 - \underbrace{\frac{\partial v_z}{\partial r}}_0 \right) \vec{e}_\theta + \left(\underbrace{\frac{\partial v_\theta}{\partial r}}_{-\frac{\Omega r_o^2}{r^2}} - \underbrace{\frac{\partial v_r}{r \partial \theta}}_0 + \underbrace{\frac{v_\theta}{r}}_{\frac{\Omega r_o^2}{r^2}} \right) \vec{e}_z = 0$$

Assim, para $r > r_o$ o escoamento é irrotacional.

Bibliografia:

White, F.M., “Mecânica dos Fluidos”, 5ª edição, Ed. McGraw Hill, 2010.

Potter, M.C.; Wiggert, D.C., “Mecânica dos Fluidos”, Ed. Thomson Learning, 2004.

Mase, G.T.; Mase, G.E., “Continuum Mechanics for Engineers”, third edition, CRC Press, 1999.