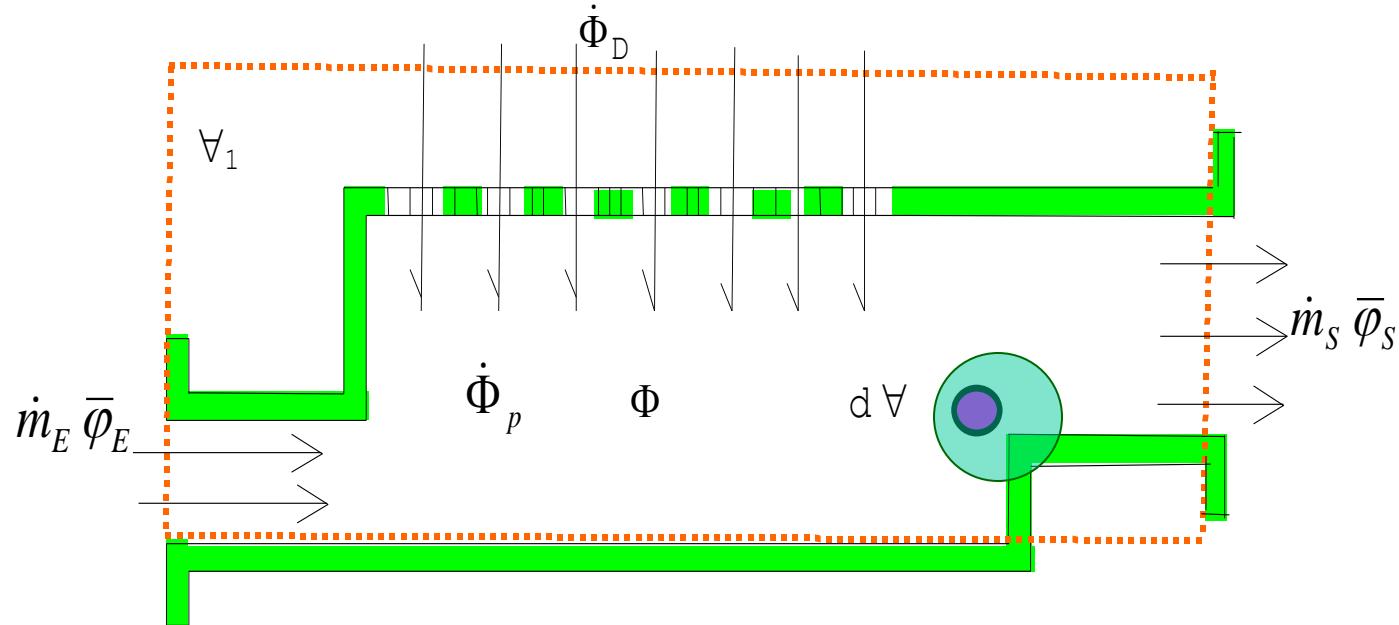


Possibilidades dos FTS

Discretização

A.G. Antunha

$$\frac{d\Phi}{dt} = \dot{m}_E \bar{\varphi}_E - \dot{m}_S \bar{\varphi}_S + \dot{\Phi}_D + \dot{\Phi}_P$$



$$\dot{m} = \bar{\rho} v_b S = \int_S \rho \vec{v} dS \quad ; \quad \dot{m} \bar{\varphi} = \bar{\rho} \varphi_b v_b S = \int_S \rho \varphi \vec{v} dS$$

$$v_b = \frac{1}{\bar{\rho} S} \int_S \rho \vec{v} dS \quad ; \quad \varphi_b = \frac{1}{\bar{\rho} v_b S} \int_S \rho \varphi \vec{v} dS$$

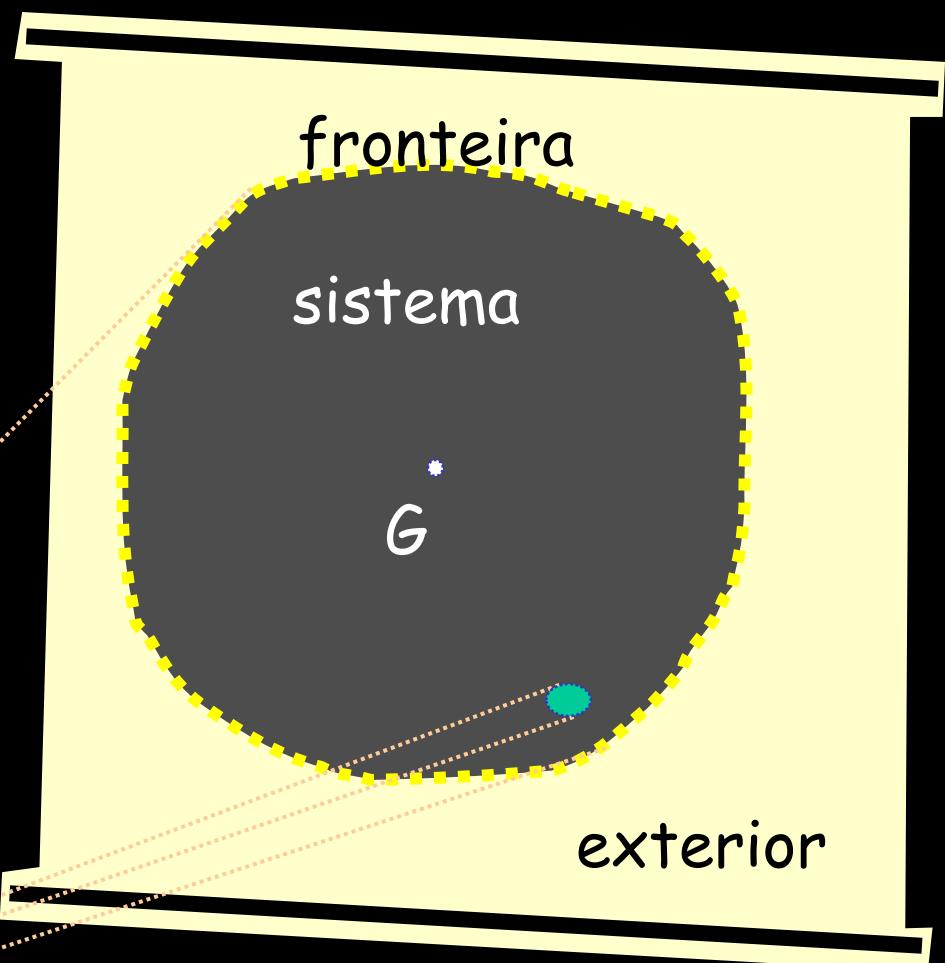
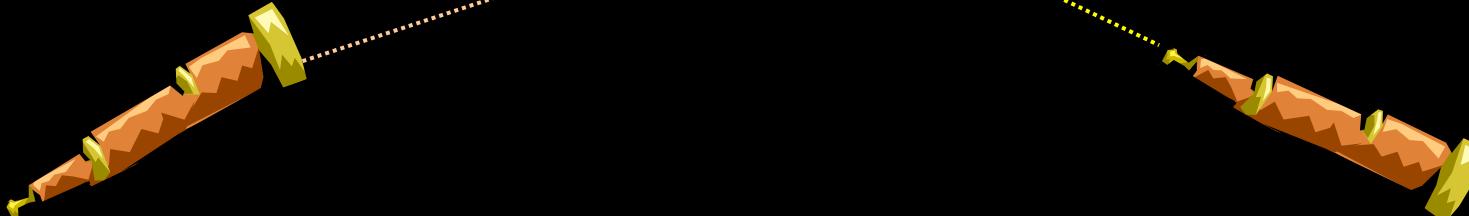
macro = ∫ micro

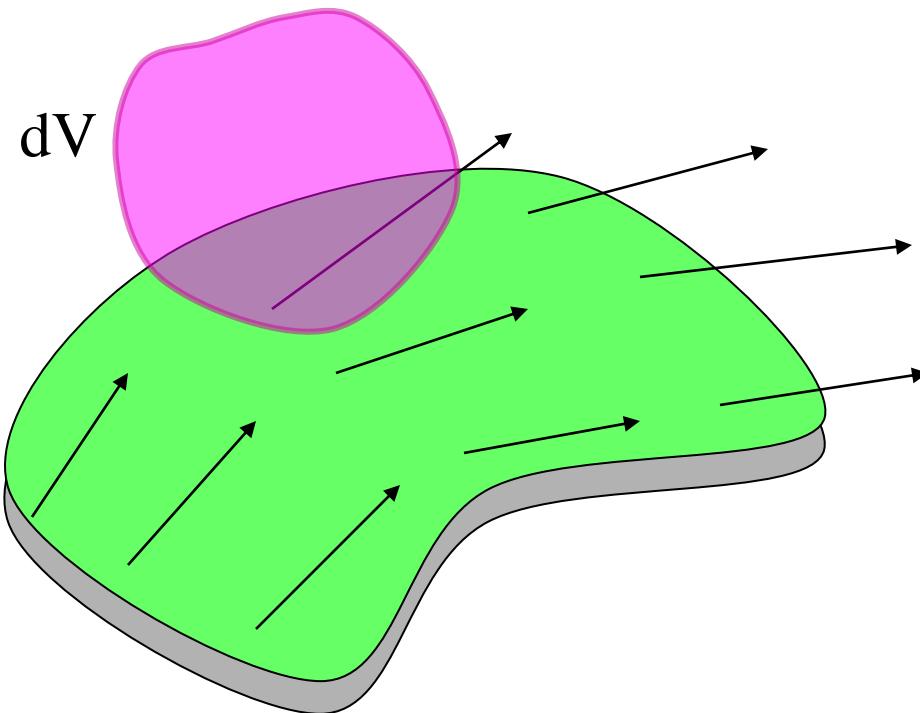
$$m = \int_{sist} dm$$

$$\forall = \int_{sist} dV$$

$$m = \int_{\forall} \rho dV$$

$$\Phi = \int_{\forall} \rho \varphi dV$$



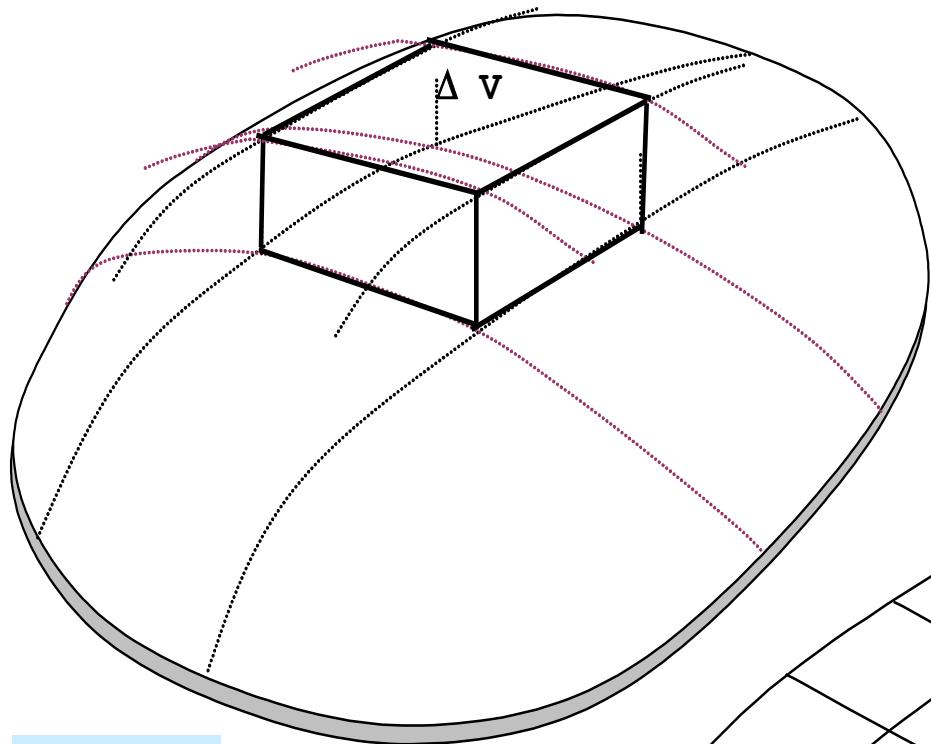


$$\rho \frac{D \phi}{Dt} = \rho \left(\frac{\partial \phi}{\partial t} + \vec{v} \cdot \vec{\text{grad}} \phi \right) = \frac{\partial \rho \phi}{\partial t} + \text{div} \rho \vec{v} \phi = \text{div} \rho \lambda_\Phi \vec{\text{grad}} \phi + \dot{\sigma}_{\forall \Phi}$$

onde $\vec{j}_\Phi = -\rho \lambda_\Phi \nabla \phi$.

sendo $\phi = 1$ a equação da continuidade:

$$\frac{\partial \rho}{\partial t} + \text{div} \rho \vec{v} = \frac{D \rho}{Dt} + \rho \text{div} \vec{v} = 0$$

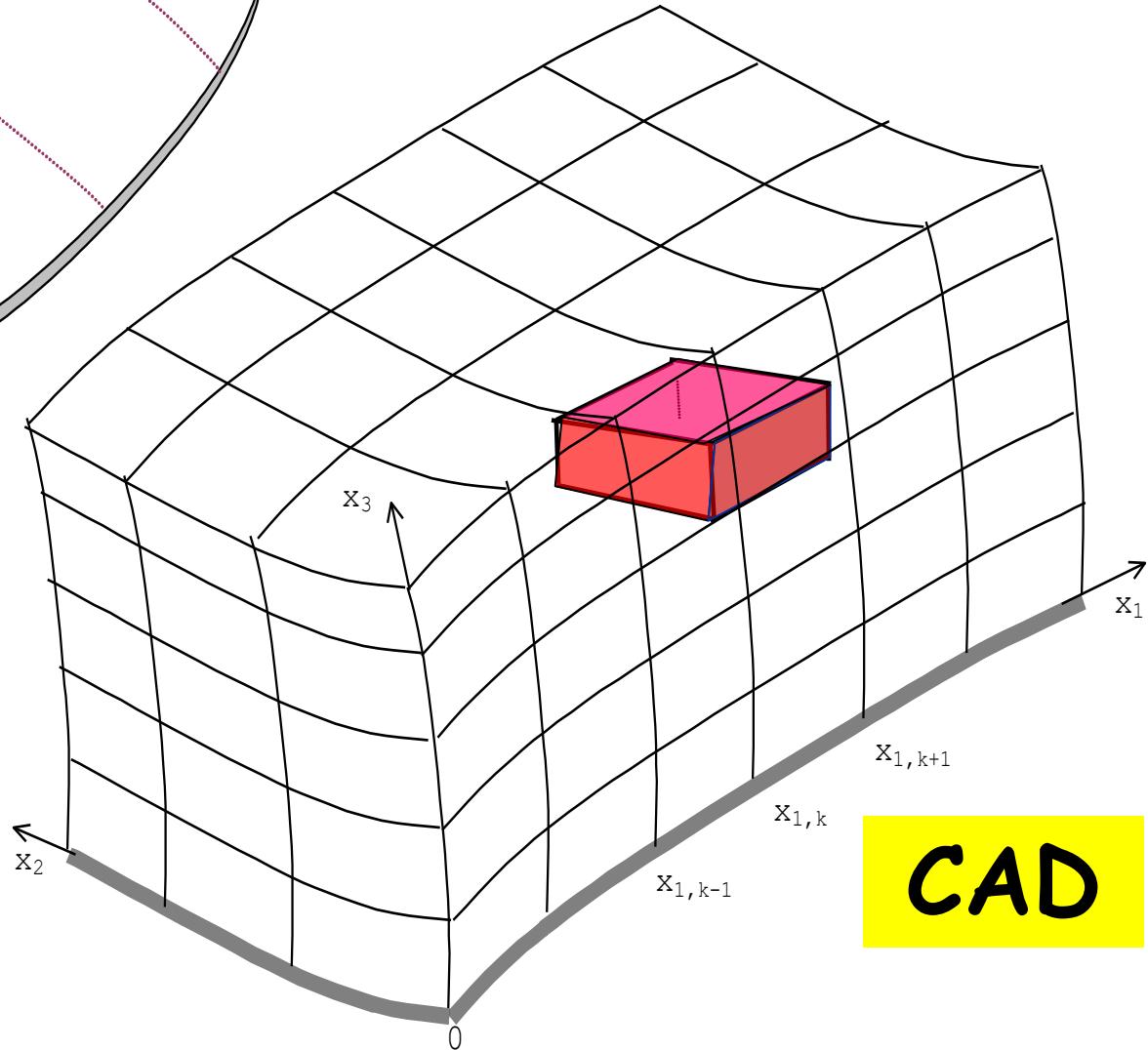


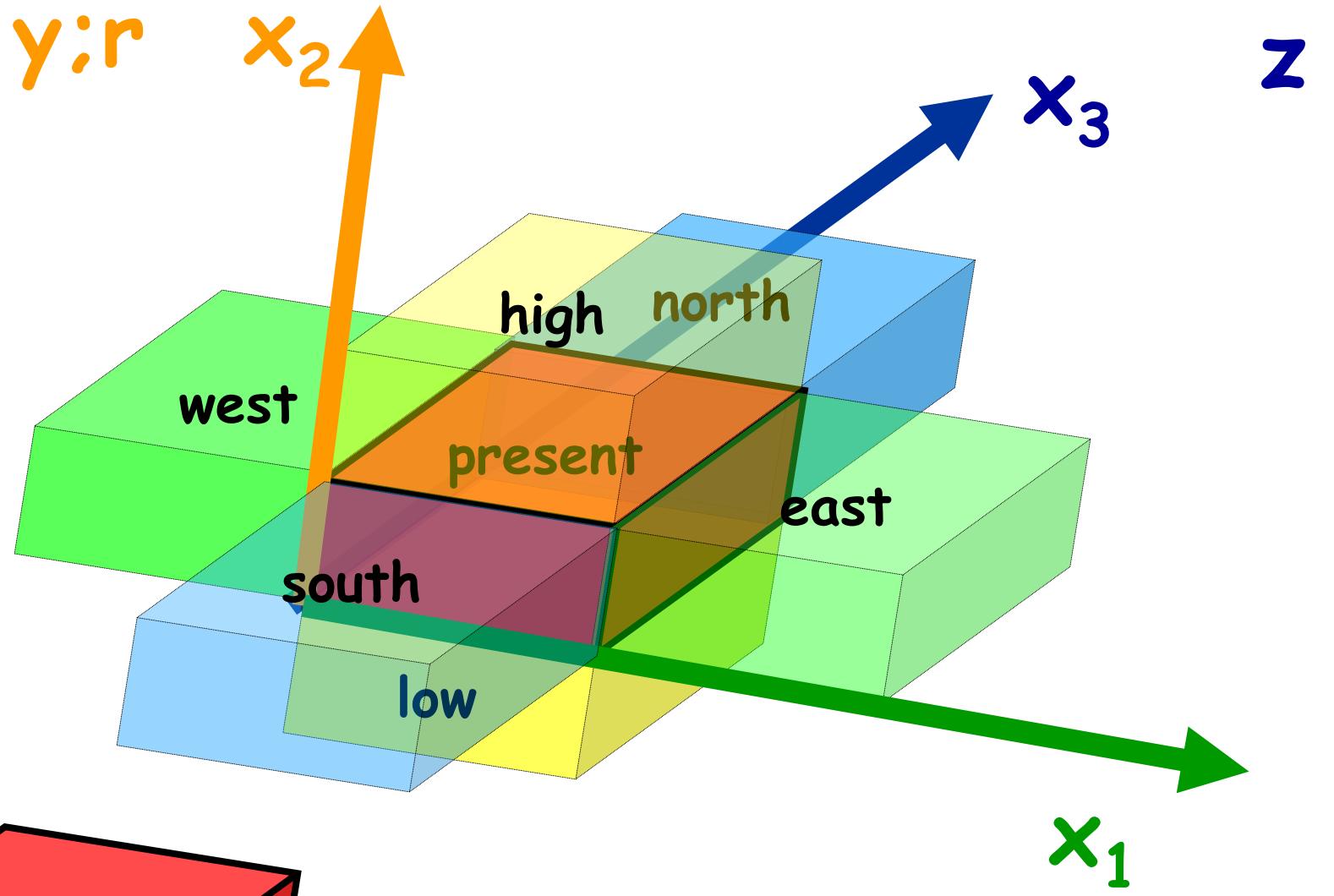
ΔV

construir a malha a
partir da fronteira

cada célula é
um batch

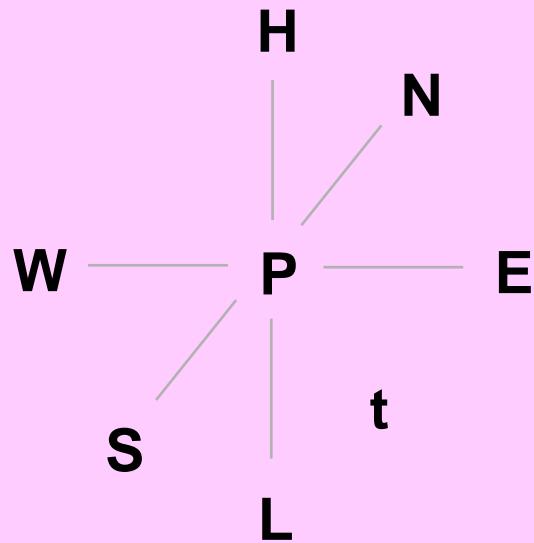
$$\dot{\sigma}_{V\Phi} = C_* (\phi_* - \phi)$$





convenções

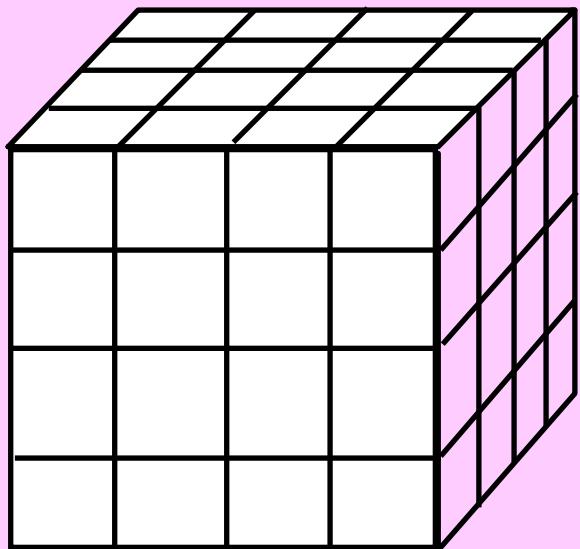
Os elementos de volume formam um conjunto ordenado de pontos no tempo e no espaço. Para o caso de elementos hexaédricos:



- P** - elemento central
- N** - elemento ao norte
- S** - elemento ao sul
- E** - elemento ao leste
- W** - elemento ao oeste
- H** - elemento acima
- L** - elemento abaixo
- t** - instante de tempo

Relação de um ponto do *grid* com suas vizinhanças:

$$A_P \phi_P = A_N \phi_N + A_S \phi_S + A_E \phi_E + A_W \phi_W + A_H \phi_H + A_L \phi_L + A_t \phi_t + \text{fontes}$$

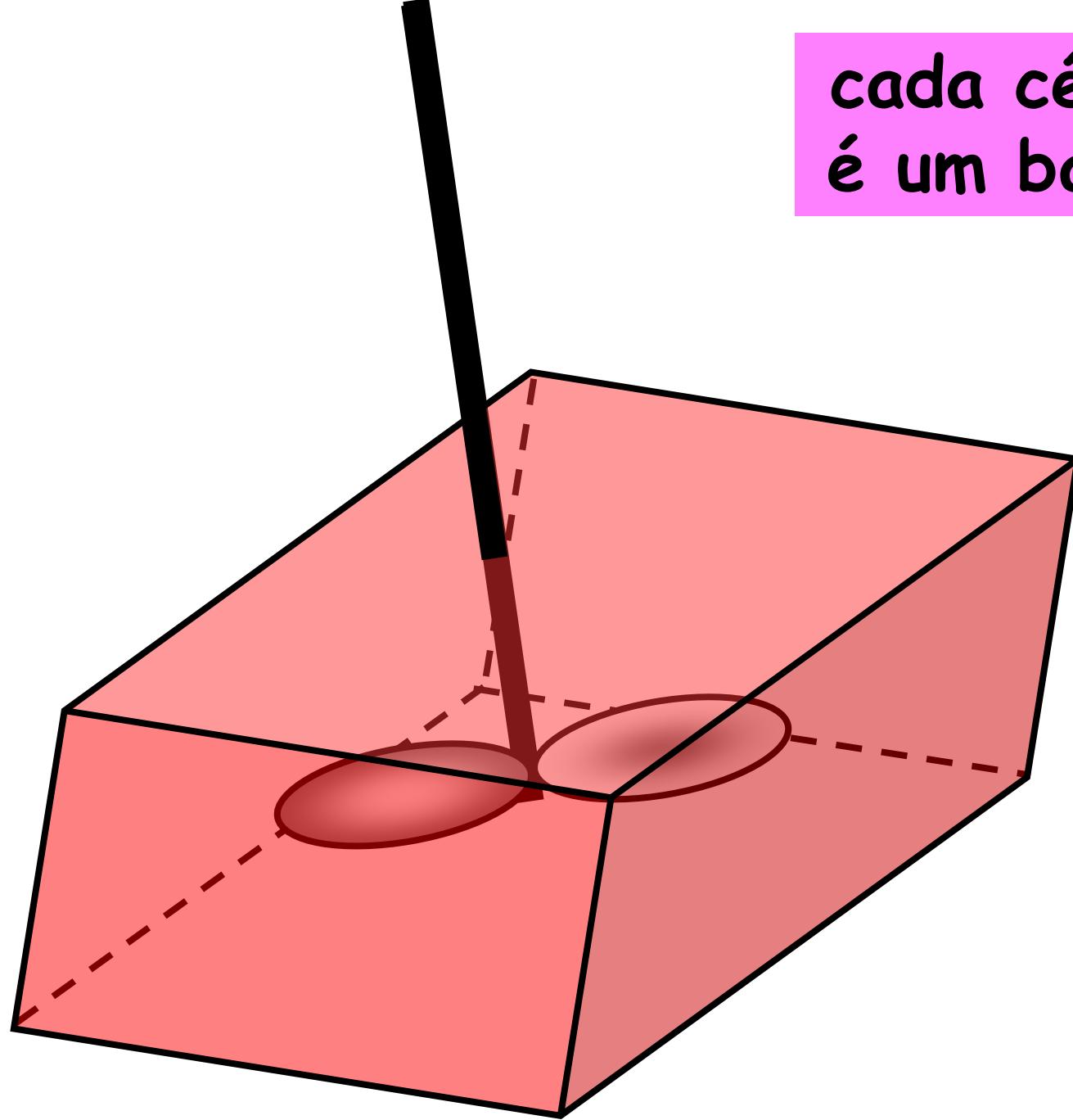


Como as equações de conservação não podem ser aplicadas na forma diferencial, elas são discretizadas em um número finito de elementos de volume, definidos de forma que preencham todo o espaço pelo qual o fluido escoa (domínio).

Quando o número de elementos de volume tende a infinito, a equação de conservação tende a forma diferencial. É, no entanto, atualmente inviável a resolução computacional de um domínio infinitesimalmente discretizado.

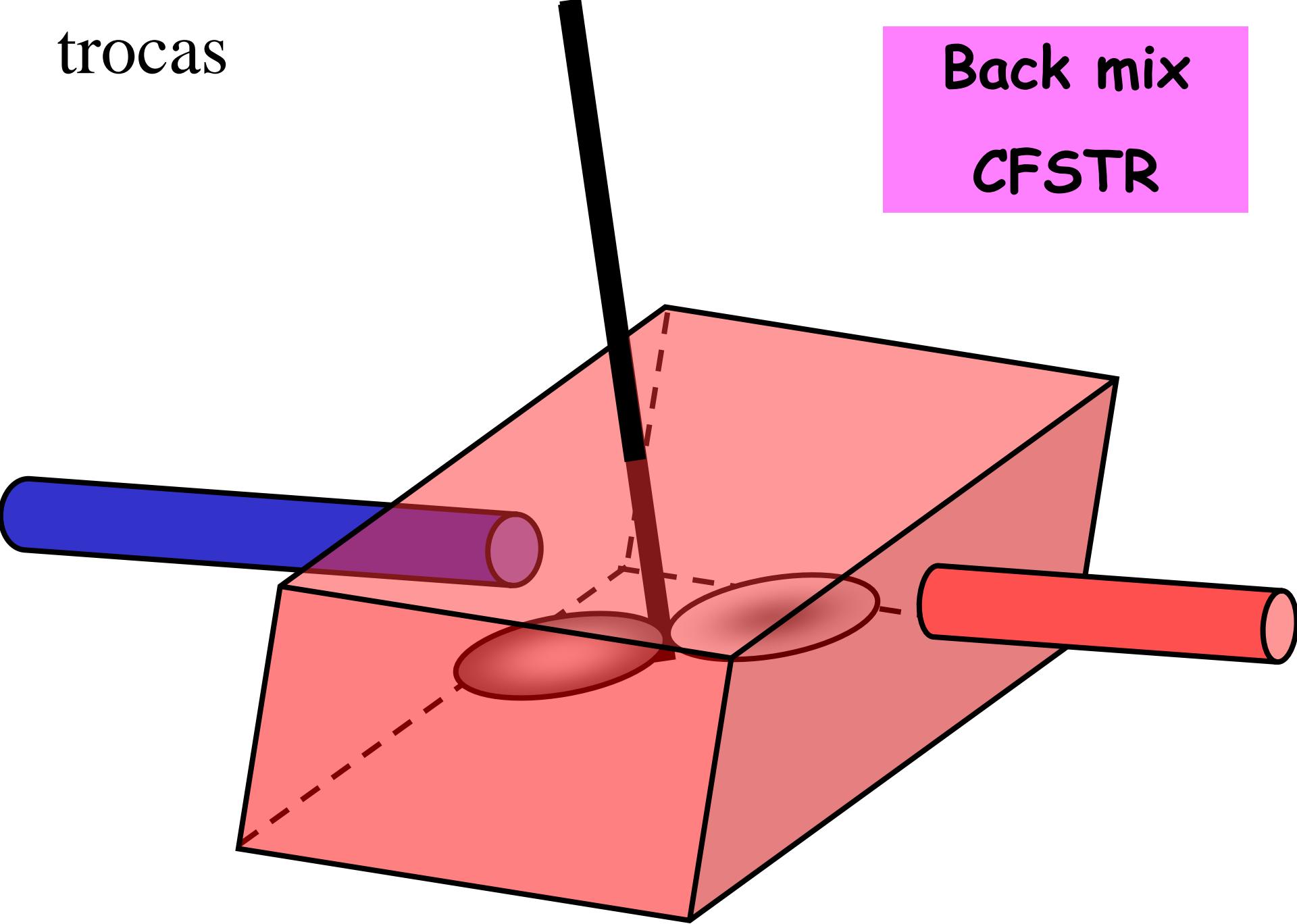
batelada

cada célula
é um batch



trocas

Back mix
CFSTR



des-divergir

$$\frac{\partial \rho \varphi}{\partial t} + \operatorname{div} \rho \vec{v} \varphi = - \operatorname{div} \vec{j}_\Phi + \dot{\sigma}_{V_\Phi}$$

$$\vec{j}_\varphi = -\rho \lambda_\Phi \operatorname{grad} \varphi$$

$$\dot{\sigma}_{V_\Phi} = c_* (\varphi_* - \varphi)$$

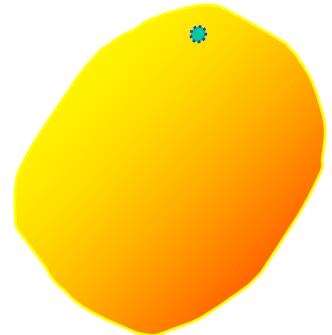
$$\frac{\partial \rho \varphi}{\partial t} = - \operatorname{div} \rho (\vec{v} \varphi - \lambda_\Phi \operatorname{grad} \varphi) + c_* (\varphi_* - \varphi)$$

A

$$\int_{\mathbb{A}} \frac{\partial \rho \varphi}{\partial t} dV = - \int_{\mathbb{A}} \operatorname{div} \rho (\vec{v} \varphi - \lambda_\Phi \operatorname{grad} \varphi) dV + \int_{\mathbb{A}} c_* (\varphi_* - \varphi) dV$$

$$\int_{\mathbb{A}} \frac{\partial \rho \varphi}{\partial t} dV = - \int_{\$} \rho (\vec{v} \varphi - \lambda_\Phi \operatorname{grad} \varphi) \cdot d\vec{S} + c_* \left(\varphi_* \int_{\mathbb{A}} dV - \int_{\mathbb{A}} \varphi dV \right)$$

$$\vec{v} \cdot d\vec{S} = \sum_{i=1}^3 v_i dS_i \quad \text{e} \quad \operatorname{grad} \varphi \cdot d\vec{S} = \sum_{i=1}^3 \frac{\partial \varphi}{\partial x_i} dS_i$$

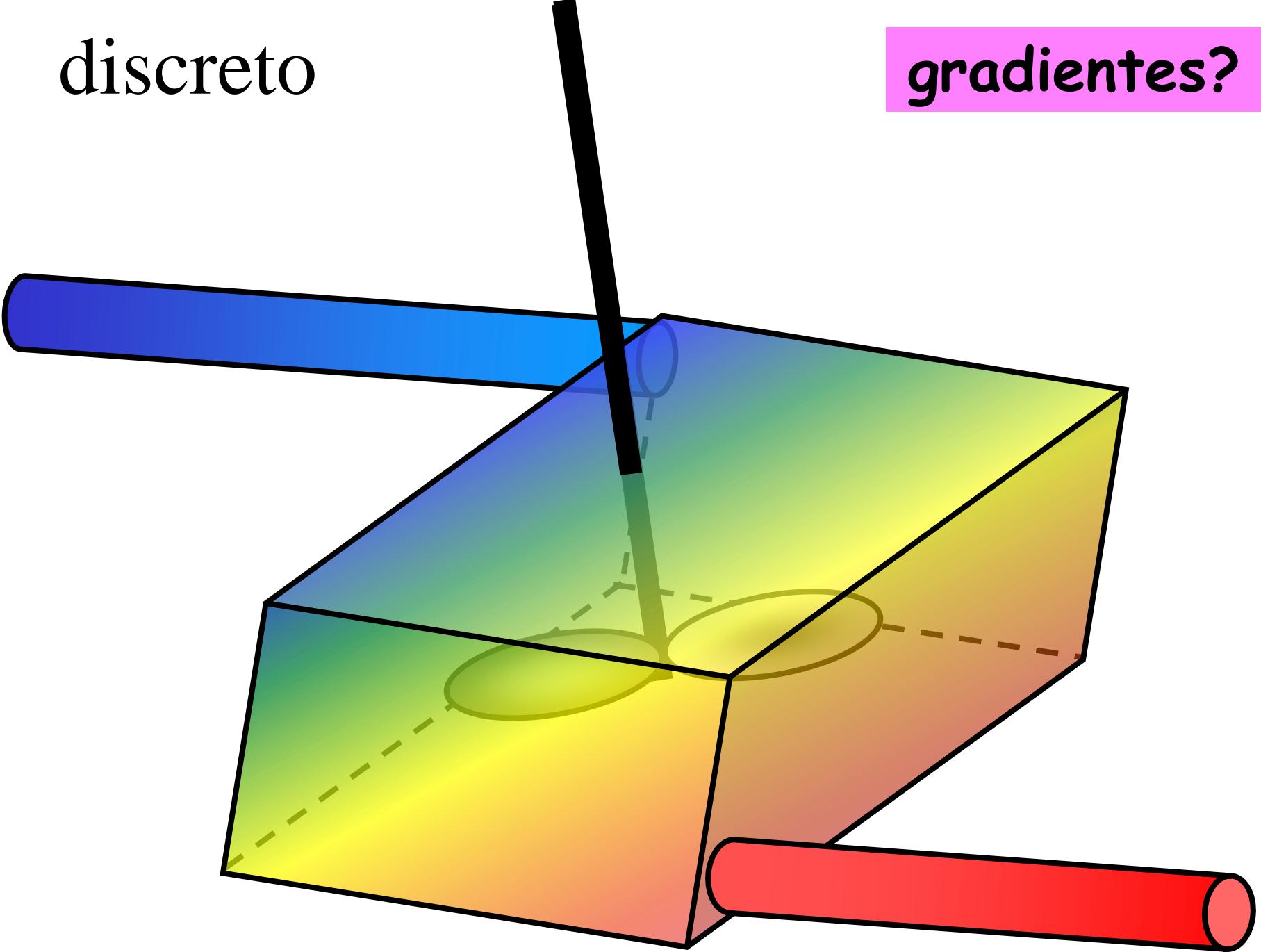


$$\int_{\mathbb{A}} \frac{\partial \rho \varphi}{\partial t} dV = - \int_{\$} \sum_{i=1}^3 \rho \left(v_i \varphi - \lambda_\Phi \frac{\partial \varphi}{\partial x_i} \right) dS_i + c_* \left(\varphi_* \int_{\mathbb{A}} dV - \int_{\mathbb{A}} \varphi dV \right)$$

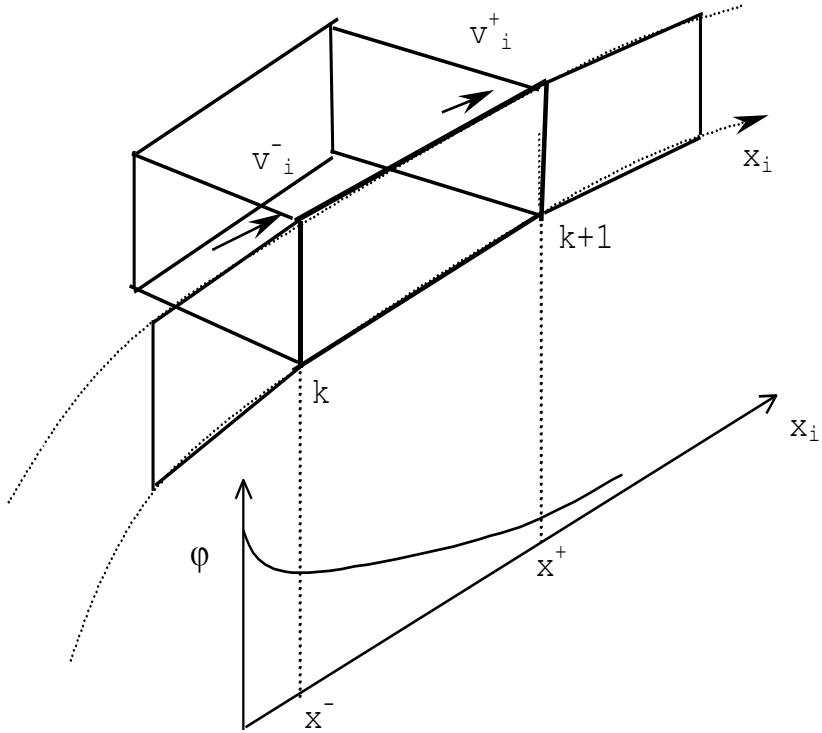
Δ

discreto

gradientes?



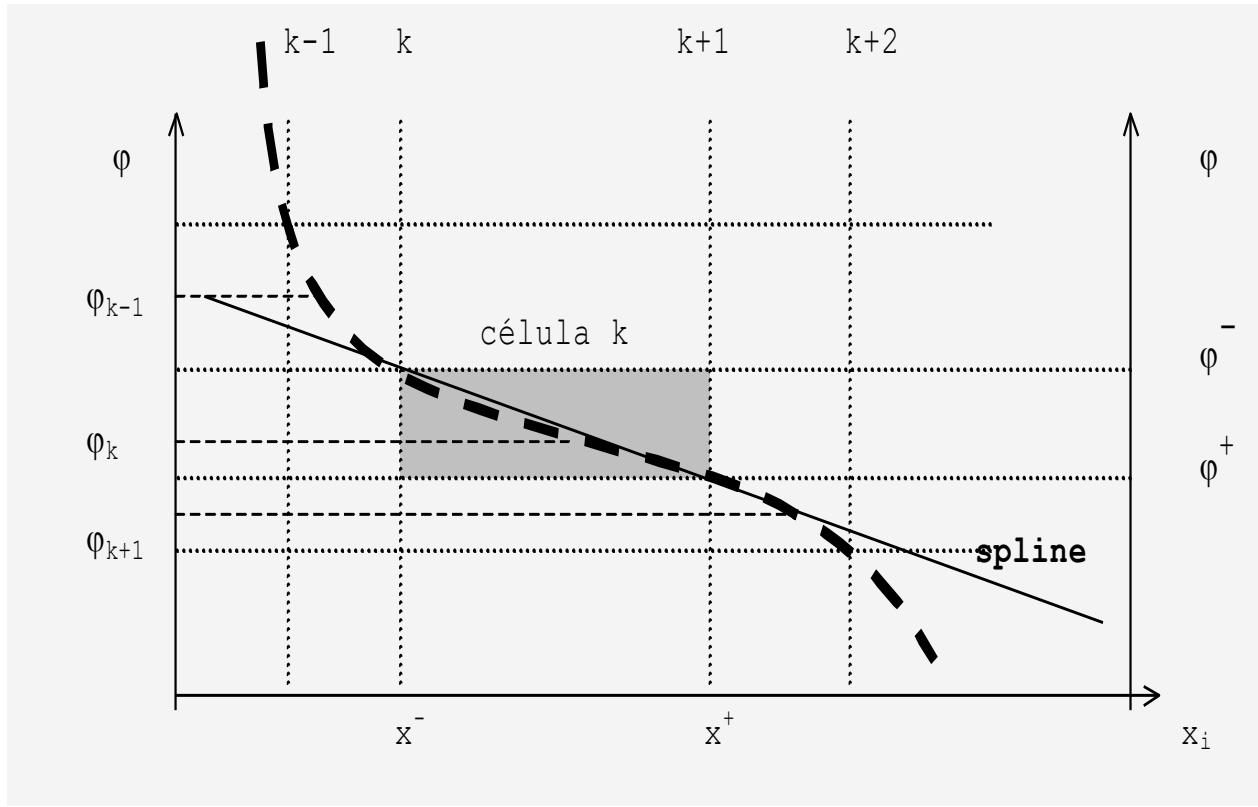
$$\int_{\mathbb{V}} \frac{\partial \rho \phi}{\partial t} dV = - \int_{\$} \sum_{i=1}^3 \rho \left(v_i^- \phi^- - \lambda_\Phi^- \frac{\partial \phi}{\partial x_i} \Big|_- \right) dS_i + c_* \left(\phi_* \int_{\mathbb{V}} dV - \int_{\mathbb{A}} \phi dV \right)$$



$$\forall = \frac{(s_i^- + s_i^+)}{2} (x_i^- - x_i^+) = \prod_{i=1}^3 (x_i^- - x_i^+)$$

$$\frac{\partial \rho \phi}{\partial t} \forall = \sum_{i=1}^3 \left\{ \rho^- \left(v_i^- \phi_i^- - \lambda_\Phi^- \frac{\partial \phi}{\partial x_i} \Big|_- \right) \$_i^- + \rho^+ \left(-v_i^+ \phi_i^+ + \lambda_\Phi^+ \frac{\partial \phi}{\partial x_i} \Big|_+ \right) \$_i^+ \right\} + c_* (\phi_* - \phi) \forall$$

$$\frac{\partial \rho \phi}{\partial t} \forall = \sum_{i=1}^3 \left\{ \rho^- \left(v_i^- \phi_i^- - \lambda_\Phi^- \left. \frac{\partial \phi}{\partial x_i} \right|_- \right) \$_i^- + \rho^+ \left(-v_i^+ \phi_i^+ + \lambda_\Phi^+ \left. \frac{\partial \phi}{\partial x_i} \right|_+ \right) \$_i^+ \right\} + c_* (\phi_* - \phi) \forall$$



$$\begin{aligned} v_i^- &= v_k \\ v_i^+ &= v_{k+1} \end{aligned}$$

$$\phi \cong \phi_k$$

$$\lambda_i^- = \frac{\lambda_k + \lambda_{k-1}}{2}$$

$$\lambda_i^+ = \frac{\lambda_{k+1} + \lambda_{k+2}}{2}$$

$$\frac{\partial \rho \phi}{\partial t} \forall = \sum_{i=1}^6 \pm \rho^\pm \left(v_i^\pm \phi^\pm - \lambda_\Phi^\pm \left. \frac{\partial \phi}{\partial x_i} \right|_\pm \right) \$_i^\pm + c_* (\phi_* - \phi) \forall$$

$$\frac{\partial \rho \varphi}{\partial t} \forall = \sum_{i=1}^6 \pm \left(\rho^\pm v_i^\pm \varphi^\pm - \rho^\pm \lambda_\Phi^\pm \frac{\partial \varphi}{\partial x_i} \Big|_\pm \right) S_i^\pm + c_* (\varphi_* - \varphi) \forall$$

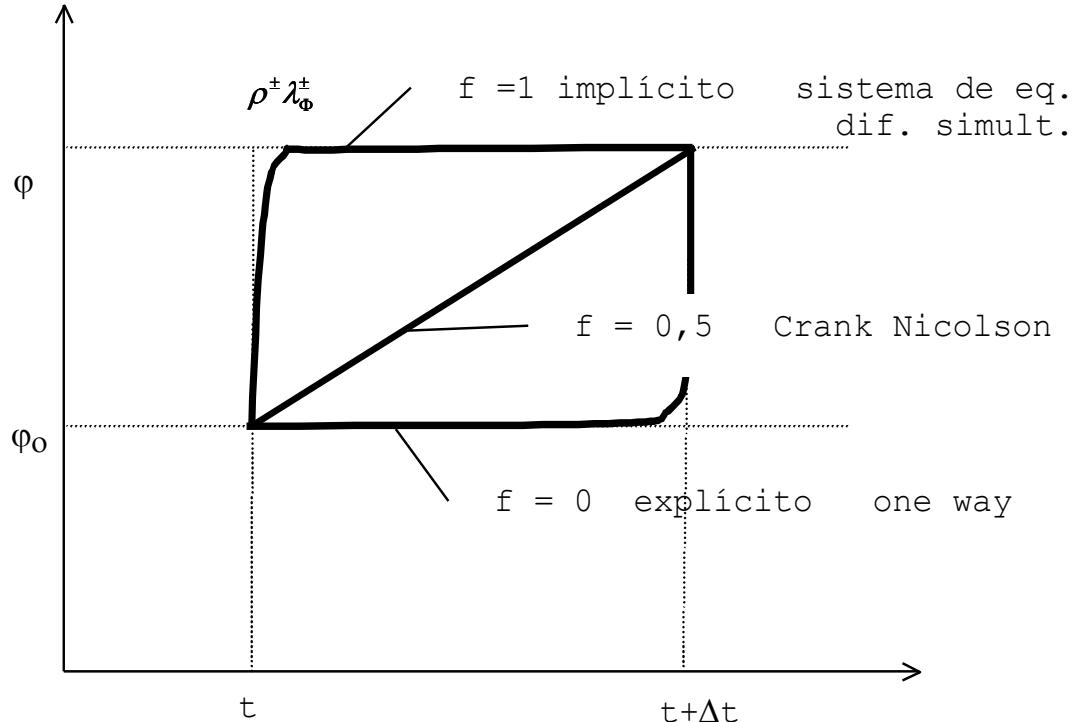
$$\Gamma_\Phi^\pm = \rho^\pm \lambda_\Phi^\pm$$

$$\frac{\partial \rho \varphi}{\partial t} \cong \frac{\Delta \rho \varphi}{\Delta t} = \frac{\rho \varphi - (\rho \varphi)_o}{\Delta t}$$

$f = 0 \rightarrow$ método **explicito**

$$\bar{\varphi} = \frac{1}{\Delta t} \int_t^{t+\Delta t} \varphi dt \cong \\ \cong f \varphi + (1 - f) \varphi_o$$

$$\Delta t < \frac{(\Delta x)^2}{2 \lambda_\Phi}$$



$$\frac{\partial \rho\phi}{\partial t} \quad \forall = \sum_{i=1}^3 \left\{ \rho^- \left(v_i^- \phi_i^- - \lambda_\Phi^- \left. \frac{\partial \phi}{\partial x_i} \right|_- \right) \$_i^- + \rho^+ \left(-v_i^+ \phi_i^+ + \lambda_\Phi^+ \left. \frac{\partial \phi}{\partial x_i} \right|_+ \right) \$_i^+ \right\} + c_* (\phi_* - \phi) \quad \forall$$

$$\begin{aligned}\lambda_i^- &= \frac{\lambda_k + \lambda_{k-1}}{2} \\ \lambda_i^+ &= \frac{\lambda_{k+1} + \lambda_{k+2}}{2}\end{aligned}$$

difusão

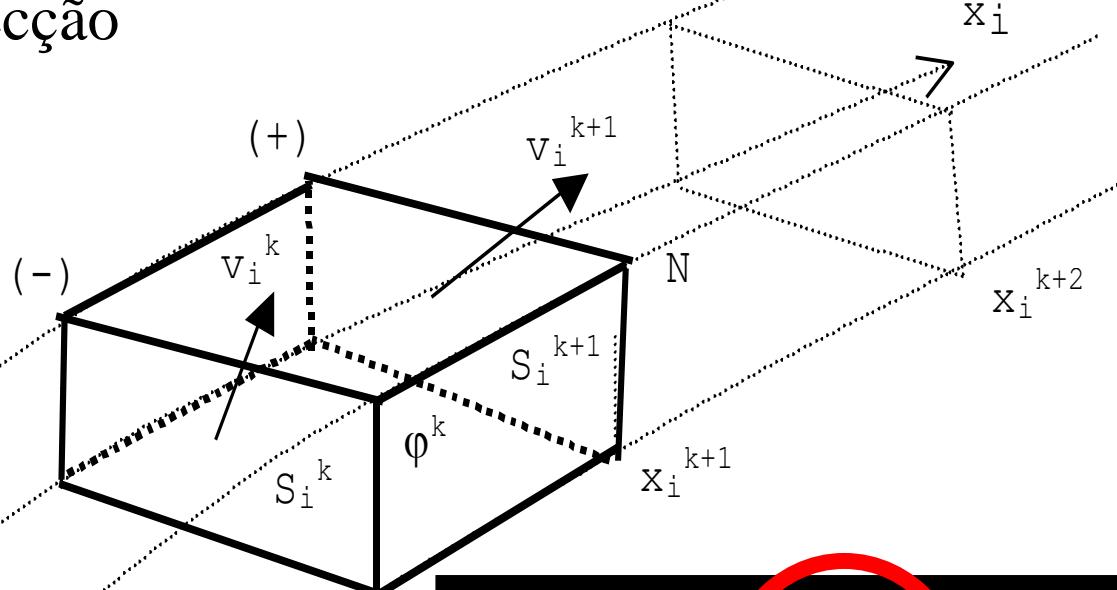
$$\frac{\partial \rho \phi}{\partial t} \quad \forall = \sum_{i=1}^6 \pm \rho^\pm \left(v_i^\pm \phi^\pm - \lambda_\Phi^\pm \left. \frac{\partial \phi}{\partial x_i} \right|_\pm \right) \$_i^\pm + c_* (\phi_* - \phi) \quad \forall$$

tempo

$$\frac{\partial \rho \varphi}{\partial t} \forall = \sum_{i=1}^6 \pm \left(\rho^\pm v_i^\pm \varphi^\pm - \rho^\pm \lambda_\Phi^\pm \left. \frac{\partial \varphi}{\partial x_i} \right|_\pm \right) S_i^\pm + c_* (\varphi_* - \varphi) \forall$$

$$\frac{\partial \rho \varphi}{\partial t} \approx \frac{\Delta \rho \varphi}{\Delta t} = \frac{\rho \varphi - (\rho \varphi)_0}{\Delta t}$$

convecção

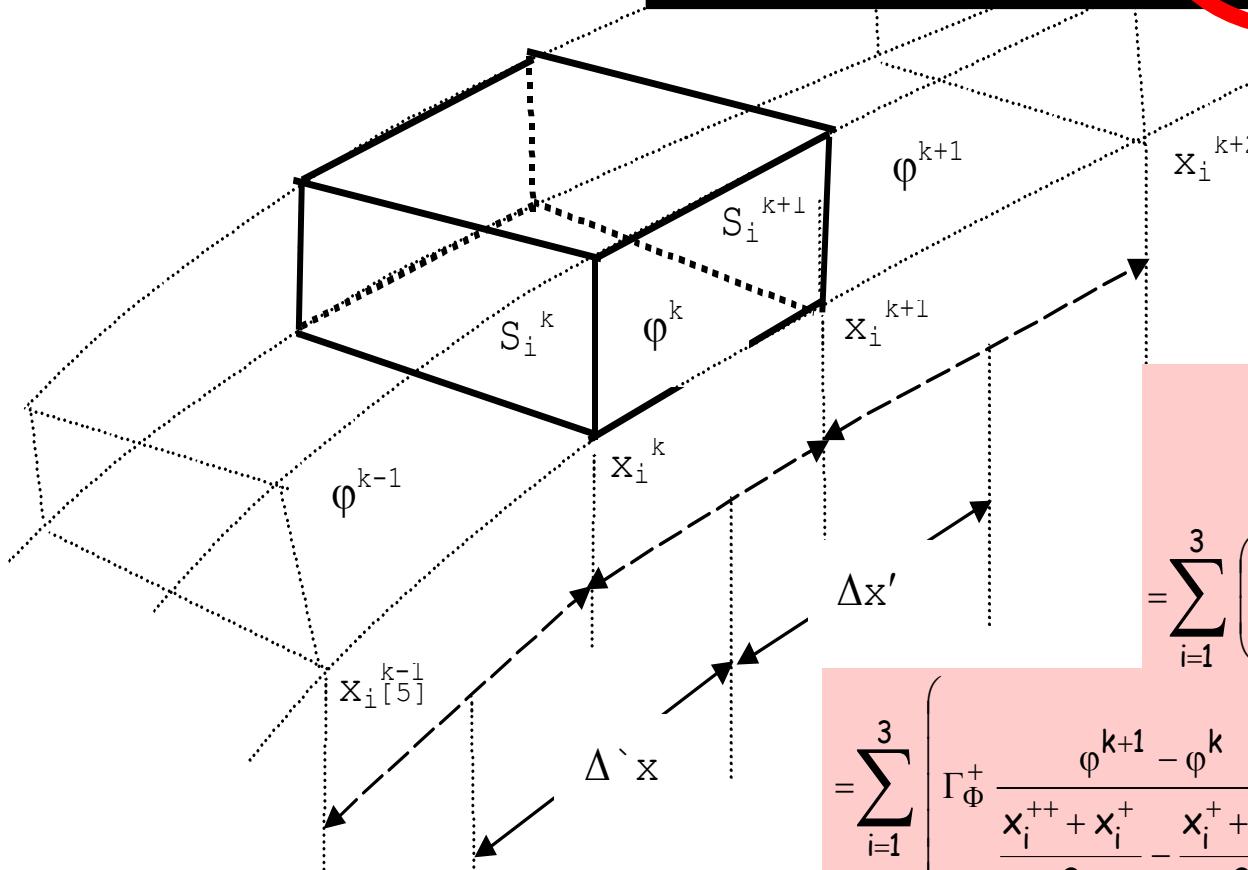


$$\frac{\partial \rho \phi}{\partial t} \nabla = \sum_{i=1}^6 \pm \left(\rho^\pm v_i^\pm \phi^\pm - \rho^\pm \lambda_\Phi^\pm \frac{\partial \phi}{\partial x_i} \right) S_i^\pm + c_* (\phi_* - \phi) \nabla$$

$$\sum_{i=1}^3 \pm \rho^\pm v_i^\pm \phi^\pm S_i^\pm \approx \sum_{i=1}^3 \left(\rho^- \phi^- v_i^- S_i^- - \rho^+ \phi^+ v_i^+ S_i^+ \right)$$

difusão

$$\frac{\partial \rho \varphi}{\partial t} \forall = \sum_{i=1}^6 \pm \left(\rho^\pm v_i^\pm \varphi^\pm - \rho^\pm \lambda_\Phi^\pm \frac{\partial \varphi}{\partial x_i} \right) S_i^\pm + c_* (\varphi_* - \varphi) \forall$$



$$\begin{aligned} \sum_{i=1}^6 \pm \Gamma_\Phi^\pm \frac{\partial \varphi}{\partial x_i} \Big|_\pm S_i^\pm &= \\ = \sum_{i=1}^3 \left(\Gamma_\Phi^+ \frac{\varphi^{k+1} - \varphi^k}{\Delta x'} S_i^+ - \Gamma_\Phi^- \frac{\varphi^k - \varphi^{k-1}}{\Delta' x} S_i^- \right) &= \\ = \sum_{i=1}^3 \left(\Gamma_\Phi^+ \frac{\varphi^{k+1} - \varphi^k}{\frac{x_i^{++} + x_i^+ - x_i^+ + x_i^-}{2}} S_i^+ - \Gamma_\Phi^- \frac{\varphi^k - \varphi^{k-1}}{\frac{x_i^+ + x_i^- - x_i^- + x_i^{--}}{2}} S_i^- \right) &= \end{aligned}$$

$$\sum_{i=1}^3 2 \left(\Gamma_\Phi^+ \frac{\varphi^{k+1} - \varphi^k}{x_i^{++} - x_i^-} S_i^+ - \Gamma_\Phi^- \frac{\varphi^k - \varphi^{k-1}}{x_i^+ - x_i^{--}} S_i^- \right)$$

difusão

juntando

$$\frac{\partial \rho \varphi}{\partial t} \forall = \sum_{i=1}^6 \pm \rho^\pm \left(v_i^\pm \varphi^\pm - \lambda_\Phi^\pm \left. \frac{\partial \varphi}{\partial x_i} \right|_\pm \right) \$_i^\pm + c_* (\varphi_* - \varphi) \forall$$

$$\frac{\partial \rho \varphi}{\partial t} \cong \frac{\rho^k \varphi^k - \rho_0^k \varphi_0^k}{\Delta t}$$

$$c_* (\varphi_* - \varphi^k) \forall$$

$$\sum_{i=1}^3 \pm \rho^\pm v_i^\pm \varphi^\pm S_i^\pm \cong \sum_{i=1}^3 \left(\rho^- \varphi^- v_i^- S_i^- - \rho^+ \varphi^+ v_i^+ S_i^+ \right)$$

$$\sum_{i=1}^3 2 \left(\Gamma_\Phi^+ \frac{\varphi^{k+1} - \varphi^k}{x_i^{++} - x_i^-} S_i^+ - \Gamma_\Phi^- \frac{\varphi^k - \varphi^{k-1}}{x_i^+ - x_i^{--}} S_i^- \right)$$

$$\frac{\rho^k \varphi^k - \rho_0^k \varphi_0^k}{\Delta t} - c_* (\varphi_* - \varphi^k) = \frac{1}{\forall} \sum_{i=1}^3 \left(\rho^- \varphi^- v_i^- S_i^- - 2 \Gamma_\Phi^- \frac{\varphi^k - \varphi^{k-1}}{x_i^+ - x_i^-} S_i^- - \rho^+ \varphi^+ v_i^+ S_i^+ + 2 \Gamma_\Phi^+ \frac{\varphi^{k+1} - \varphi^k}{x_i^{++} - x_i^-} S_i^+ \right)$$

$$\frac{\rho^k \phi^k - \rho_0^k \phi_0^k}{\Delta t} - c_* (\phi_*^k - \phi^k) = \frac{1}{\forall} \sum_{i=1}^3 \left(\rho^- \phi^- v_i^- S_i^- - 2\Gamma_\Phi^- \frac{\phi^k - \phi^{k-1}}{x_i^+ - x_i^{--}} S_i^- - \rho^+ \phi^+ v_i^+ S_i^+ + 2\Gamma_\Phi^+ \frac{\phi^{k+1} - \phi^k}{x_i^{++} - x_i^-} S_i^+ \right)$$

$$\phi^- = \frac{\phi^k + \phi^{k-1}}{2} \quad ; \quad \phi^+ = \frac{\phi^{k+1} + \phi^k}{2}$$

$$\frac{\rho^k \phi^k - \rho_0^k \phi_0^k}{\Delta t} - c_* (\phi_*^k - \phi^k) = \frac{1}{\forall} \sum_{i=1}^3 \left(\frac{\phi^k + \phi^{k-1}}{2} \rho^- v_i^- S_i^- - 2\Gamma_\Phi^- \frac{\phi^k - \phi^{k-1}}{x_i^+ - x_i^{--}} S_i^- - \frac{\phi^{k+1} + \phi^k}{2} \rho^+ v_i^+ S_i^+ + 2\Gamma_\Phi^+ \frac{\phi^{k+1} - \phi^k}{x_i^{++} - x_i^-} S_i^+ \right)$$

$$\begin{aligned} & \phi^k \left[c_* + \frac{\rho^k}{\Delta t} + \sum_{i=1}^3 \underbrace{\left(\frac{2\Gamma_\Phi^+}{\forall} \frac{S_i^+}{x_i^{++} - x_i^-} + \frac{\rho^+ v_i^+ S_i^+}{2\forall} + \frac{2\Gamma_\Phi^-}{\forall} \frac{S_i^-}{x_i^+ - x_i^{--}} - \frac{\rho^- v_i^- S_i^-}{2\forall} \right)}_{a_k} \right] = \\ & = c_* \underbrace{\phi_*^k}_{a_*} + \underbrace{\frac{\rho_0^k}{\Delta t} \phi_0^k}_{a_0} + \sum_{i=1}^3 \left[\underbrace{\phi^{k+1} \left(\frac{2\Gamma_\Phi^+}{\forall} \frac{S_i^+}{x_i^{++} - x_i^-} - \frac{\rho^+ v_i^+ S_i^+}{2\forall} \right)}_{a_i^+} \underbrace{\phi^{k-1} \left(\frac{2\Gamma_\Phi^-}{\forall} \frac{S_i^-}{x_i^+ - x_i^{--}} + \frac{\rho^- v_i^- S_i^-}{2\forall} \right)}_{a_i^-} \right] \end{aligned}$$

$$\varphi^k = \frac{a_* \varphi_*^k + a_0 \varphi_0^k + \sum_{i=1}^3 (a_i^+ \varphi^{k+1} + a_i^- \varphi^{k-1})}{a_* + a_0 + \sum_{i=1}^3 (a_i^+ + a_i^-)} = \frac{\sum_{j=1}^8 a_j \varphi_j}{\sum_{j=1}^8 a_j}$$

$$\varphi = \frac{a_0 \varphi_0 + a_N \varphi_N + a_S \varphi_S + a_E \varphi_E + a_W \varphi_W + a_H \varphi_H + a_L \varphi_L + a_* \varphi_*}{a_0 + a_N + a_S + a_E + a_W + a_H + a_L + a_*}$$

para a discretização não há mais diferença entre entradas e produções

Condições de Contorno ou Fontes

$$0 = \sum_{i=1}^9 S_\phi = \sum_{i=1}^9 CO(VAL - \phi)$$

fixed value
 $CO_{BC} = 10^{20}$

fixed flux
 $CO_{BC} = 10^{-20}$;
 $VAL_{BC} = 10^{20} \cdot \text{flux}$

$$0 = \sum_{i=1}^9 10^{20} (VAL_{BC} - \phi)$$
$$\phi = VAL_{BC}$$

$$S_{\phi*} = 10^{-20} (10^{20} \cdot \text{flux} - \phi)$$

$$S_{\phi*} = \text{flux}$$

$\varphi = P1, U1, V1, W1, H1, KE, EP, A, \dots$ (~ 10)

$x, y, z = (100 . 100 . 100)$ $t = (1000)$

vector PHI com $\sim 10^7$ elementos

$$\sum_{i=1}^7 a_i \varphi_i = C_i$$

present	west	east	low	high	south	north	sources
---------	------	------	-----	------	-------	-------	---------

$$a_{111} \varphi_{111} + a_{011} \varphi_{011} + a_{211} \varphi_{211} + a_{101} \varphi_{101} + a_{121} \varphi_{121} + a_{110} \varphi_{110} + a_{112} \varphi_{112} = C_{\varphi_{111}}$$

$$a_{211} \varphi_{211} + a_{111} \varphi_{111} + a_{311} \varphi_{311} + a_{201} \varphi_{201} + a_{221} \varphi_{221} + a_{210} \varphi_{210} + a_{212} \varphi_{212} = C_{\varphi_{211}}$$

...

$$a_{xyz} \varphi_{xyz} + a_{(x-1)yz} \varphi_{(x-1)yz} + a_{(x+1)yz} \varphi_{(x+1)yz} + \dots + = C_{\varphi_{xyz}}$$

• • •

$$z_{111} \psi_{111} + z_{011} \psi_{011} + z_{211} \psi_{211} + z_{101} \psi_{101} + z_{121} \psi_{121} + z_{110} \psi_{110} + z_{112} \psi_{112} = C_{\psi_{111}}$$

$$z_{211} \psi_{211} + z_{111} \psi_{111} + z_{311} \psi_{311} + z_{201} \psi_{201} + z_{221} \psi_{221} + z_{210} \psi_{210} + z_{212} \psi_{212} = C_{\psi_{211}}$$

...

$$z_{xyz} \psi_{xyz} + z_{(x-1)yz} \psi_{(x-1)yz} + z_{(x+1)yz} \psi_{(x+1)yz} + \dots + = C_{\psi_{xyz}}$$

BC

matrizes

$$a_{ij} \varphi_i = C_j$$

Considerando os termos de entrada e saída como fontes eles vão para o vetor C . Resolvendo por slabs bi dimensionais a matriz a se torna tri-diagonal.

$$\begin{bmatrix} a_1 & a_2 & a_3 & \dots & 0 & 0 \\ 0 & a_4 & a_5 & \dots & 0 & 0 \\ 0 & 0 & a_7 & \dots & 0 & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & \dots & a_{3N-3} & 0 \\ 0 & 0 & 0 & \dots & a_{3N-1} & a_{3N} \end{bmatrix} \cdot \begin{bmatrix} \varphi_1 \\ \varphi_2 \\ \dots \\ \dots \\ \dots \\ \varphi_N \end{bmatrix} = \begin{bmatrix} C_1 \\ C_2 \\ \dots \\ \dots \\ \dots \\ C_N \end{bmatrix}$$

TRIDIAGONAL

PHI

S+BC

em vez de uma matriz $10^4 \times 10^4 = 10^8$ “apenas” 3×10^4 são necessários, assim a solução é possível em PCs.

Convergência

$$\Delta t < \frac{(\Delta x)^2}{2 \lambda_\Phi}$$

implícito x explícito

$$\bar{\varphi} \cong f\varphi + (1-f)\varphi_0$$

relaxação on line

linear f

$$S_{\phi_{\text{relax}}} = f(\varphi_{\text{old}} - \varphi_p) \frac{\rho V}{\Delta t}$$

"dt falso"

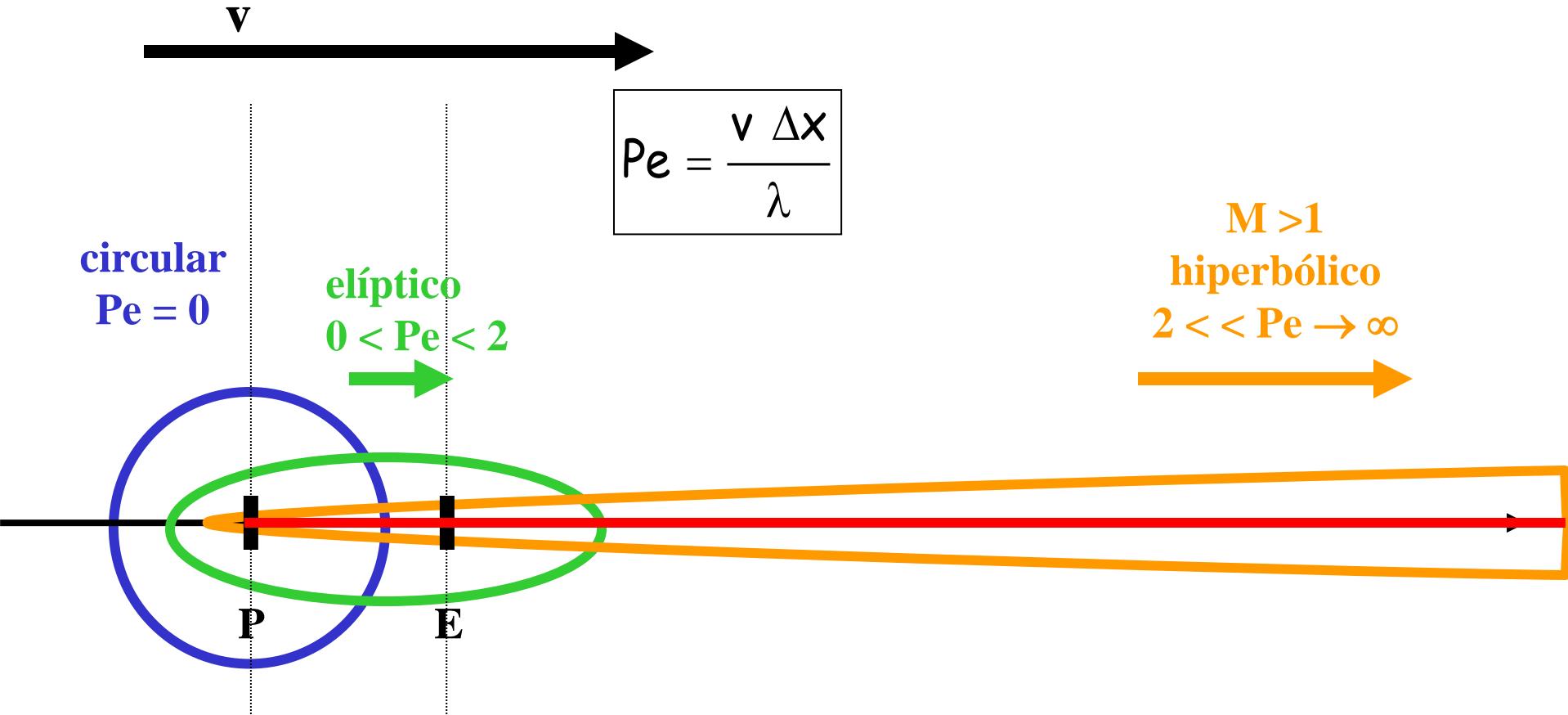
$$S_{\phi_{\text{relax}}} = \left[\frac{\rho V}{dt_f} \right] (\varphi_{\text{old}} - \varphi_p)$$

up-wind

"Pe falso"

$$a_i^\pm = \frac{\bar{s}_i \bar{\lambda}}{\Delta x} \begin{cases} 1 \pm \frac{Pe}{2} & \text{se } Pe < 2 \\ Pe & \text{se } Pe > 2 \end{cases}$$

upwind

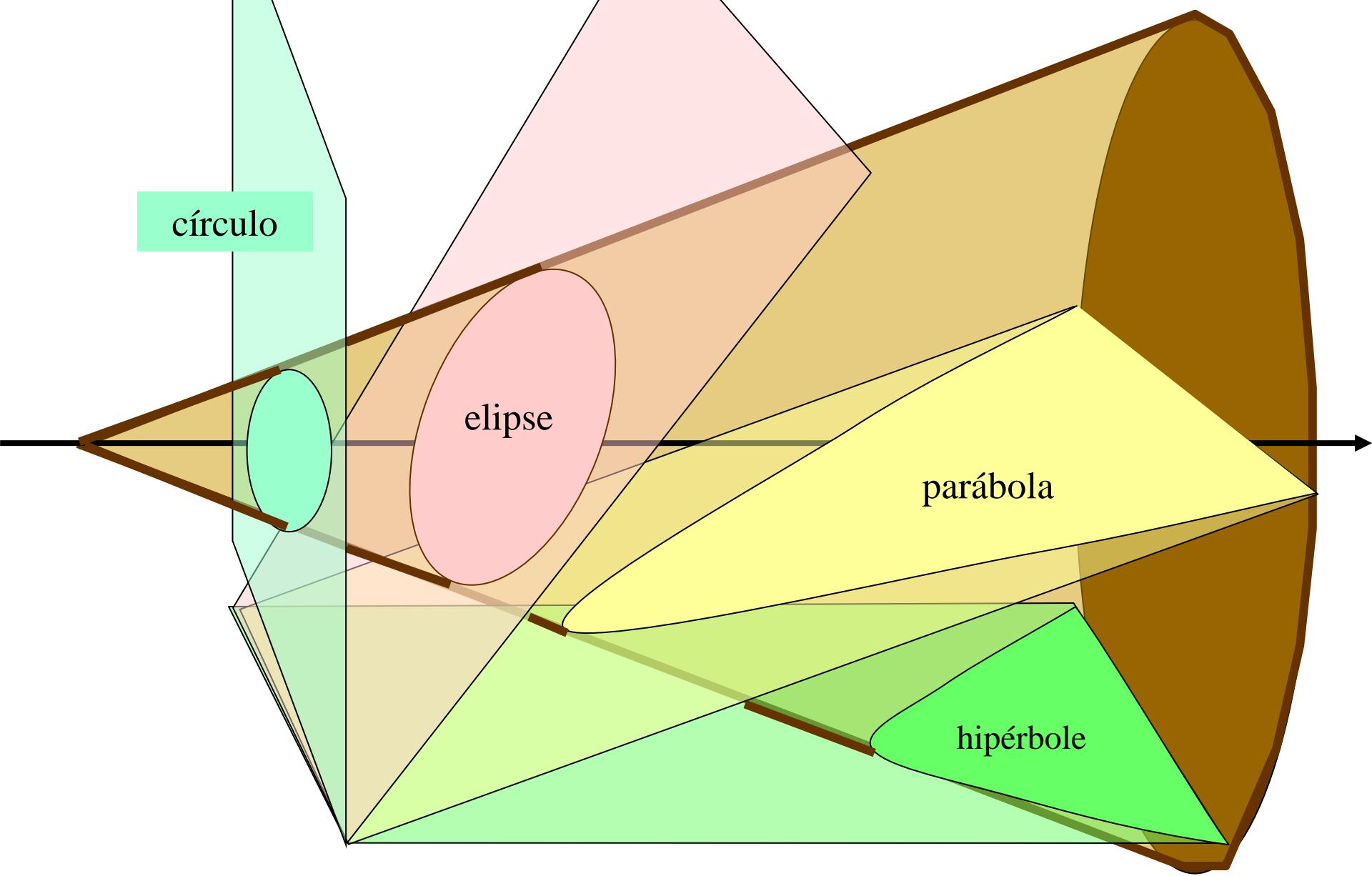


$$a_i^{\pm} = \frac{\bar{S}_i \bar{\lambda}}{\Delta x} \left[1 \pm \frac{Pe}{2} \right] \quad \rightarrow \quad a_i^{\pm} = \frac{\bar{S}_i \bar{\lambda}}{\Delta x} \max \left\{ \pm Pe ; 0 ; \left[1 \pm \frac{Pe}{2} \right] \right\}$$

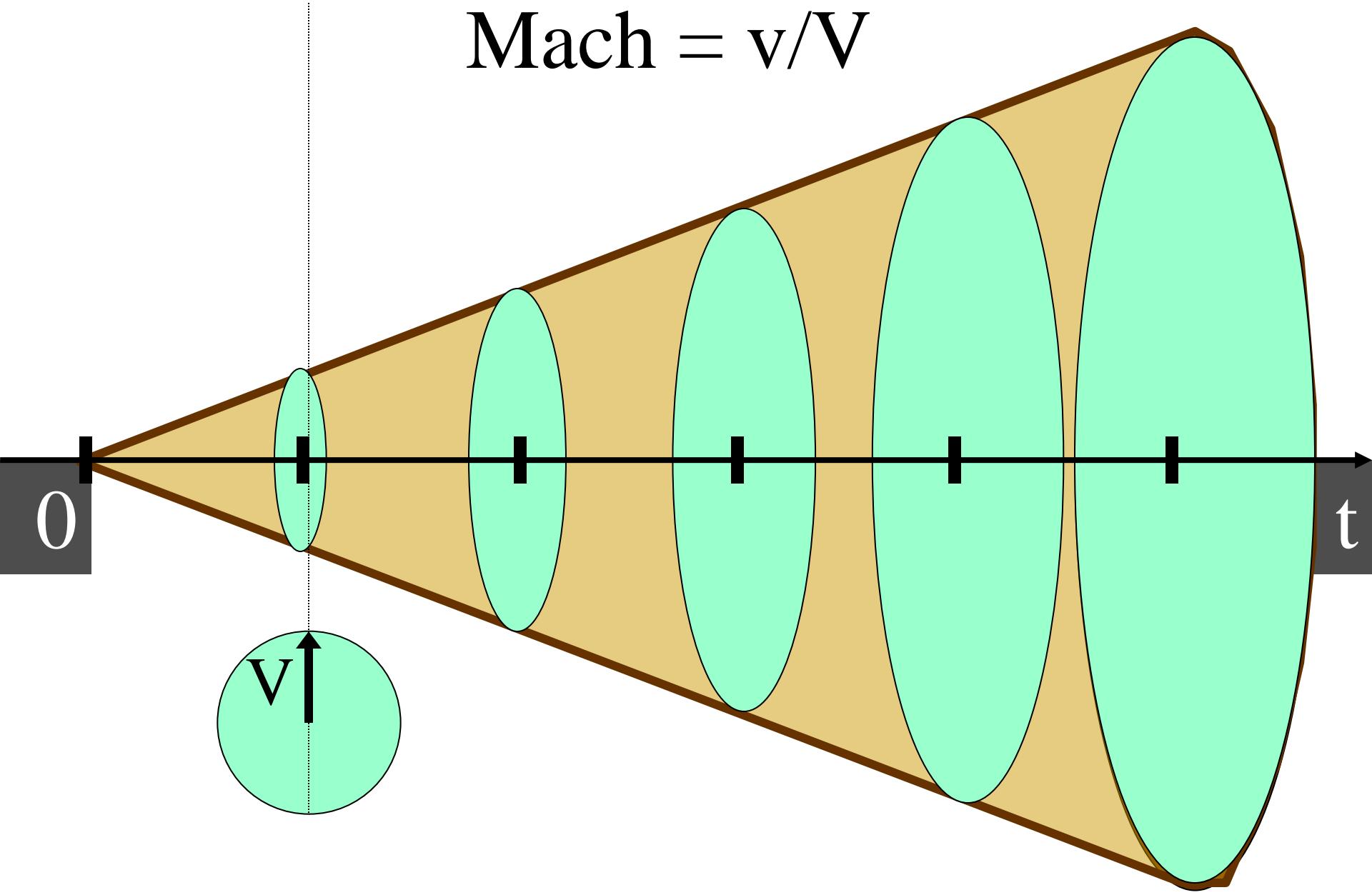
híbrido

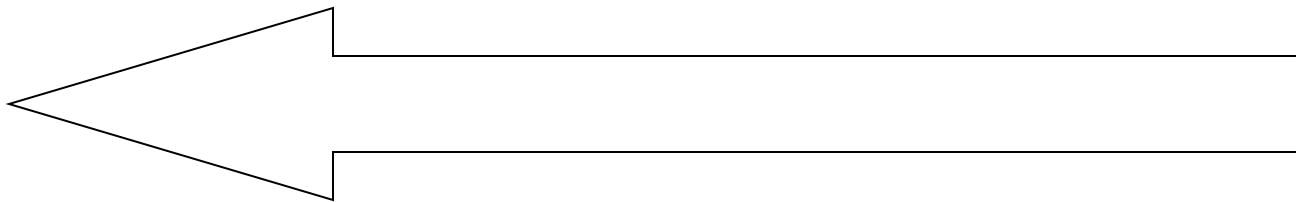
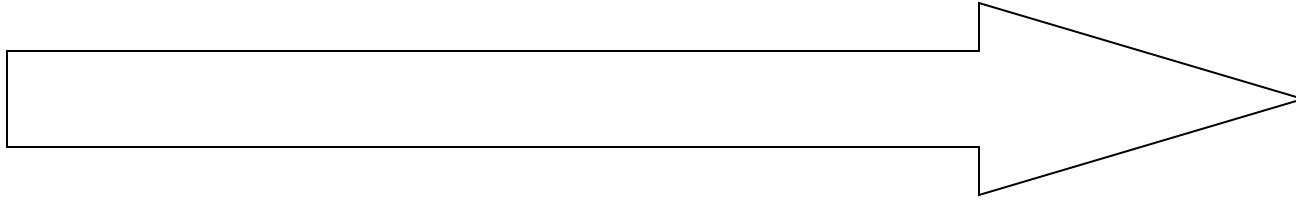
ver pág 182 Versteeg

cônicas



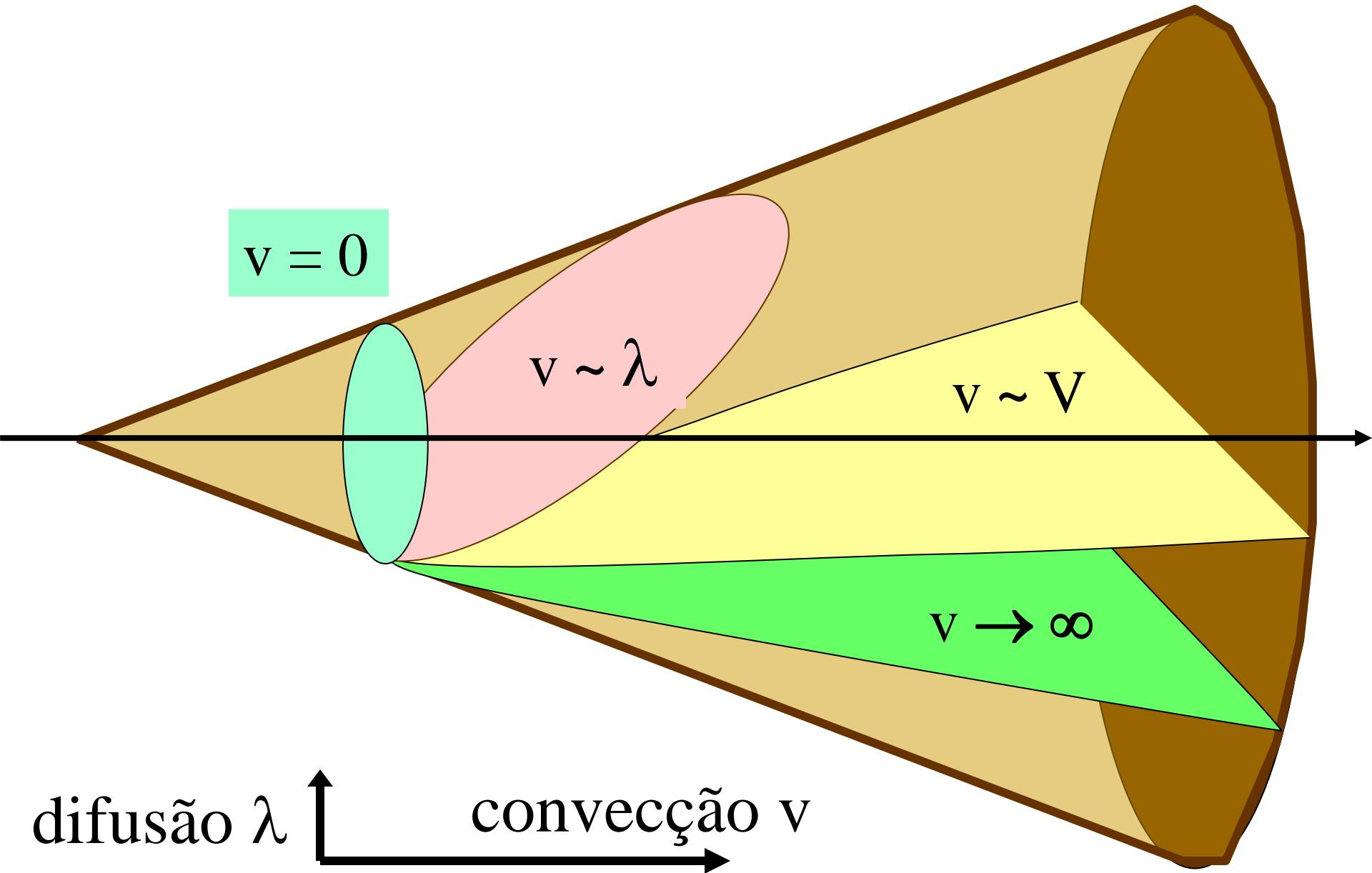
$$\text{Mach} = v/V$$

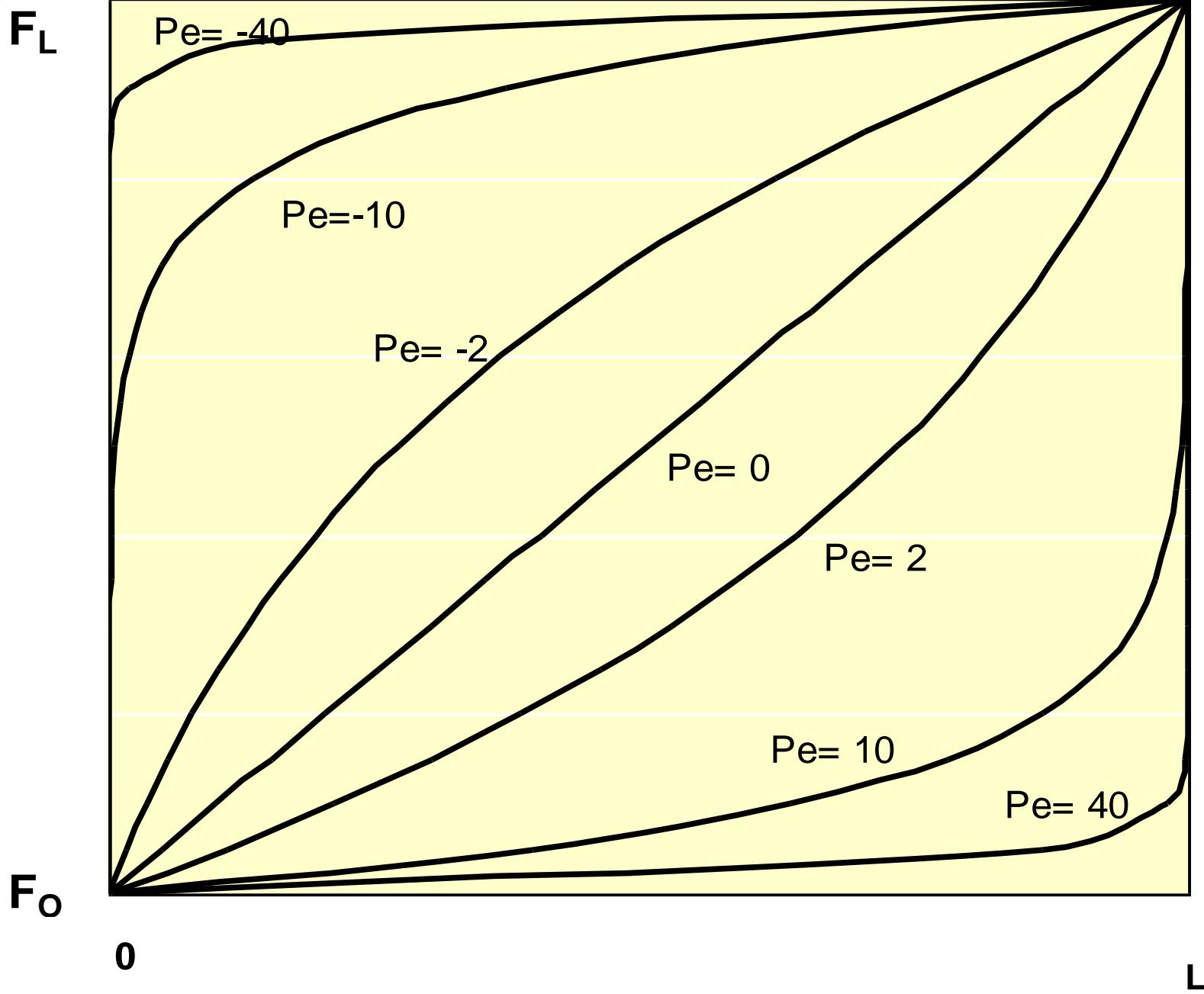




caçada

upwind





coordinate direction. The discretised equation that covers all cases is given by

$$a_P \phi_P = a_W \phi_W + a_E \phi_E + a_S \phi_S + a_N \phi_N + a_B \phi_B + a_T \phi_T \quad (5.43)$$

with central coefficient

$$a_P = a_W + a_E + a_S + a_N + a_B + a_T + \Delta F$$

and the coefficients of this equation for the **hybrid differencing scheme** are as follows:

	One-dimensional flow	Two-dimensional flow	Three-dimensional flow
a_W	$\max\left[F_w, \left(D_w + \frac{F_w}{2}\right), 0\right]$	$\max\left[F_w, \left(D_w + \frac{F_w}{2}\right), 0\right]$	$\max\left[F_w, \left(D_w + \frac{F_w}{2}\right), 0\right]$
a_E	$\max\left[-F_e, \left(D_e - \frac{F_e}{2}\right), 0\right]$	$\max\left[-F_e, \left(D_e - \frac{F_e}{2}\right), 0\right]$	$\max\left[-F_e, \left(D_e - \frac{F_e}{2}\right), 0\right]$
a_S	-	$\max\left[F_s, \left(D_s + \frac{F_s}{2}\right), 0\right]$	$\max\left[F_s, \left(D_s + \frac{F_s}{2}\right), 0\right]$
a_N	-	$\max\left[-F_n, \left(D_n - \frac{F_n}{2}\right), 0\right]$	$\max\left[-F_n, \left(D_n - \frac{F_n}{2}\right), 0\right]$
a_B	-	-	$\max\left[F_b, \left(D_b + \frac{F_b}{2}\right), 0\right]$
a_T	-	-	$\max\left[-F_t, \left(D_t - \frac{F_t}{2}\right), 0\right]$
ΔF	$F_e - F_w$	$F_e - F_w + F_n - F_s$	$F_e - F_w + F_n - F_s + F_t - F_b$

In the above expressions the values of F and D are calculated with the following formulae:

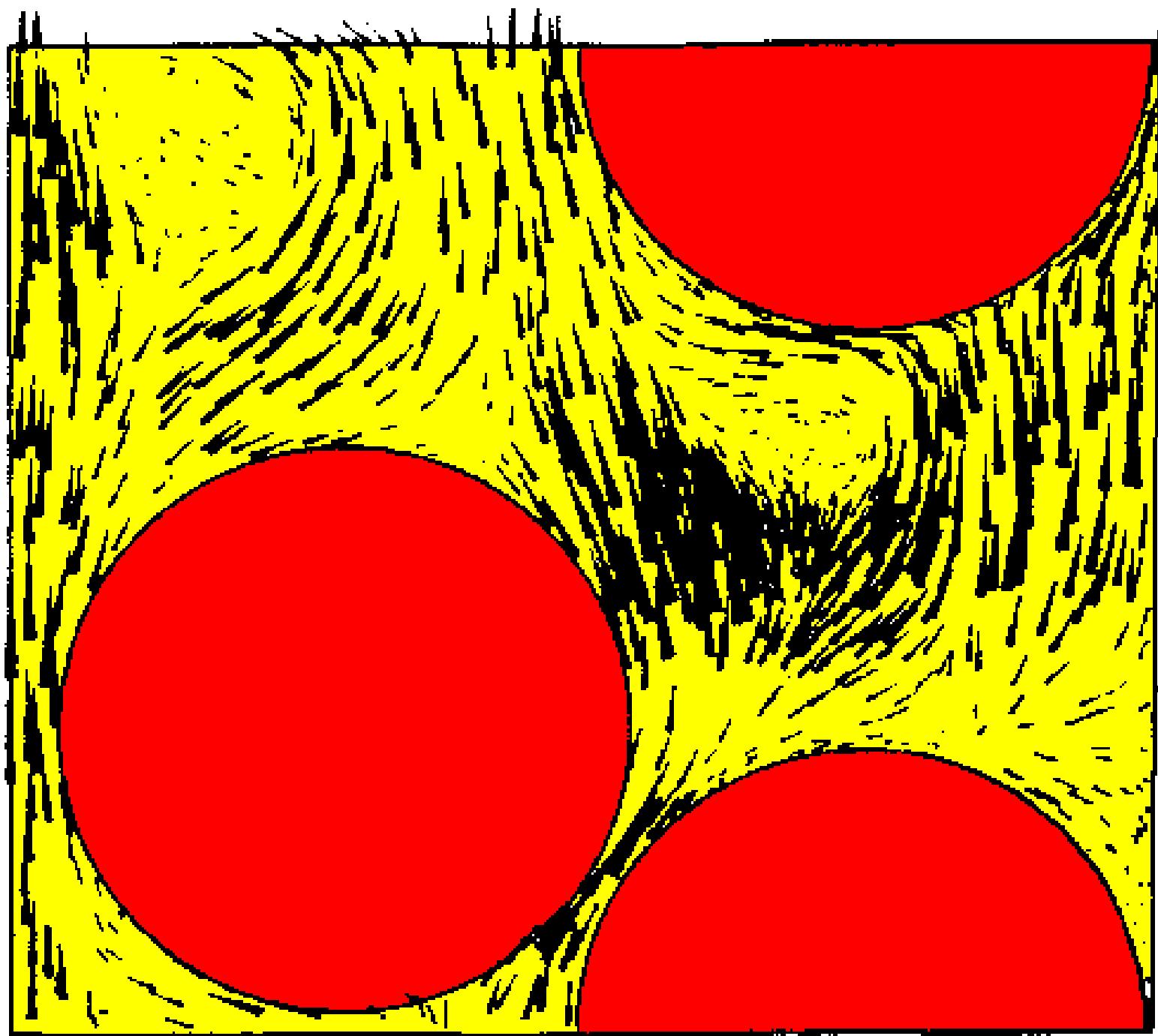
Face	w	e	s	n	b	t
F	$(\rho u)_w A_w$	$(\rho u)_e A_e$	$(\rho v)_s A_s$	$(\rho v)_n A_n$	$(\rho w)_b A_b$	$(\rho w)_t A_t$
D	$\frac{\Gamma_w}{\delta x_{WP}} A_w$	$\frac{\Gamma_e}{\delta x_{PE}} A_e$	$\frac{\Gamma_s}{\delta y_{SP}} A_s$	$\frac{\Gamma_n}{\delta y_{PN}} A_n$	$\frac{\Gamma_b}{\delta z_{PN}} A_b$	$\frac{\Gamma_t}{\delta z_{PT}} A_t$

Modifications to these coefficients to cater for boundary conditions in two and three dimensions are available in the form of expressions such as (5.40).

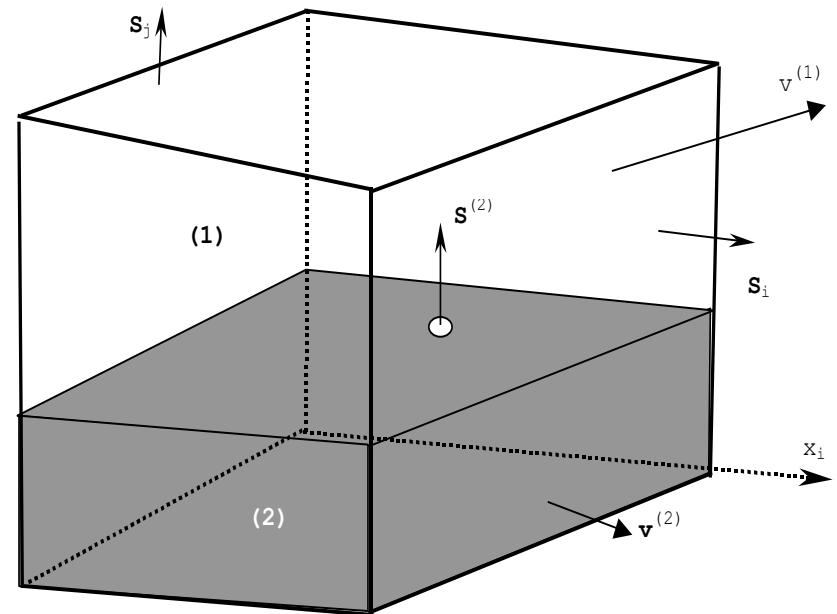
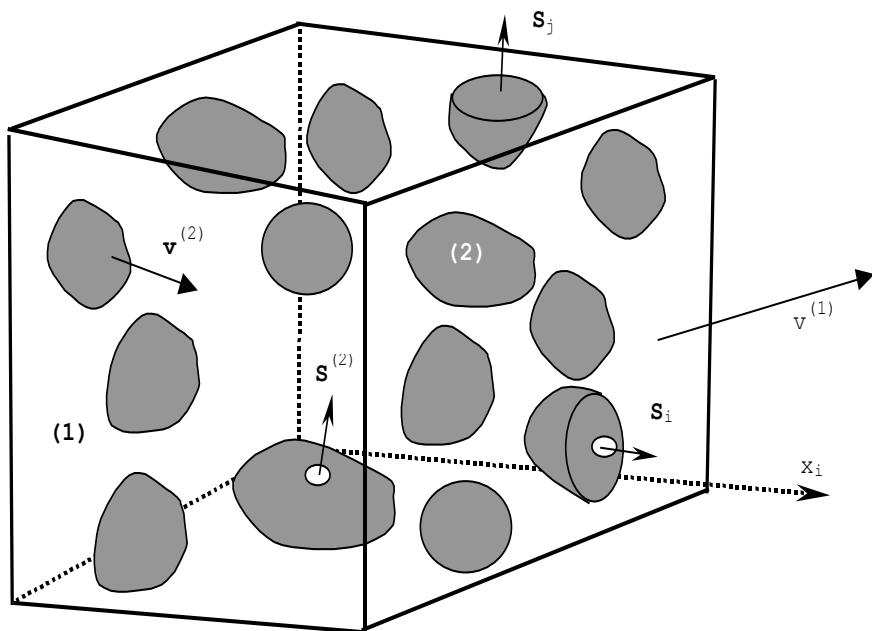


Bifásico

célula
mínima



bifásicos

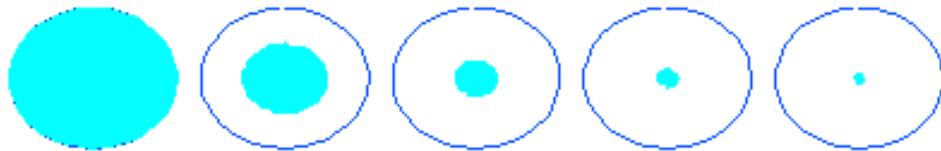


$$\frac{A_{\text{efetiva}}^{(1,2)}}{\forall} \approx N_p \phi_p \pi \langle D_p \rangle^2$$

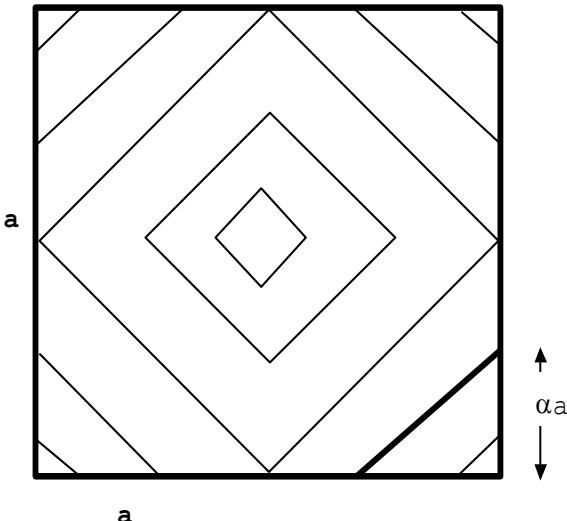
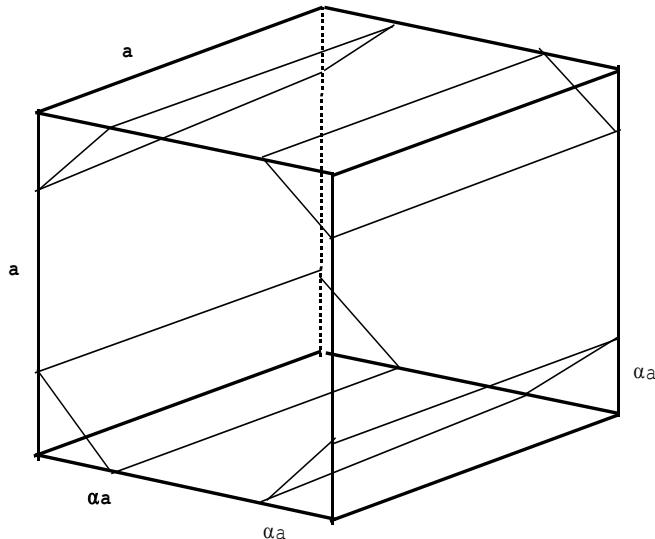
$$r = N_p \frac{4 \pi}{3} \left(\frac{D_p}{2} \right)^3 = N_p \frac{\pi}{6} \langle D_p \rangle^3$$

Particle Size Calculation

In many practical cases, the particle size will vary throughout the domain, as the result of combustion or evaporation/condensation.



- This can be dealt with by use of the SHADOW technique. This uses a third phase (the SHADOW phase), which behaves like the disperse phase (usually phase 2), but without interphase mass transfer.
- On the picture, the solid particles are phase 2, and the open ones the shadow phase.
- Then changes in particle size can be calculated from local volume fraction ratios:
$$\frac{D_p}{D_{p,in}} = \left(\frac{R_2}{R_S} \right)^{1/3}$$



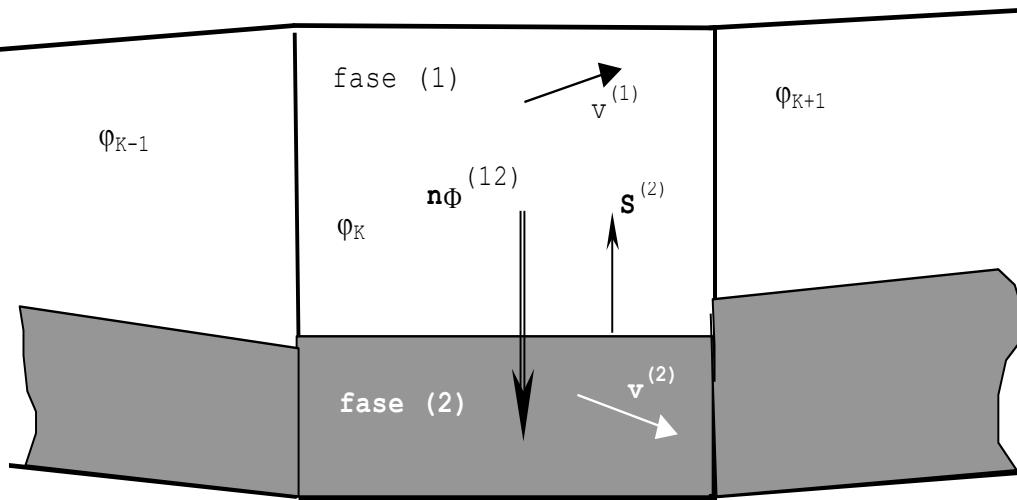
$$r = \frac{V^{(\alpha)}}{V} = \frac{4 \frac{\alpha a \alpha a}{2} a}{a^3} \rightarrow \alpha = \sqrt{\frac{r}{2}}$$

$$\frac{A^{(12)}}{\$} = 4 a \left(2 \alpha a \cos 45^\circ \right) = 4 a^2 \sqrt{r} \rightarrow \frac{A^{(12)}}{\$} = \frac{4 a^2 \sqrt{r}}{6 a^2} = \frac{2\sqrt{r}}{3}$$

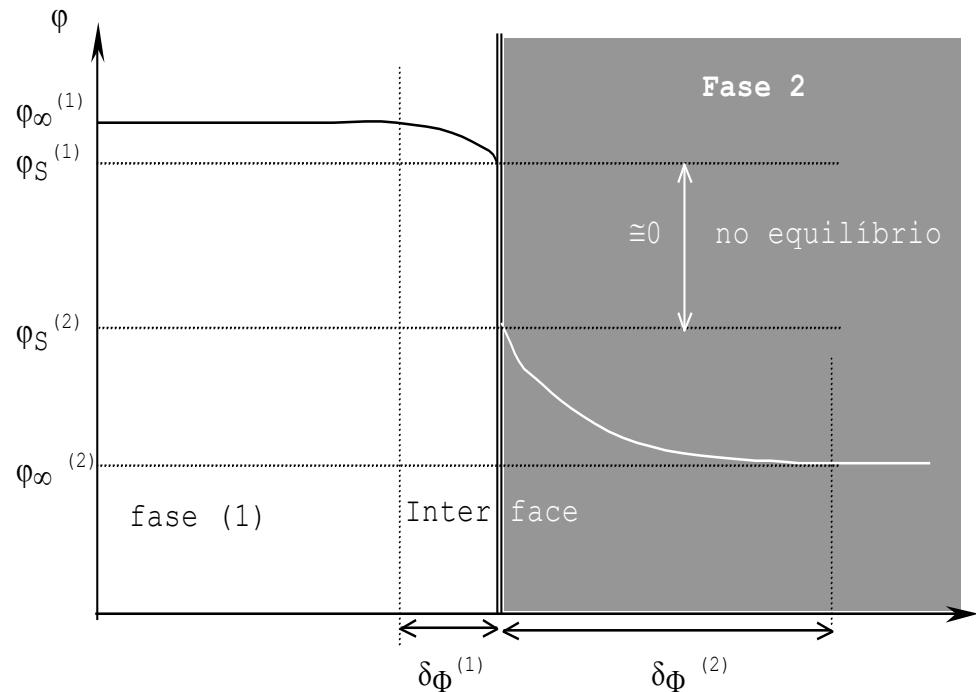
$$\frac{A^{(12)}}{\$} \approx \begin{cases} 2/3 \sqrt{r} & p / 0 \leq r \leq 0,5 \\ 2/3 \sqrt{(1 - r)} & p / 0,5 \leq r \leq 1 \end{cases}$$

$$\frac{A^{(12)}}{\$} \cong 4(r - r^2)$$

$$a_p = \frac{A_{\text{efetiva}}^{(12)}}{A^{(12)}} \cong \frac{\frac{N_p}{\$} \phi_p \pi \langle D_p \rangle^2}{4(r - r^2)} = \frac{\frac{N_p}{\$} \phi_p \pi \langle D_p \rangle^2}{r \$ 4(1 - r)}$$



$$\left| \vec{n}_\Phi^{(12)} \right| \approx \lambda_\Phi^{(1)} \frac{\phi^{(1)} - \phi_S^{(1)}}{\delta_\Phi^{(1)}} \approx \lambda_\Phi^{(2)} \frac{\phi^{(2)} - \phi_S^{(2)}}{\delta_\Phi^{(2)}}$$



$$\left| \vec{n}_{\Phi}^{(12)} \right| = c_{\Phi}^{(12)} \left(\varphi_{\infty}^{(1)} - \varphi_{\infty}^{(2)} \right) \approx \lambda_{\Phi}^{(1)} \frac{\varphi_{\infty}^{(1)} - \varphi_{\text{S}}^{(1)}}{\delta_{\Phi}^{(1)}} \approx \lambda_{\Phi}^{(2)} \frac{\varphi_{\text{S}}^{(2)} - \varphi_{\infty}^{(2)}}{\delta_{\Phi}^{(2)}}$$

$$c_{\Phi}^{(12)} = \left[\frac{\delta_{\Phi}^{(1)}}{\lambda_{\Phi}^{(1)}} + \frac{\delta_{\Phi}^{(2)}}{\lambda_{\Phi}^{(2)}} \right]^{-1}$$

$$\left| \vec{n}_{\Phi}^{(12)} \right| A_{\text{efetiva}}^{(12)} = c_{\Phi}^{(12)} \left(\varphi^{(1)} - \varphi^{(2)} \right) a_p A^{(12)}$$

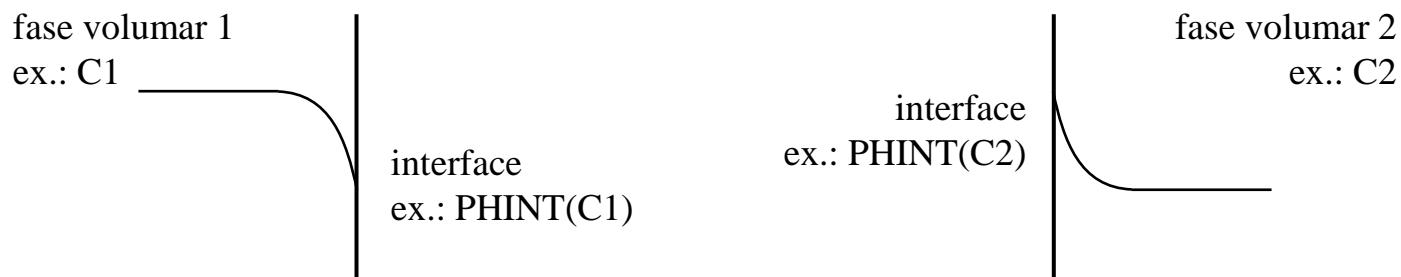
$$S_{\Phi}^{(\alpha)} = c_{\Phi}^{(12)} \left(\varphi^{(1)} - \varphi^{(2)} \right) a_p A^{(12)} = c_{\Phi}^{(\alpha)} \left(\varphi^{(1)} - \varphi^{(2)} \right) r \forall$$

$$c_{\Phi}^{(\alpha)} = \pm \frac{c_{\Phi}^{(12)} a_p}{r} A^{(12)}$$

coeficiente

INTER-PHASE-SLIP ALGORITHM (IPSA) – III

A interface:



Os coeficientes de transferência da fase volumar para a interface são definidos por $CINT(\Phi)$ (ex.: $CINT(C1)$) - diferentes possibilidades de configuração.

Outros termos relevantes:

CMDOT, CFIPS

poderosa bifásica

termo transitório

termo de
convecção

termo de difusão

termo de fonte

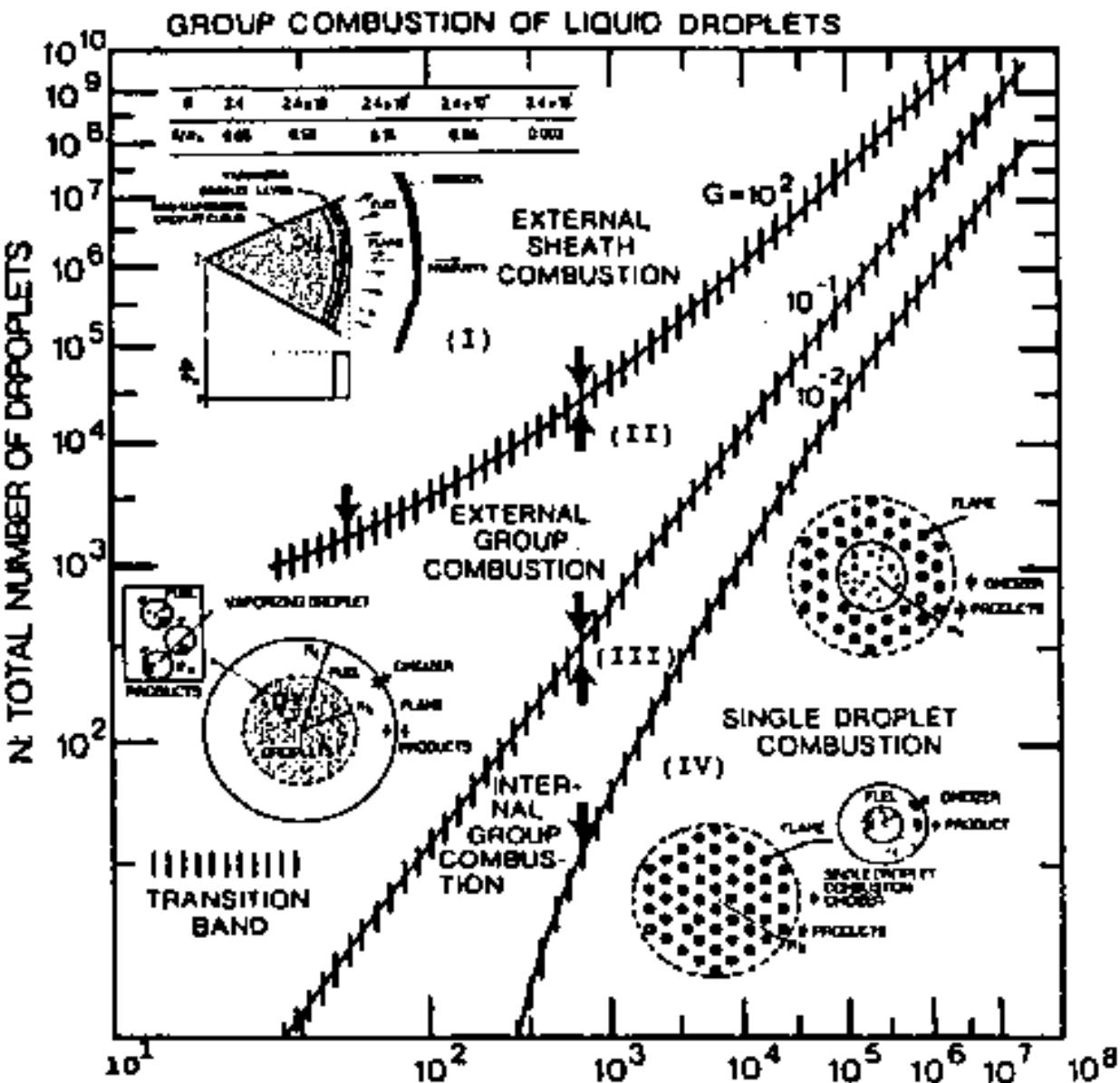
$$\frac{\partial(r_i\rho_i.\Phi)}{\partial t} + \operatorname{div}(r_i.\rho_i.\vec{v}_i.\Phi - r_i.\Gamma.\operatorname{grad}\Phi) = r_i.S_\Phi$$

$$\frac{\partial r \rho \phi}{\partial t} = \sum_{i=1}^6 \pm \frac{r^\pm \rho^\pm}{\forall} \left(v_i^\pm \phi^\pm - \lambda_\Phi^\pm \left. \frac{\partial \phi}{\partial x_i} \right|_\pm \right) \$_i^\pm + r c_* (\phi_* - \phi) + r c_\Phi^{(\alpha)} (\phi^{(\alpha)} - \phi)$$

$$\bar{\phi} = \frac{a_O \bar{\phi}_O + a_N \bar{\phi}_N + a_S \bar{\phi}_S + a_E \bar{\phi}_E + a_W \bar{\phi}_W + a_H \bar{\phi}_H + a_L \bar{\phi}_L + a_* \bar{\phi}_* + a^{(\alpha)} \bar{\phi}^{(\alpha)}}{a_O + a_N + a_S + a_E + a_W + a_H + a_L + a_* + a^{(\alpha)}}$$

spray
combustion

four groups
combustion
modes of a
droplet cloud



$$S = \frac{0,05}{\left(1 + 0,276 Re^{0,5} Pr^{0,33}\right)} \left(\frac{d}{r_e}\right)$$

CFD COMERCIAIS

chemtech

PHOENICS

Prof. Brian
Spalding



FLUENT, ...

Introdução ao PHOENICS

slides da CHEMTECH

As equações diferenciais resolvidas em CFD, na sua forma mais geral, podem ser escritas do seguinte modo:

$$\frac{\partial}{\partial t}(\rho\phi) + \frac{\partial}{\partial x}(\rho u\phi) + \frac{\partial}{\partial y}(\rho v\phi) + \frac{\partial}{\partial z}(\rho w\phi) = \frac{\partial}{\partial x}\left(\Gamma^\phi \frac{\partial\phi}{\partial x}\right) + \frac{\partial}{\partial y}\left(\Gamma^\phi \frac{\partial\phi}{\partial y}\right) + \frac{\partial}{\partial z}\left(\Gamma^\phi \frac{\partial\phi}{\partial z}\right) + S^\phi$$

Onde: **u**, **v** e **w** são as componentes da velocidade nas direções **x**, **y** e **z**. **ρ** é a densidade do fluido, Γ^ϕ é o coeficiente de transferência, ϕ é a variável do escoamento e S^ϕ é o termo fonte.

Na realidade, os códigos de CFD nunca resolvem equações diferenciais. Só resolvem as algébricas que, quando o número de volumes é grande o suficiente, possuem as mesmas implicações que as diferenciais.

Equação de conservação	ϕ	Γ^ϕ	S^ϕ
Continuidade	1	0	0
Momento em x	u	μ	$B_x + \frac{\partial}{\partial x} \left(\mu \frac{\partial u}{\partial x} - \frac{2}{3} \mu \nabla \cdot \vec{V} \right) + \frac{\partial}{\partial y} \left(\mu \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial z} \left(\mu \frac{\partial w}{\partial x} \right) - \frac{\partial P}{\partial x}$
Momento em y	v	μ	$B_y + \frac{\partial}{\partial y} \left(\mu \frac{\partial v}{\partial y} - \frac{2}{3} \mu \nabla \cdot \vec{V} \right) + \frac{\partial}{\partial x} \left(\mu \frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial z} \left(\mu \frac{\partial w}{\partial y} \right) - \frac{\partial P}{\partial y}$
Momento em z	w	μ	$B_z + \frac{\partial}{\partial z} \left(\mu \frac{\partial w}{\partial z} - \frac{2}{3} \mu \nabla \cdot \vec{V} \right) + \frac{\partial}{\partial x} \left(\mu \frac{\partial u}{\partial z} \right) + \frac{\partial}{\partial y} \left(\mu \frac{\partial v}{\partial z} \right) - \frac{\partial P}{\partial z}$
Energia	T	$\frac{k}{C_p}$	$\frac{1}{C_p} \frac{DP}{Dt} + \frac{\mu}{C_p} \Phi$
Massa de um componente i	C	ρD	0

$\frac{\partial}{\partial t}(\rho\phi)$	TERMO TEMPORAL
$\frac{\partial}{\partial x}(\rho u \phi) + \frac{\partial}{\partial y}(\rho v \phi) + \frac{\partial}{\partial z}(\rho w \phi)$	TERMO CONVECTIVO
$\frac{\partial}{\partial x}\left(\Gamma^\phi \frac{\partial \phi}{\partial x}\right) + \frac{\partial}{\partial y}\left(\Gamma^\phi \frac{\partial \phi}{\partial y}\right) + \frac{\partial}{\partial z}\left(\Gamma^\phi \frac{\partial \phi}{\partial z}\right)$	TERMO DIFUSIVO
S^ϕ	TERMO FONTE

Laval

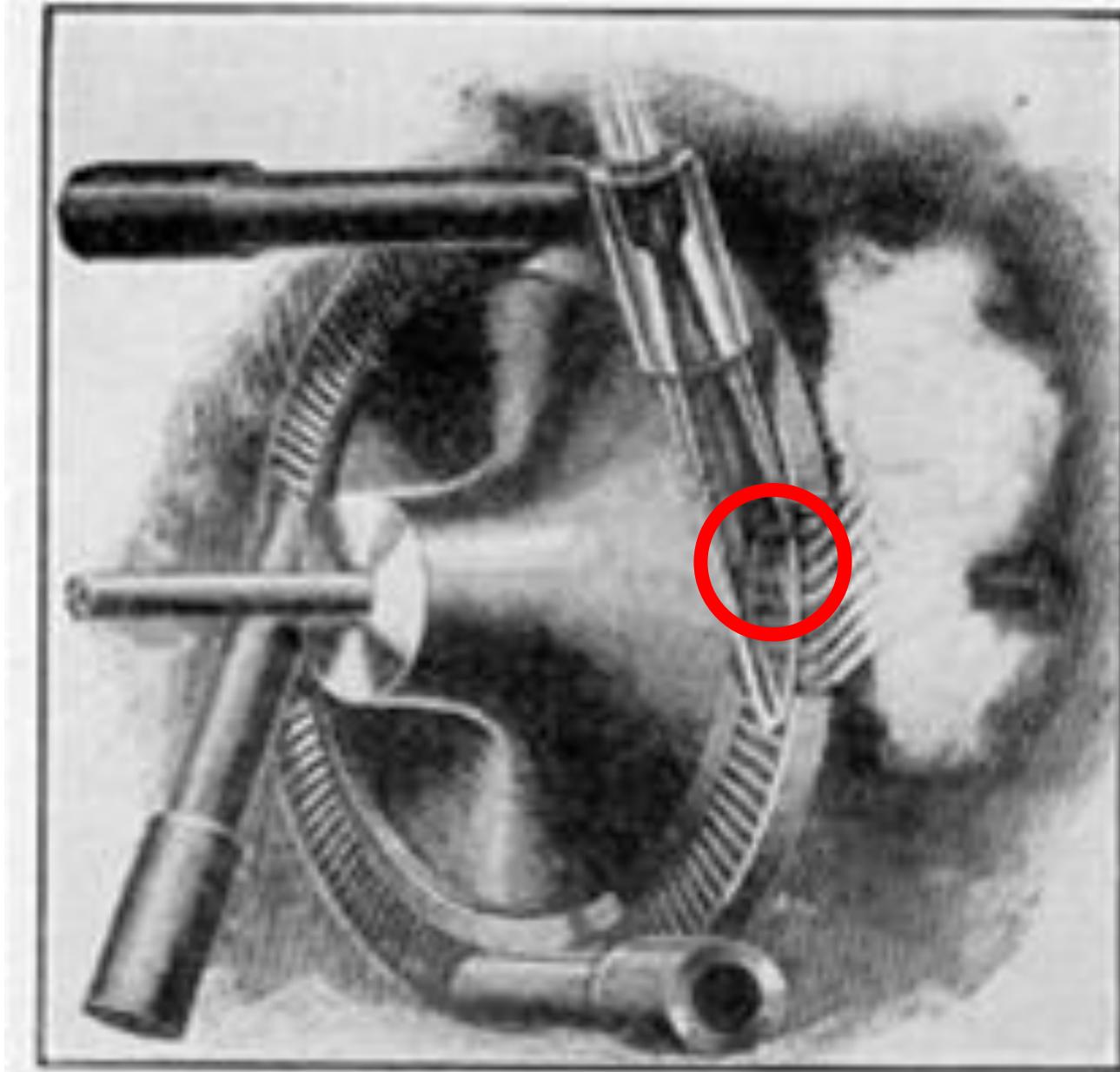
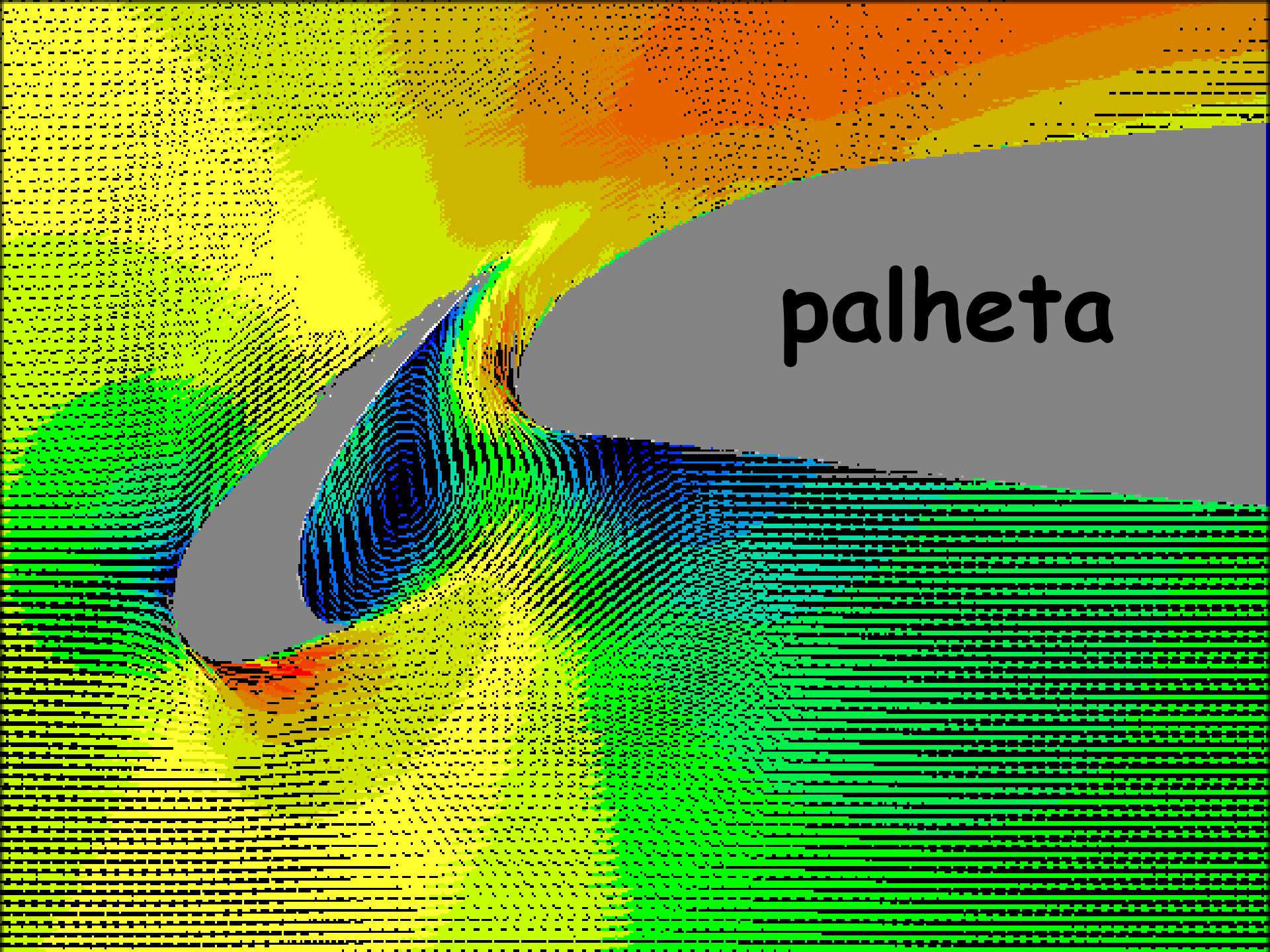
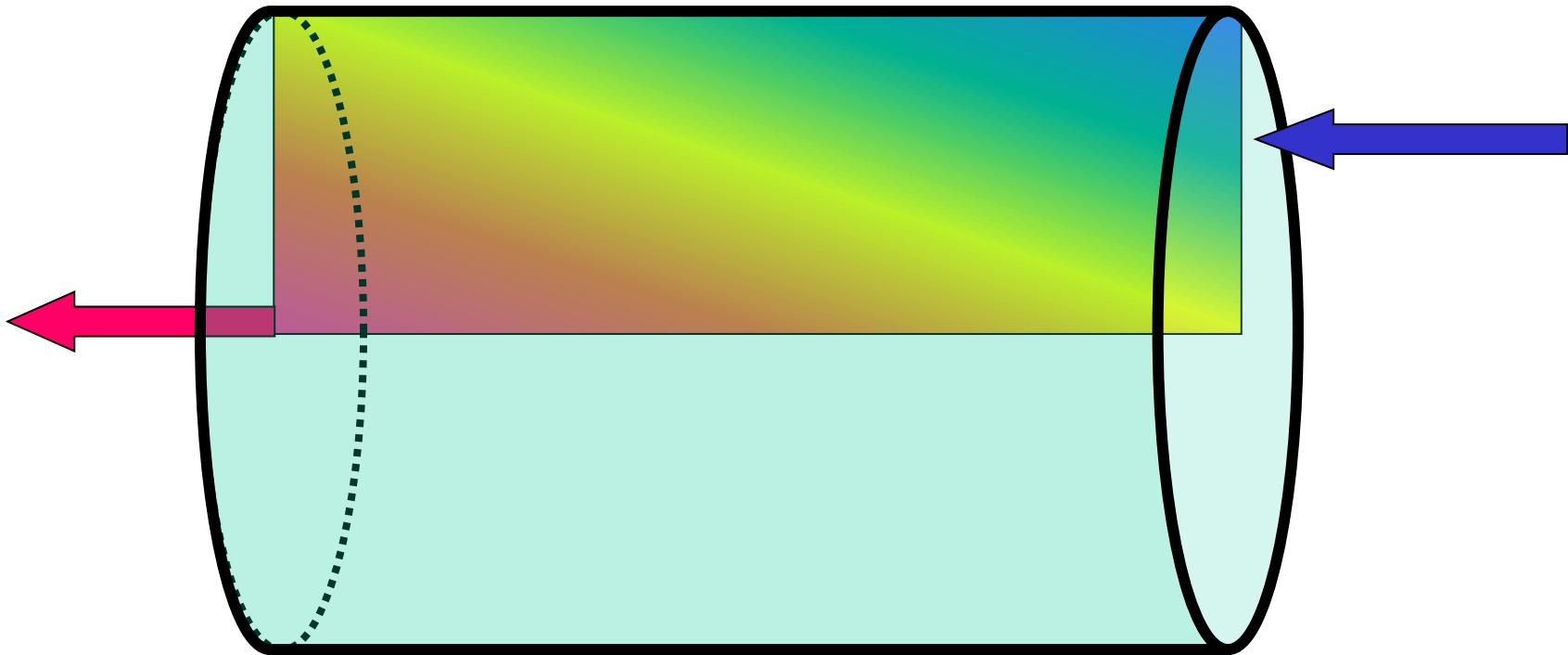


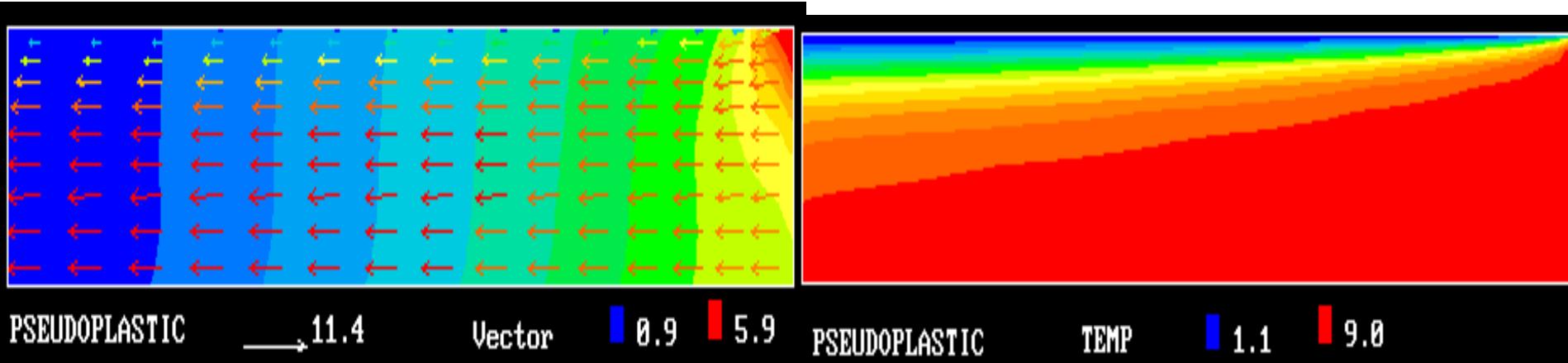
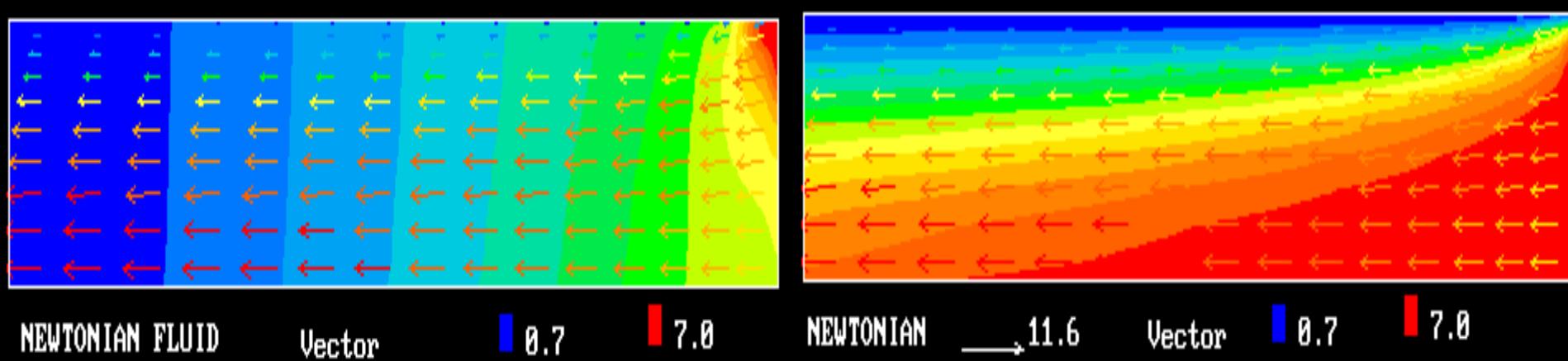
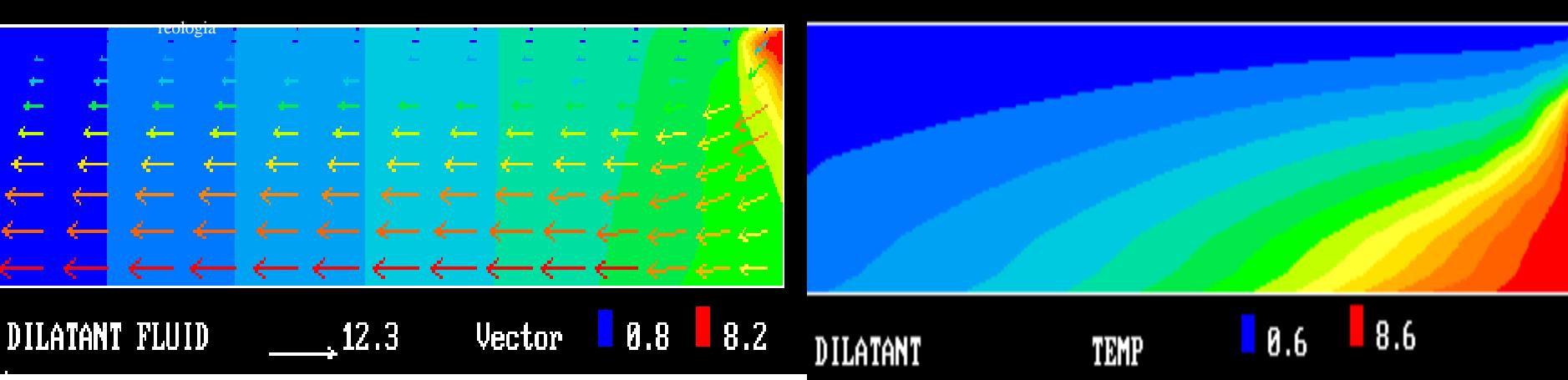
Fig. 4. Dr. de Laval's Turbloc.



palheta

Caso 0

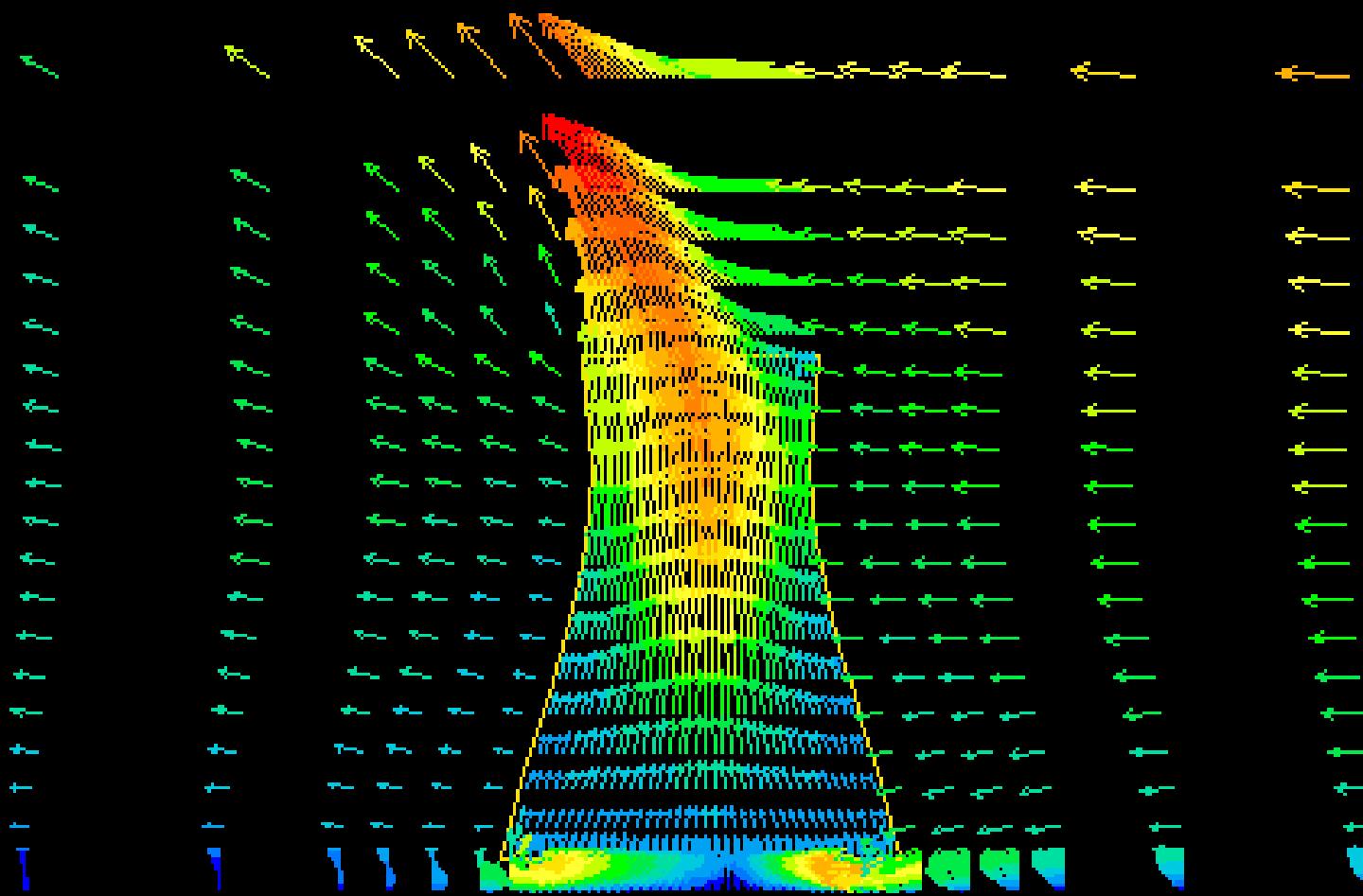




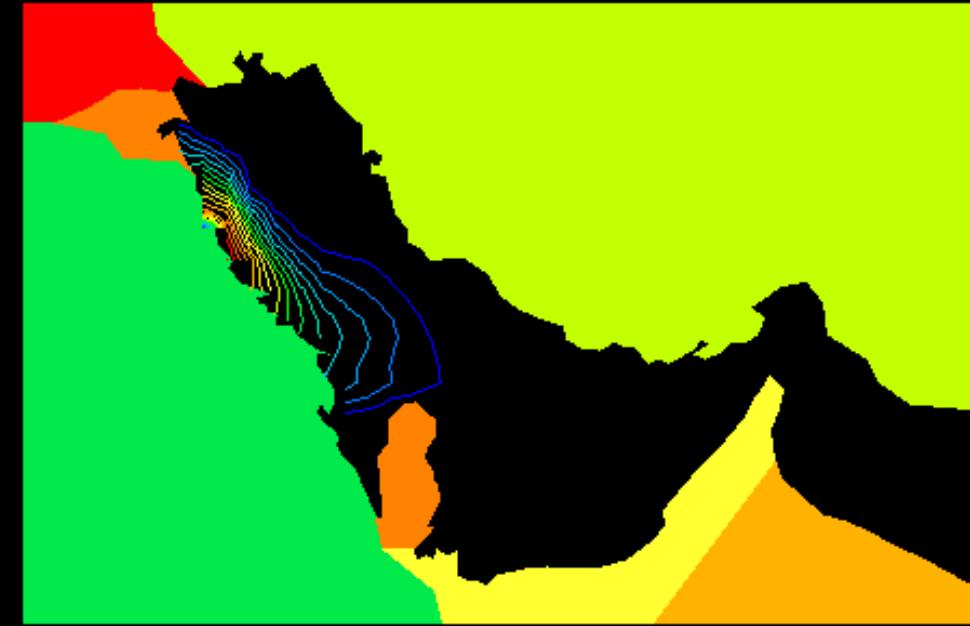
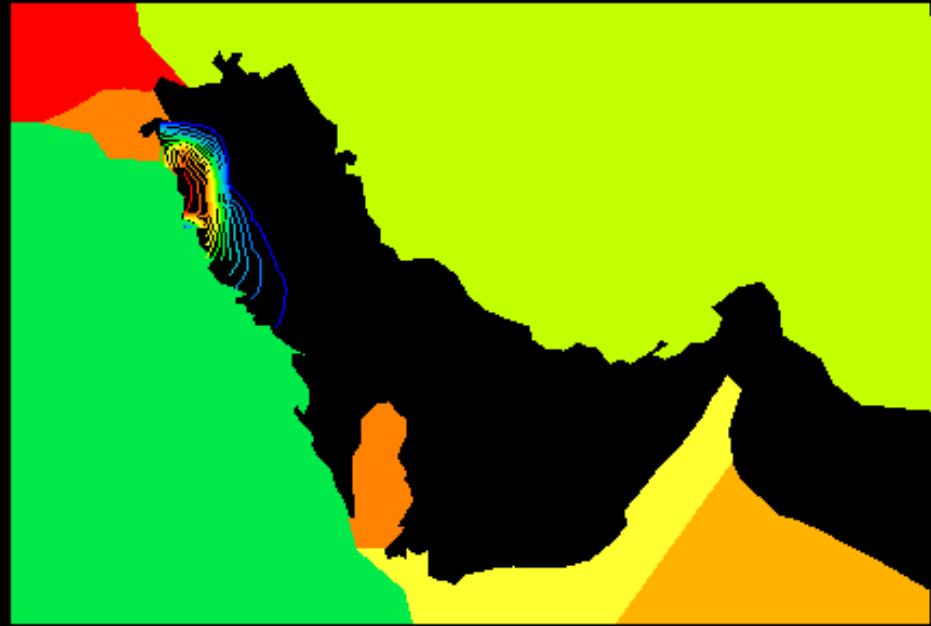
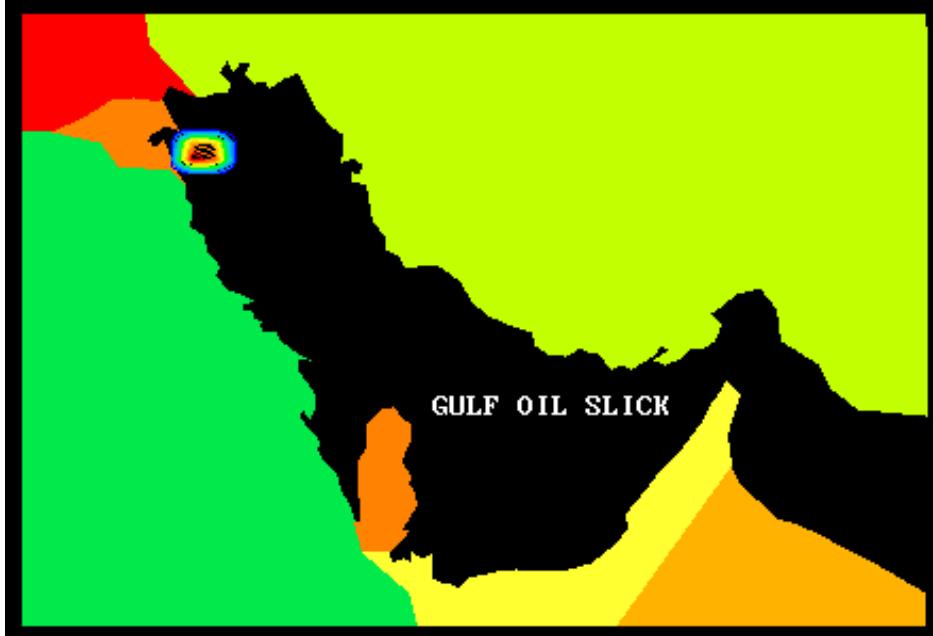


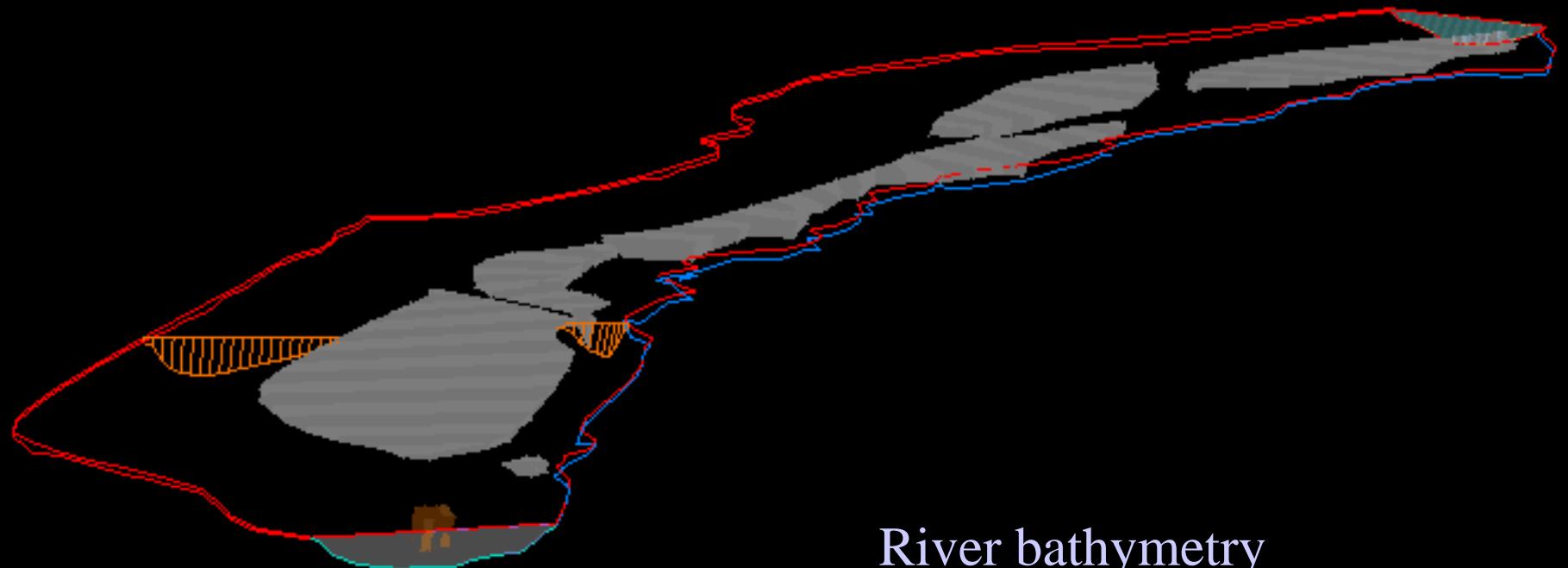
COOLING TOWER SIMULATION

←--- wind 1.5m/s at 10m

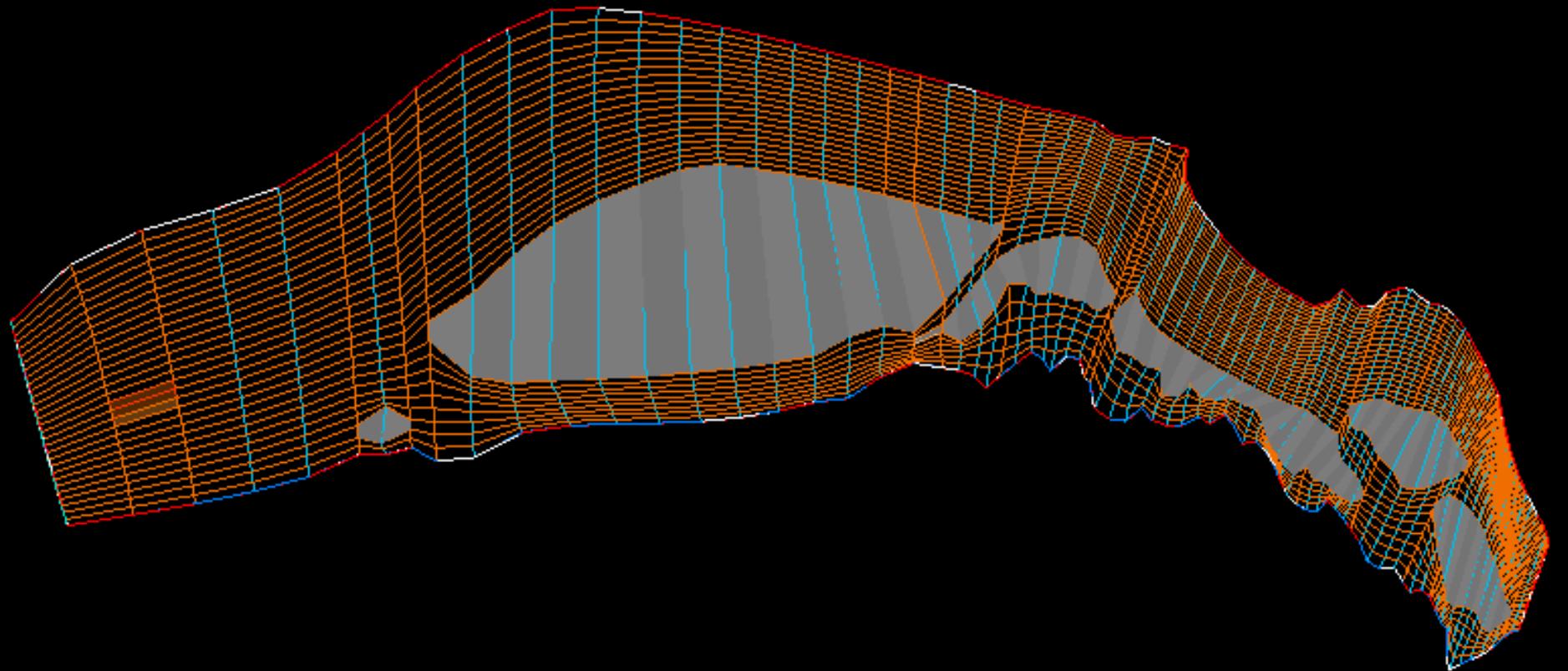


→ : 7.31 m/s. Min : 1.9599E-02 Max : 5.8843E+00

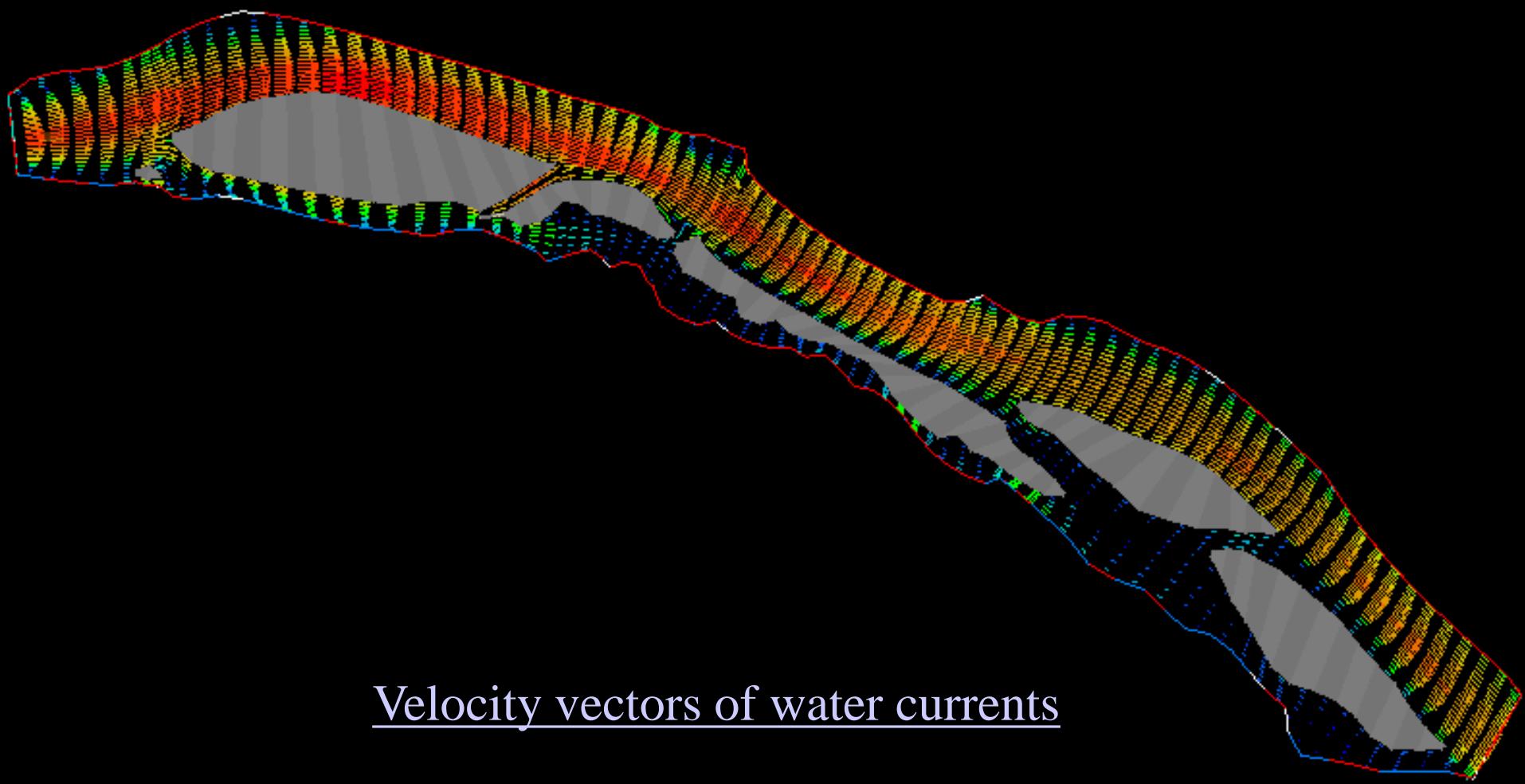




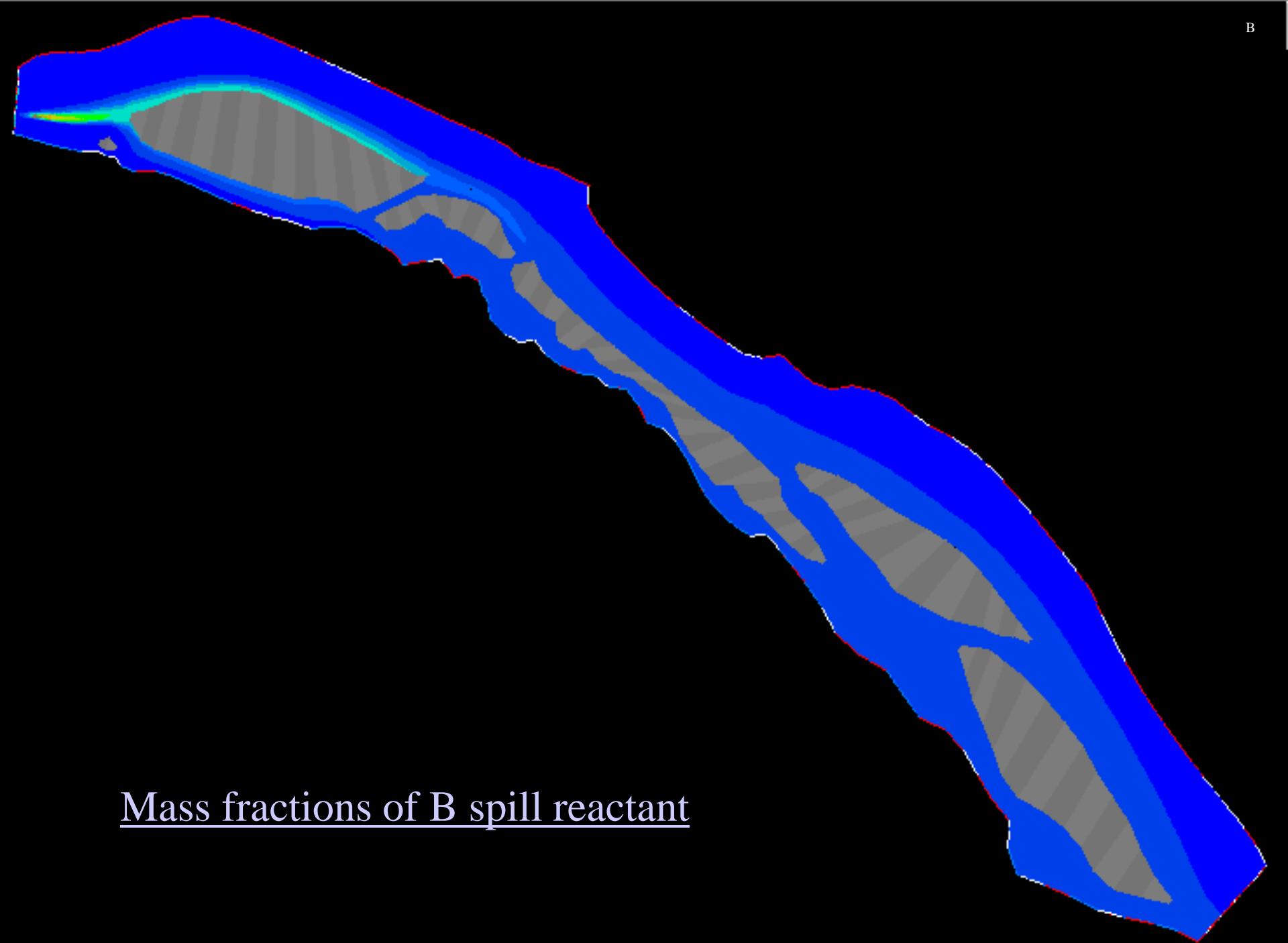
River bathymetry



Plan view on computational grid

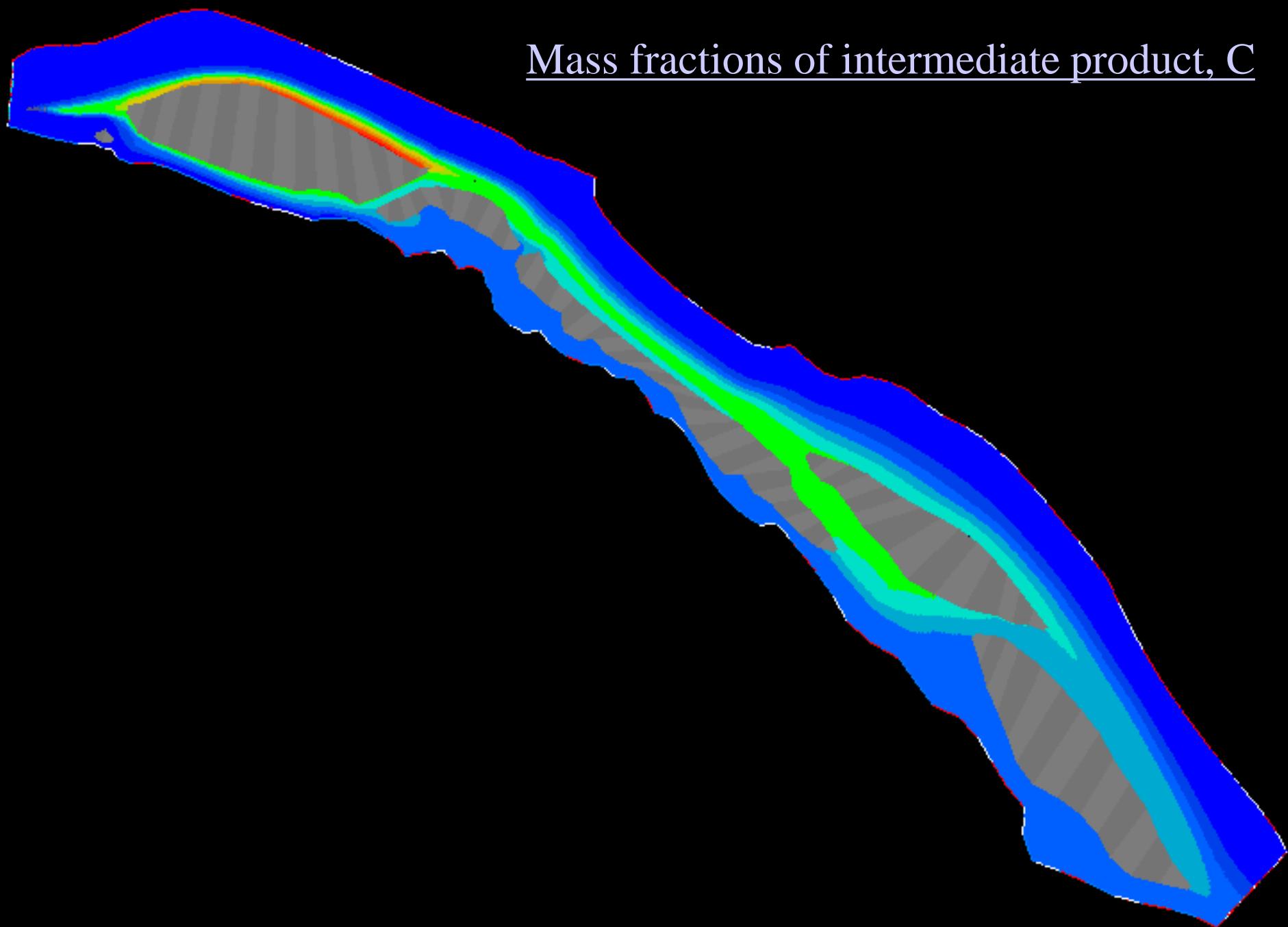


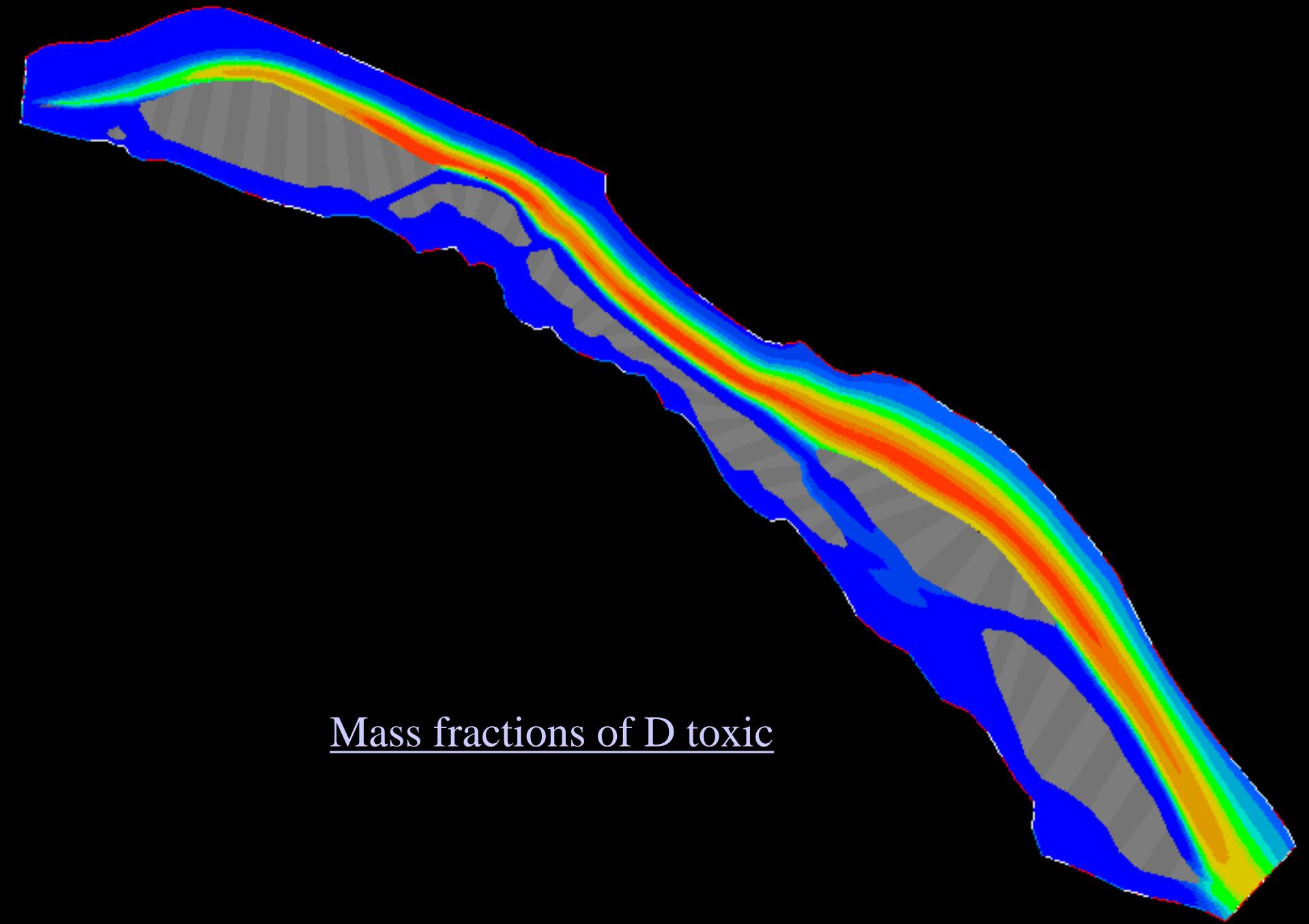
Velocity vectors of water currents



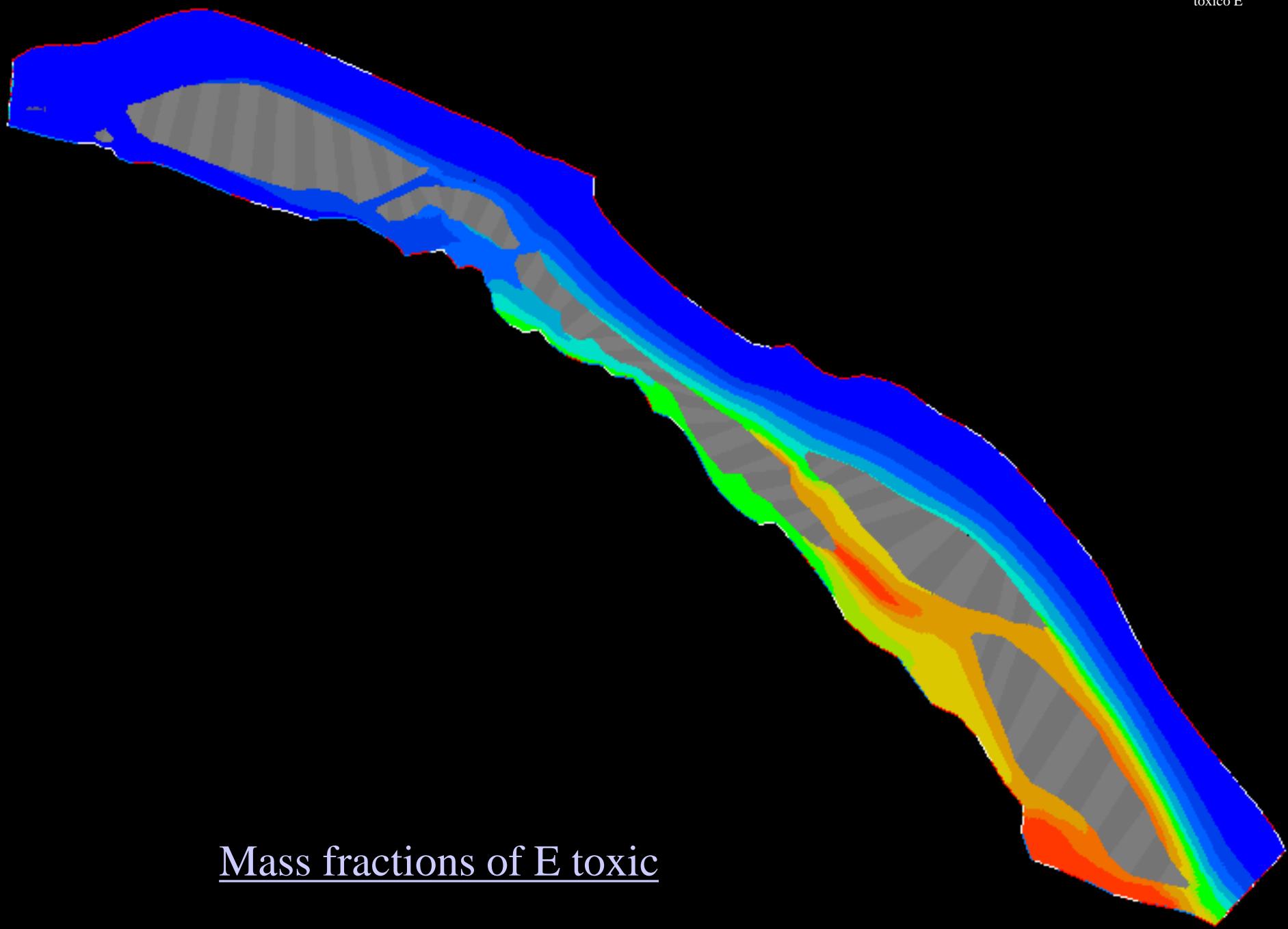
Mass fractions of B spill reactant

Mass fractions of intermediate product, C



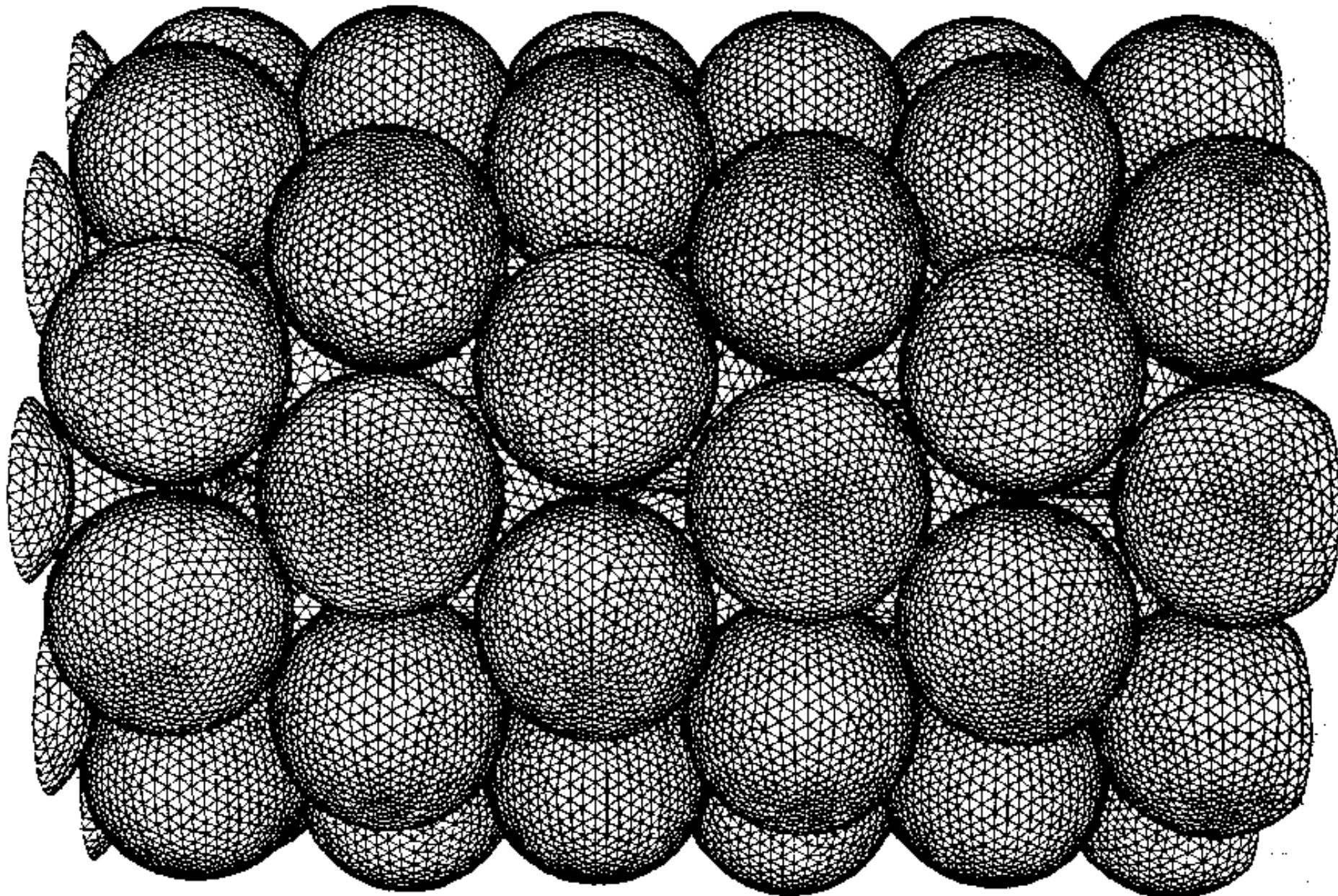


Mass fractions of E toxic



queimador





Exemplo: combustão de óleo

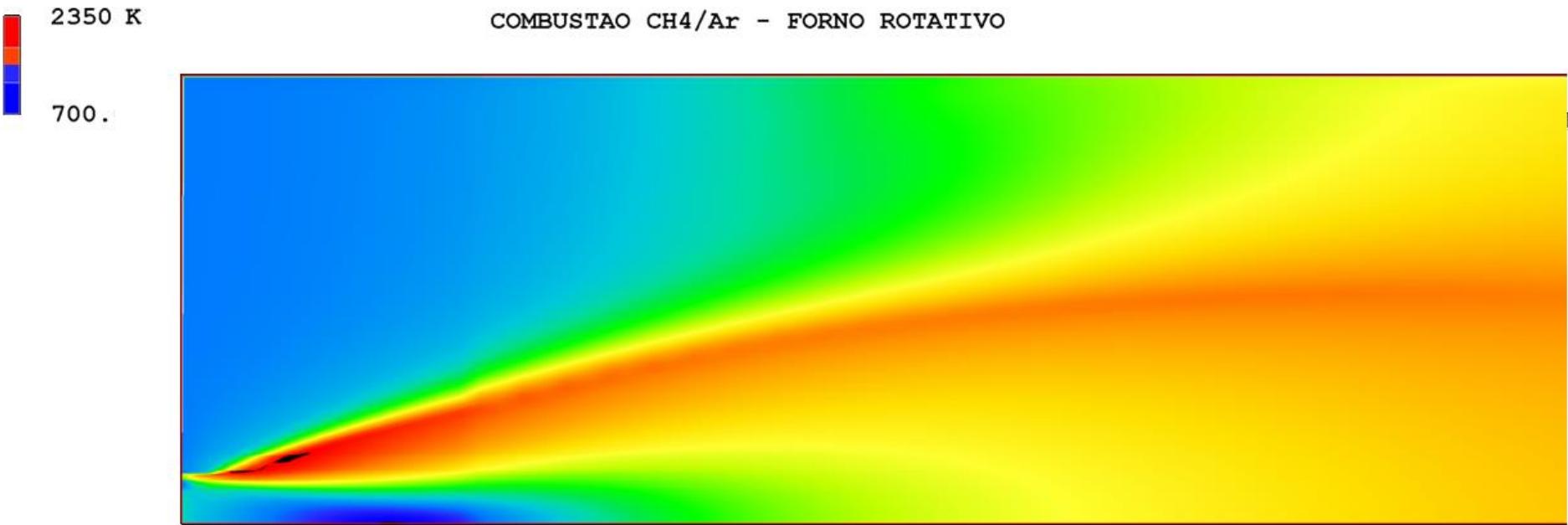


teste

A photograph of a large industrial furnace or kiln. The structure is made of dark, textured material, possibly refractory brick. A massive, turbulent plume of orange and yellow flames erupts from a circular opening on the left side. The fire is extremely intense, filling the right half of the frame with bright light and smoke. The overall atmosphere is one of raw energy and heat.

chama

simulação



MODELOS DE RADIAÇÃO NO PHOENICS

- Modelo *composite-flux (6-flux)* de Schuster & Hamaker, conforme formulado por Spalding (1980).
- Modelo *composite-radiosity* de Spalding (1994); este modelo é similar ao modelo P-1 harmônico-esférico de Ozisik (1973).
- Modelo de difusão de *Rosseland* (1936).
- Modelo IMMERSOL.
- Modelo de radiação superfície-para-superfície.

perfil



eddy break up

comburente

jato



fuel

mistura

reação

comburente

TURBULÊNCIA de mistura

6. Turbulent Reaction-Rate Sources

6.1 Eddy-breakup

For turbulent flows, the second form (see section 2.6 above) of the eddy-breakup reaction-rate is provided. The resulting source per unit phase mass is:

$$S_{\text{mfp}} = -a \times \min\{ m_{fu}, m_{ox} \times s \} \times \frac{EP}{KE}$$

This is activated by:

PATCH (CHSO, PHASEM, IF,IL, JF,JL, KF,KL, TF,TL)
COVAL (CHSO, FUEL, GRND9, GRND9)

The reaction rate constant, a and the stoichiometric ratio,s, are passed via

$$\text{CHSOB} = a \quad \text{CHSOA} = 1 / (1 + s) [= \text{fstoic}]$$

The rate controlling parameter CHSOB is commonly set to unity.

Rodando o
CFD

PFR

n = 1

$$r = k c_A$$

$$\frac{\tilde{c}_s}{\tilde{c}_e} = e^{-\frac{Lk}{v}}$$

L=1m; v=1m/s; ce=1gmole/m3; k = 1m3/gmole/s -> cs/ce=0,37



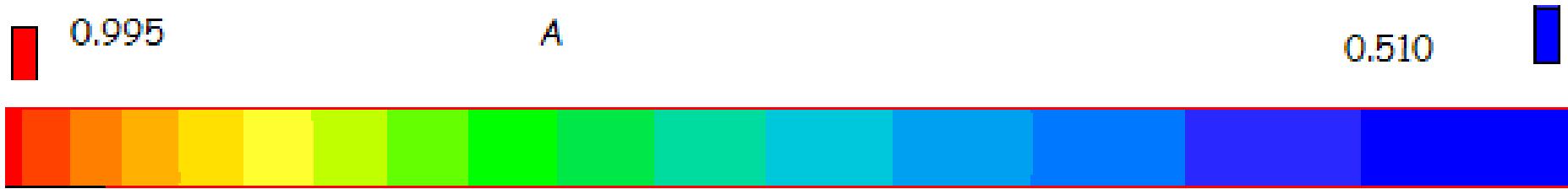
PFR L=1 v=1 ce=1 k= + 1

PFR $n = 2$

$$r = k c_A^2$$

$$\frac{\tilde{c}_s}{\tilde{c}_e} = \frac{1}{\frac{Lk\tilde{c}_e}{v} + 1}$$

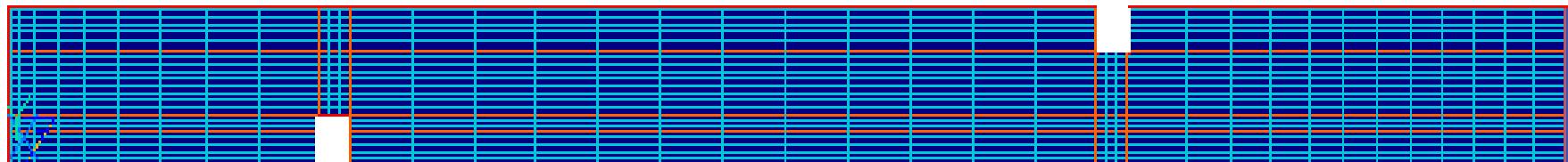
$L=1\text{m}$; $v=1\text{m/s}$; $ce=1\text{gmole/m}^3$; $k = 1\text{m}^3/\text{gmole/s}$ -> $cs/ce=0,50$



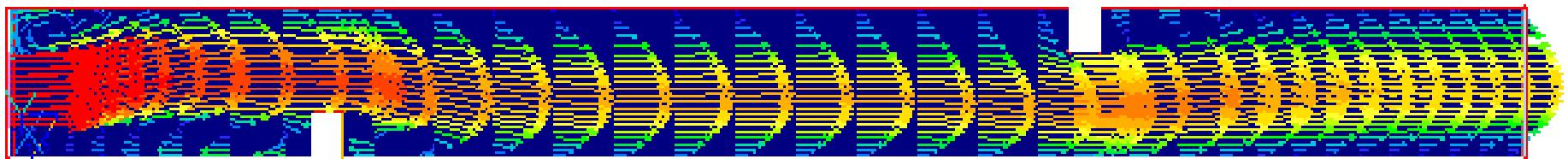
PFR $L=1$ $v=1$ $ce=1$ $k=-1$

CFD

malha



velocidades

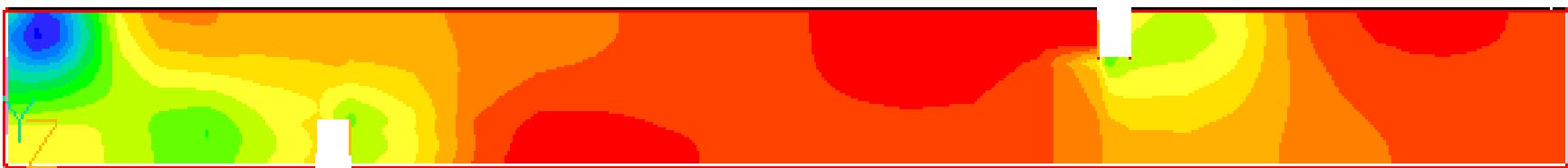


Pa

pressão

-0.37

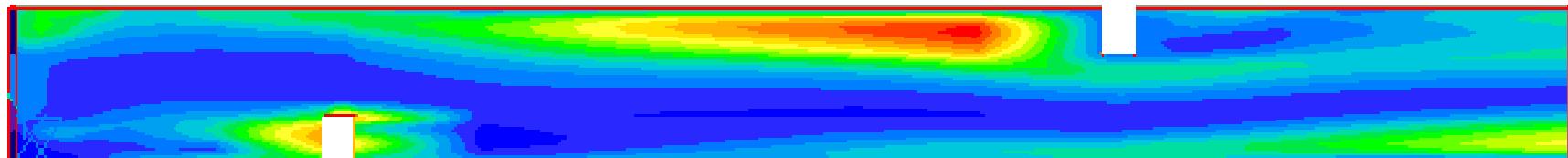
0.027



2.2E-6

viscosidade turbulenta

8.8E-5



33

EPKE

0



25

TEMPERATURA

80



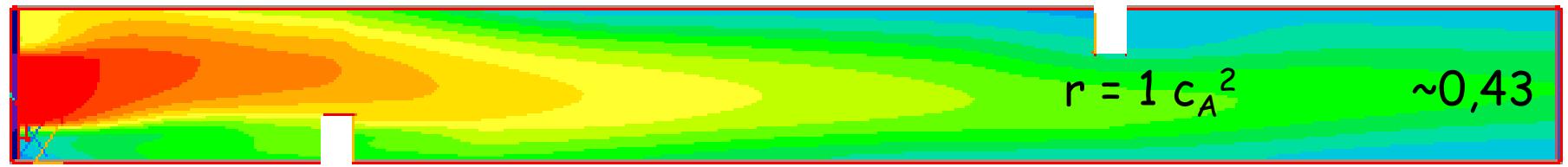
1

$k = -1$

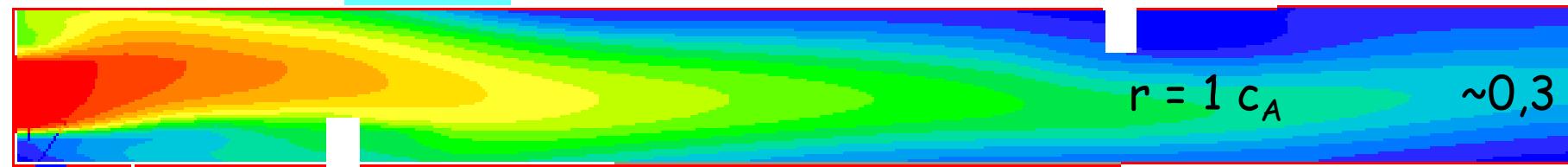
CONCENTRAÇÃO

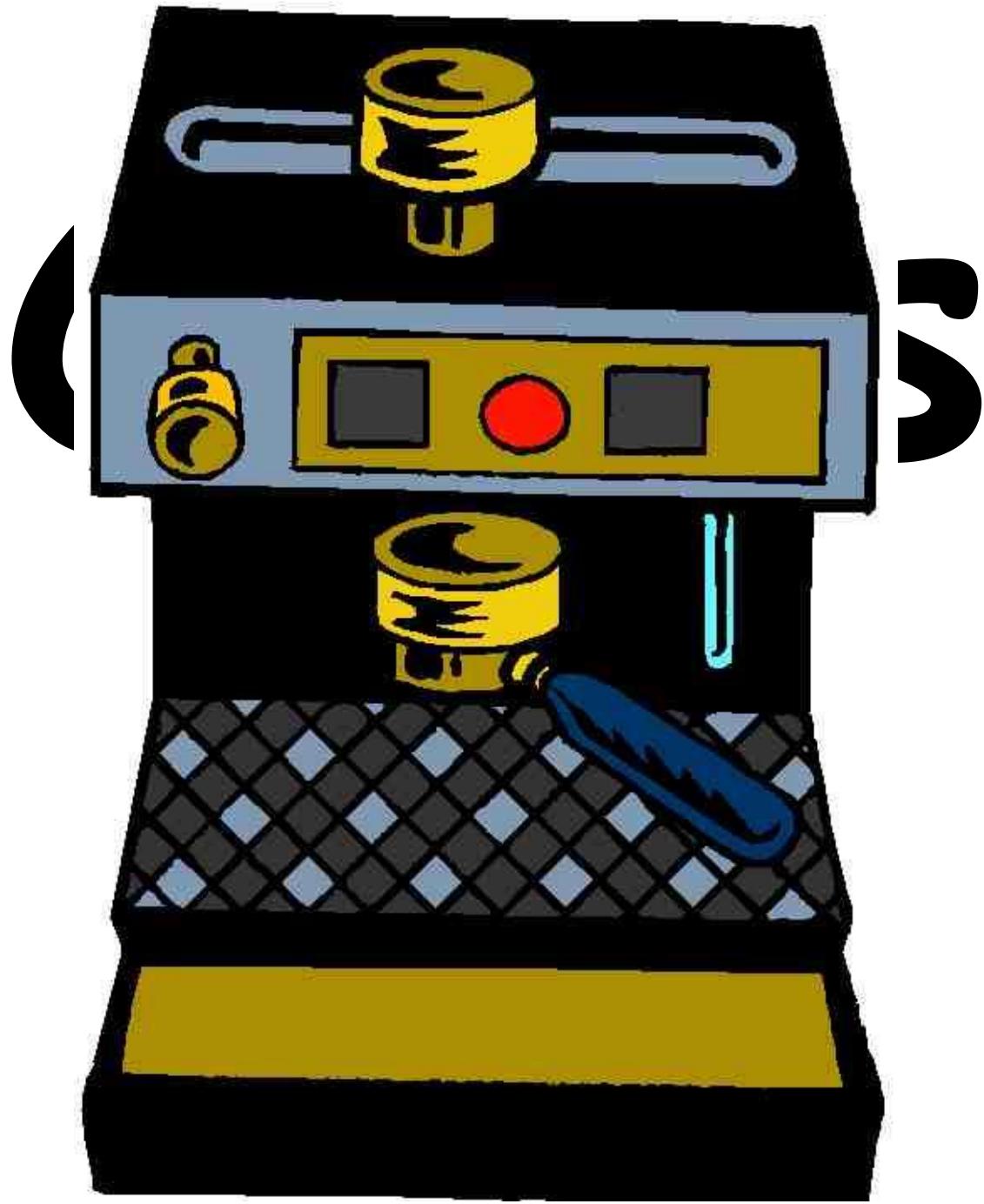
0

1

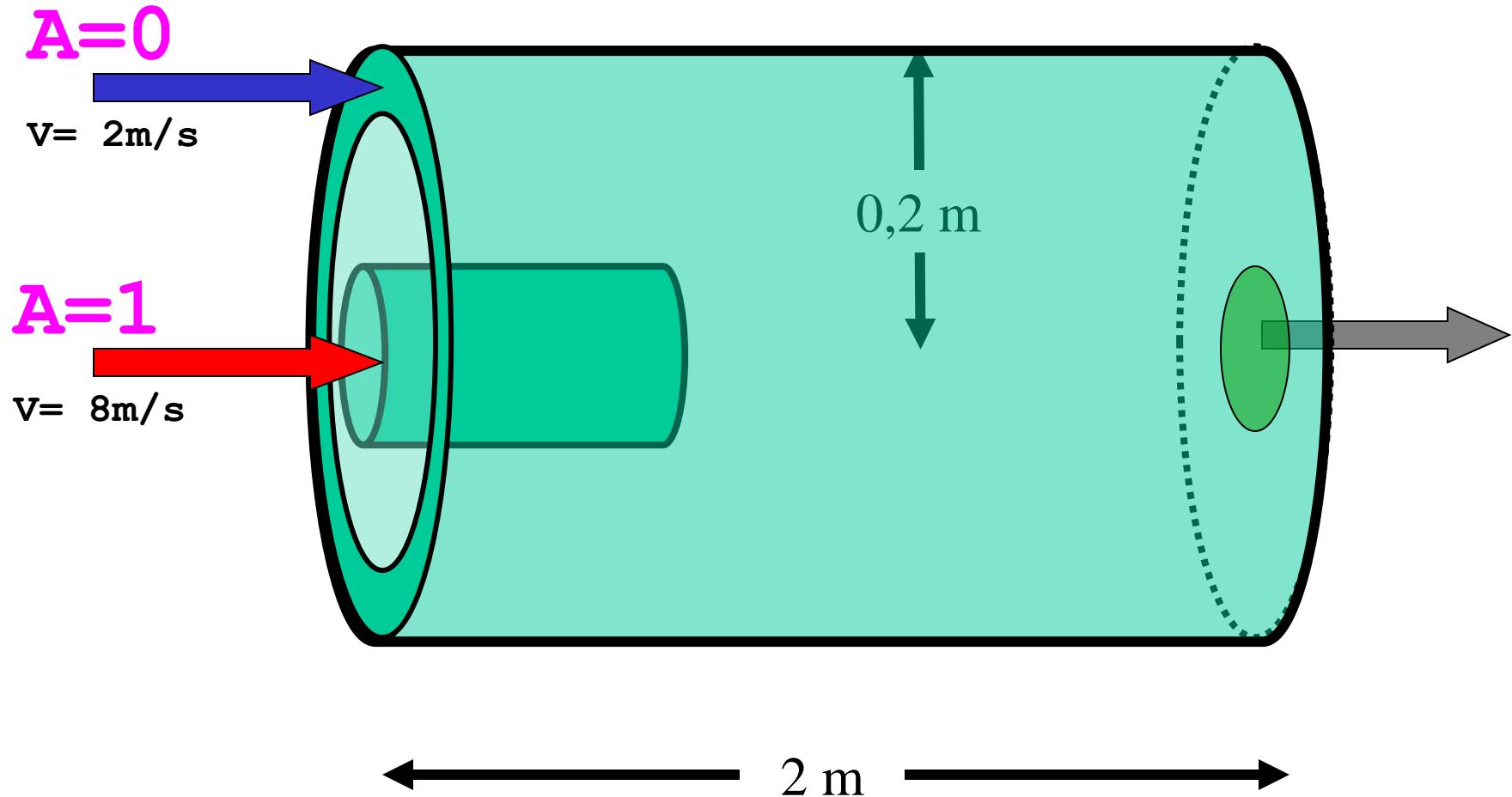


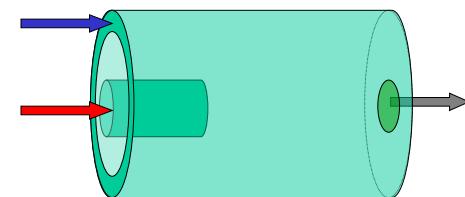
$k = +1$





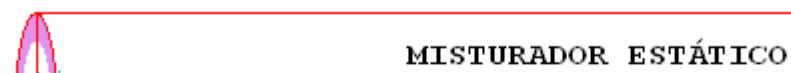
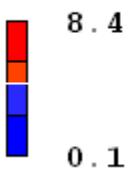
Caso 1



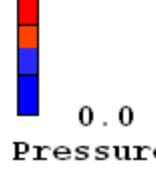


Caso 1

Velocity

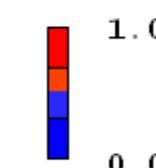


Pressure



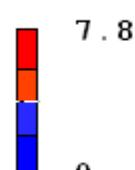
MISTURADOR ESTÁTICO

A



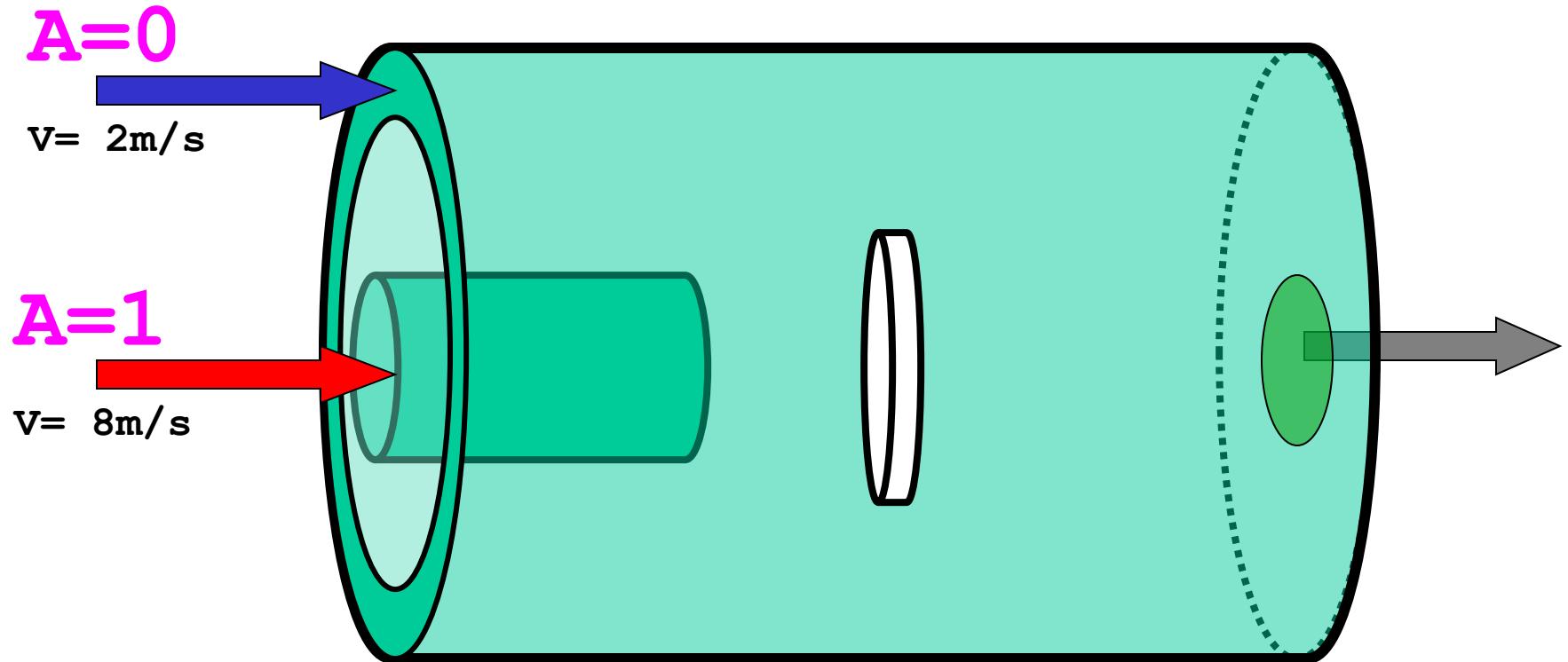
MISTURADOR ESTÁTICO

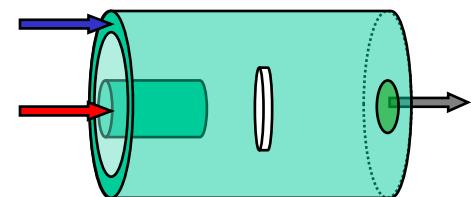
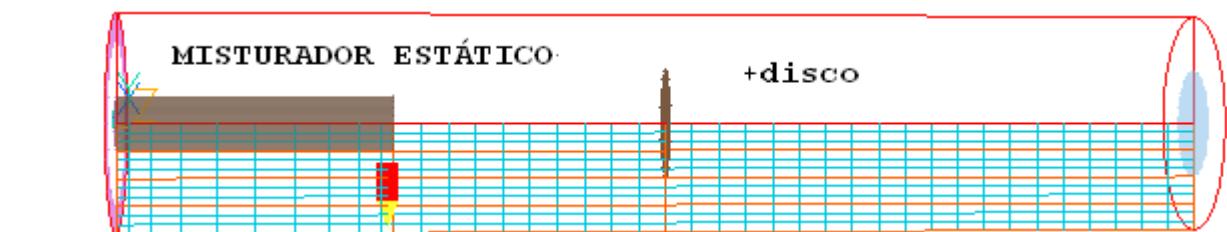
KE



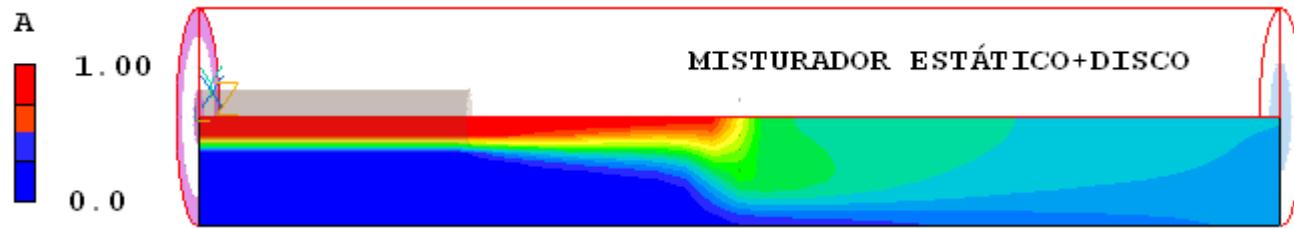
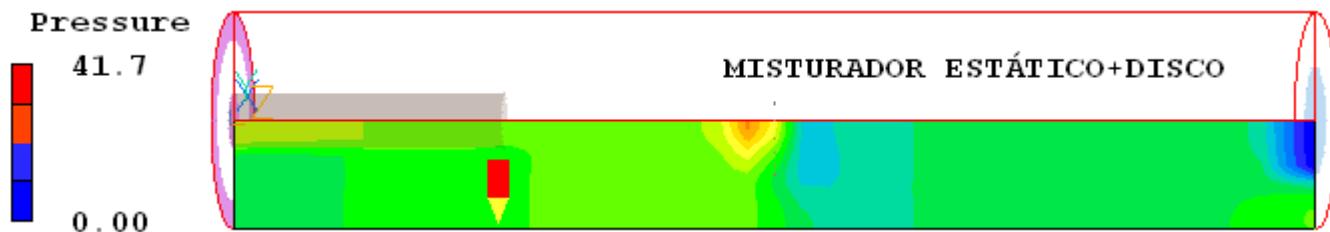
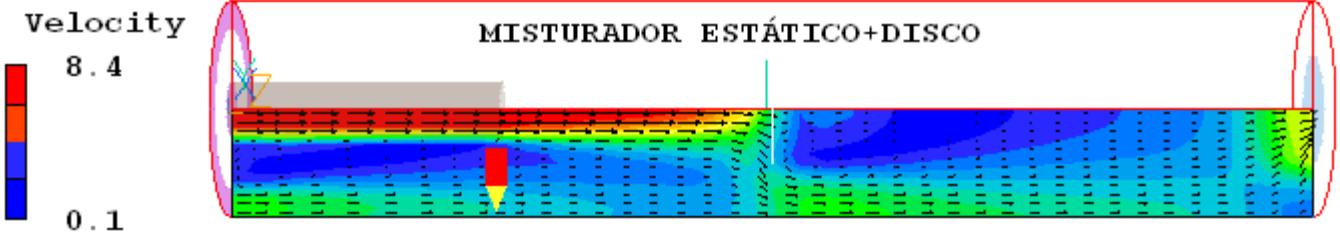
MISTURADOR ESTÁTICO

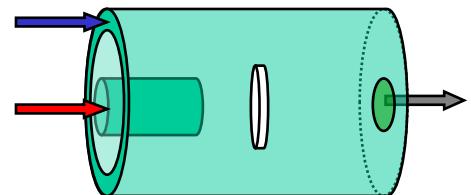
Caso 2



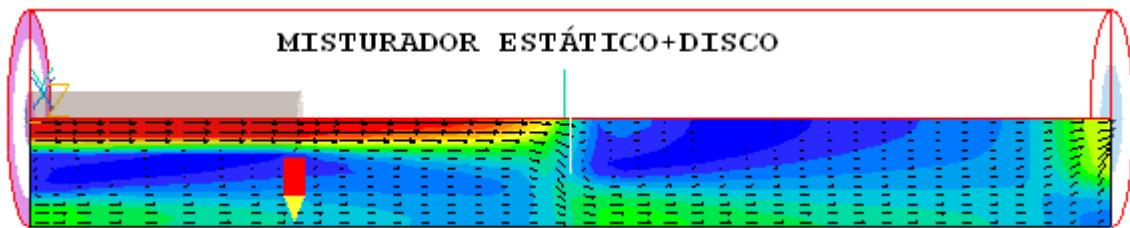
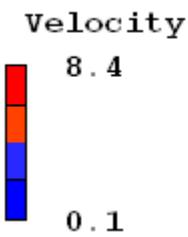


Caso 2

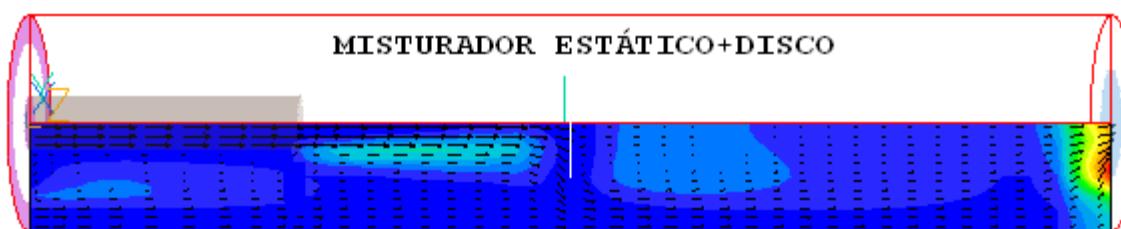
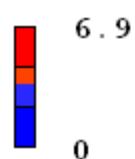




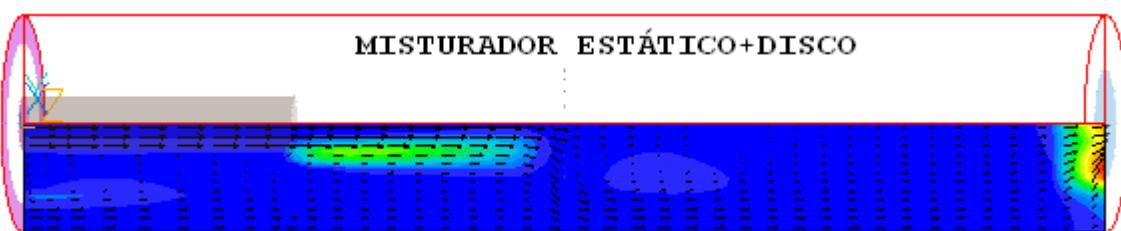
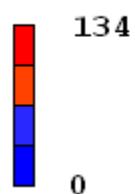
Caso 2



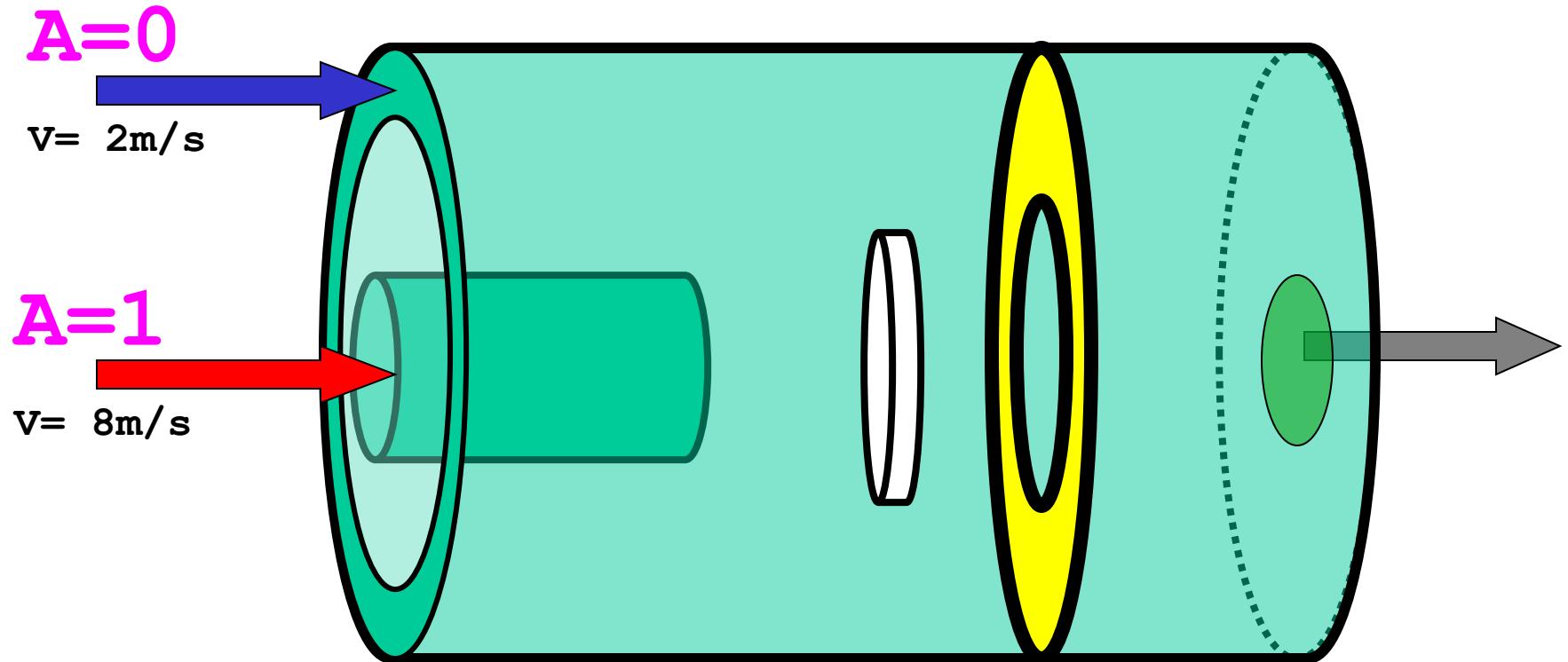
KE

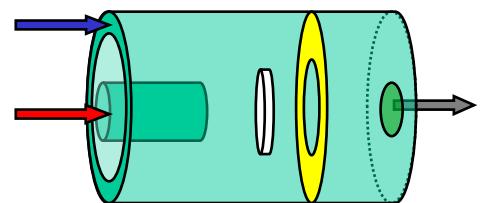


EP



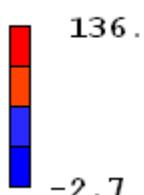
Caso 3



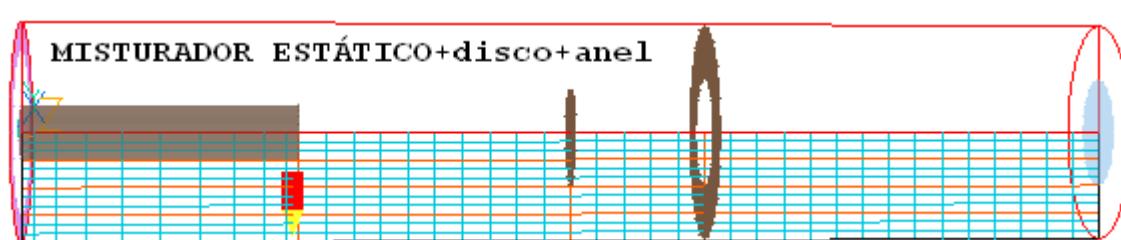


Caso 3

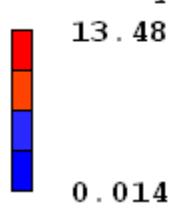
Pressure



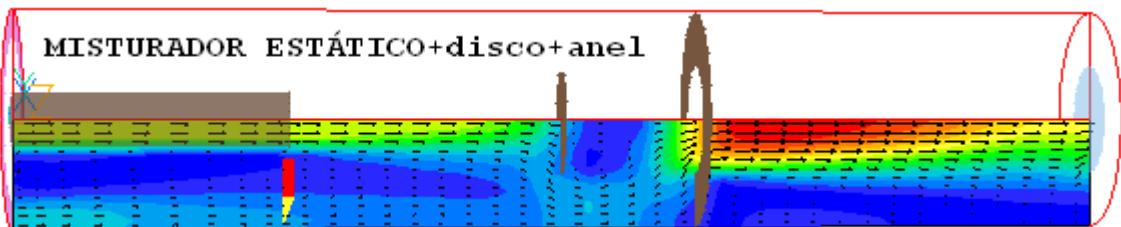
MISTURADOR ESTÁTICO+disco+anel



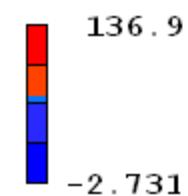
Velocity



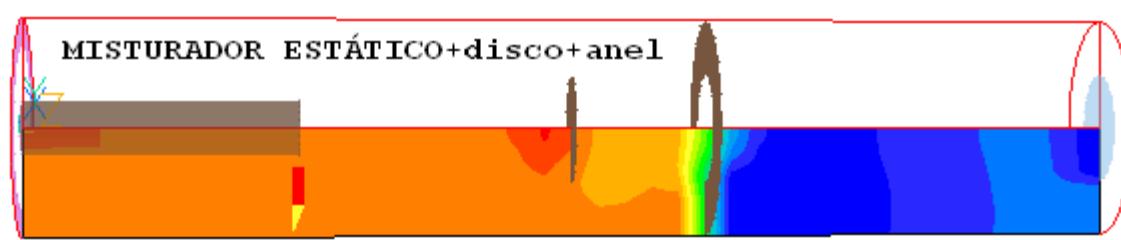
MISTURADOR ESTÁTICO+disco+anel



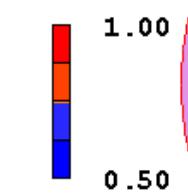
Pressure



MISTURADOR ESTÁTICO+disco+anel



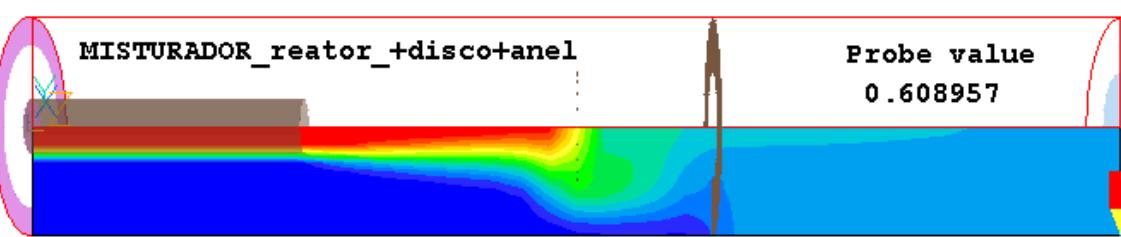
A



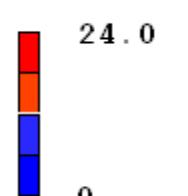
MISTURADOR_reator+_discotanel

Probe value

0.608957



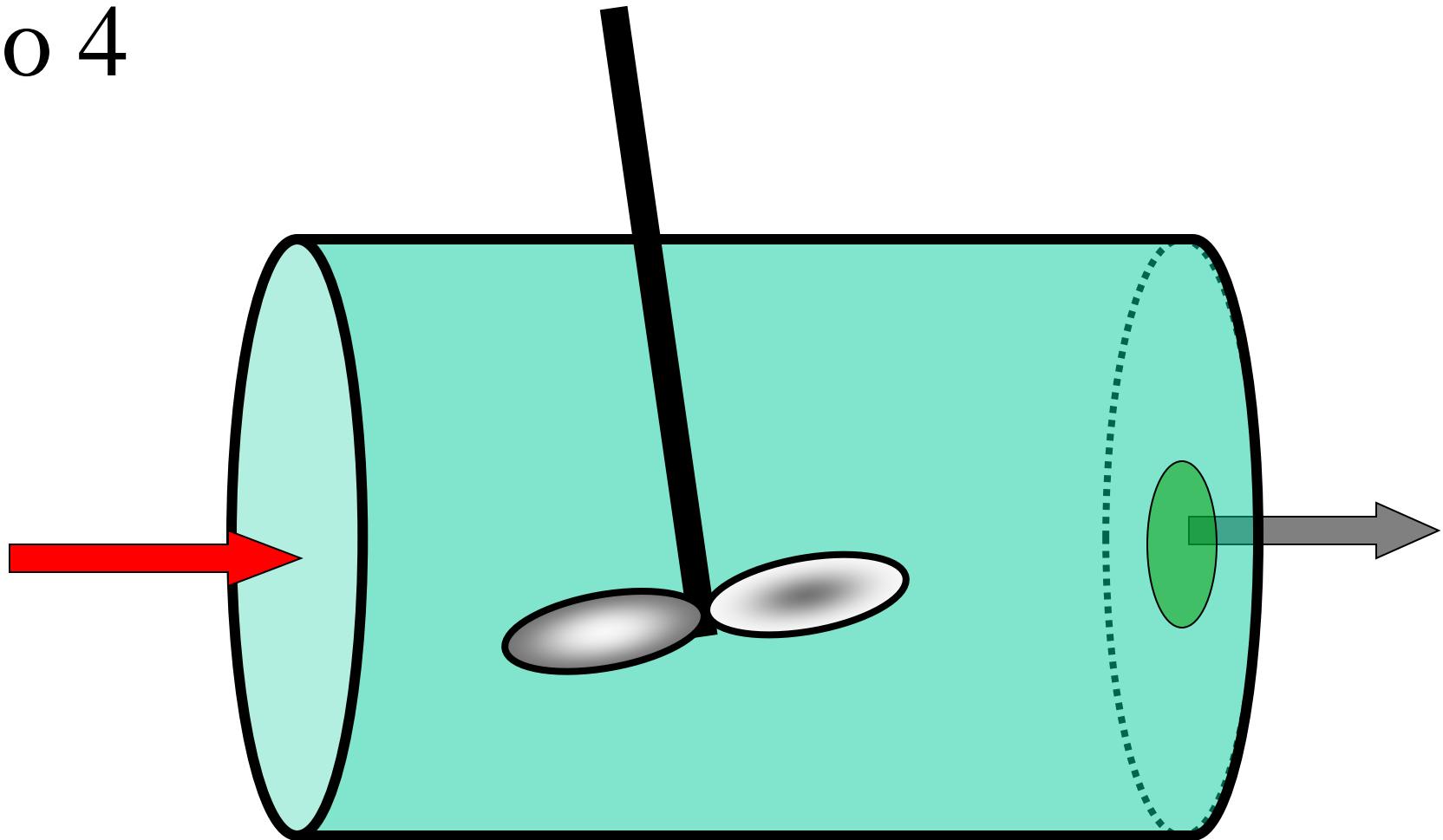
KE



MISTURADOR ESTÁTICO+disco+anel

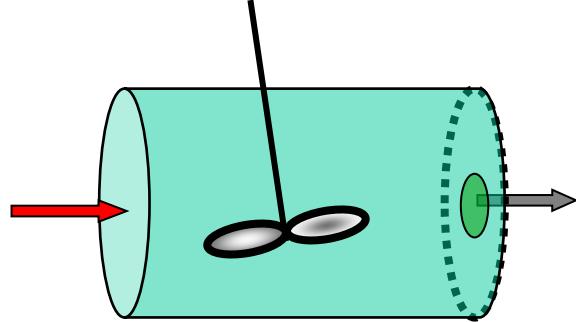


Caso 4

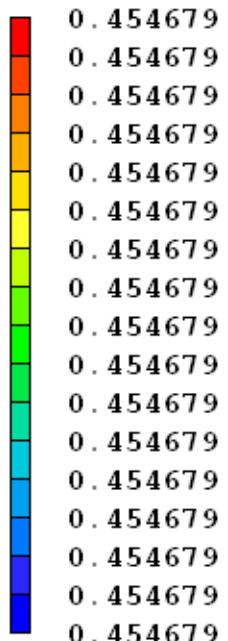


CFSTR

Caso 4



A



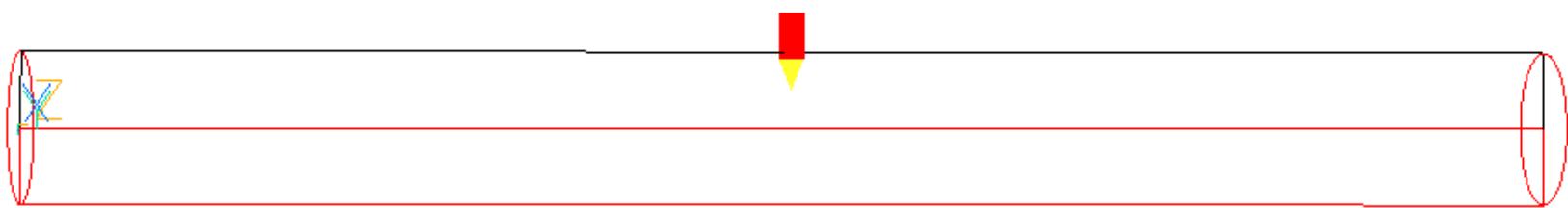
REATOR AR K=3 V=1 MISTURA

Probe value

0.454679

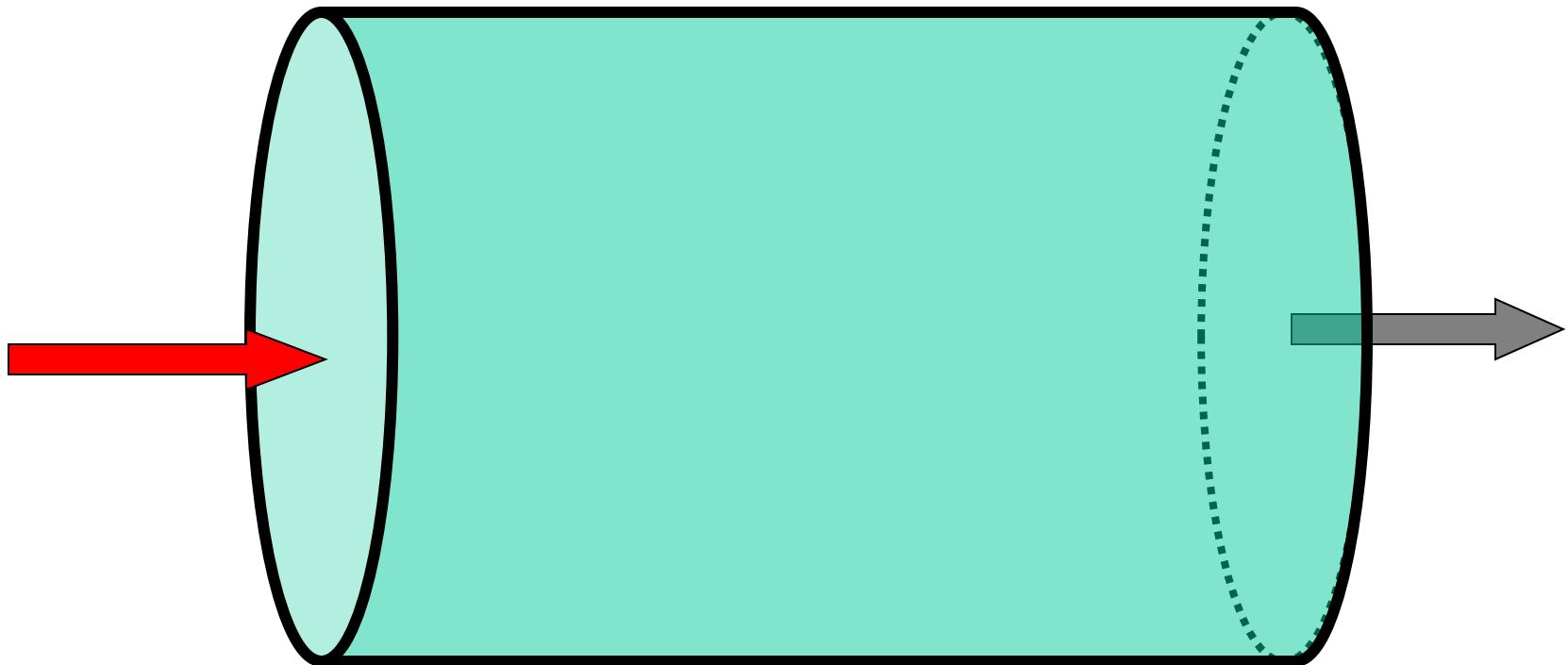
Average value

0.454679



CFSTR

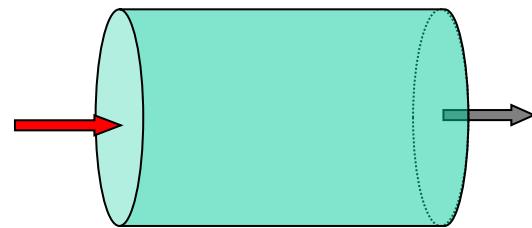
Caso 5



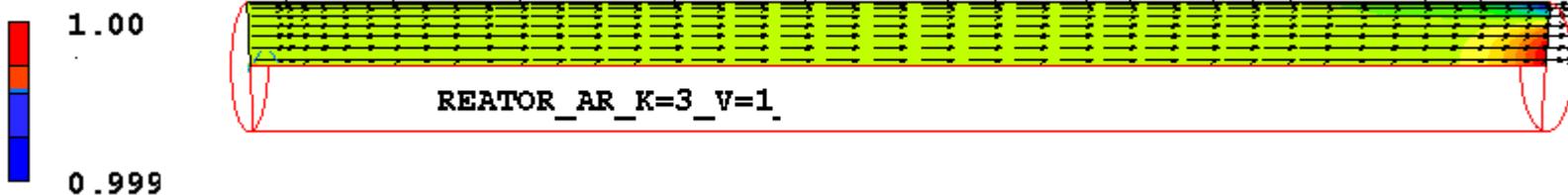
PFR

Caso 5

PFR



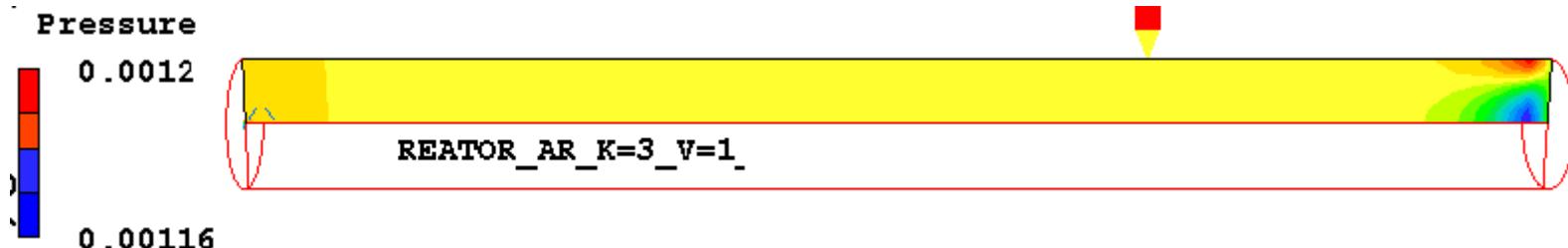
Velocity



A



Pressure



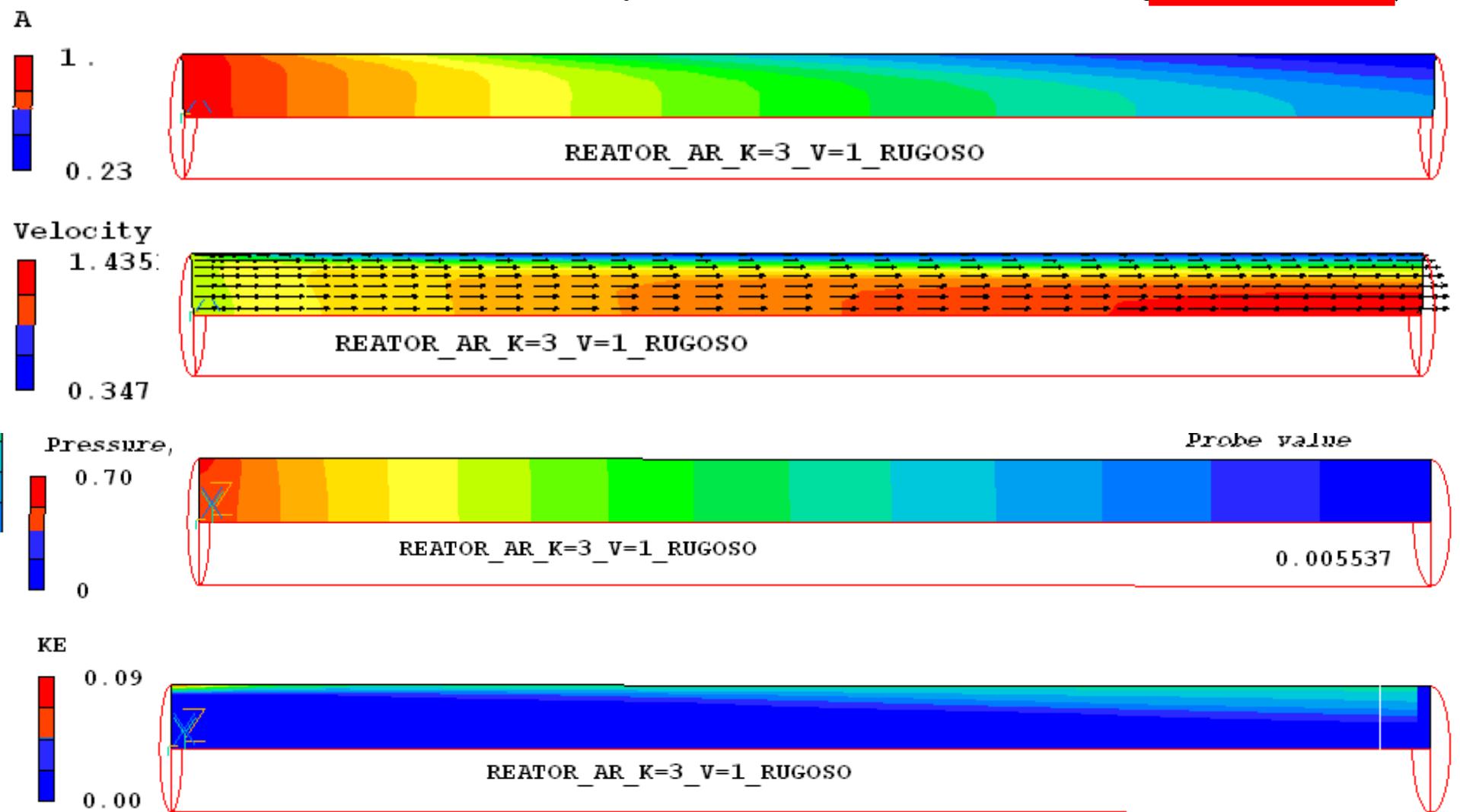
Caso 6



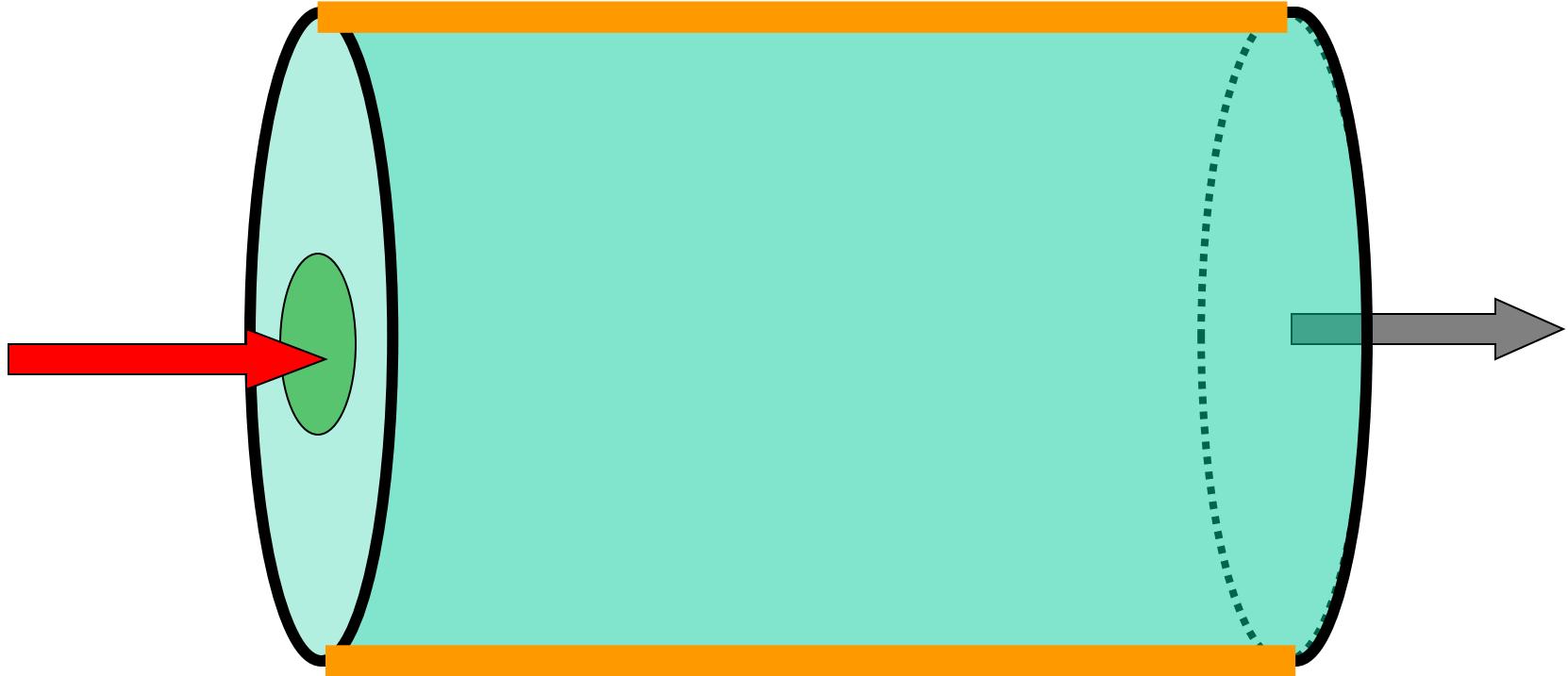
PFR com atrito na parede

Caso 6

PFR com atrito na parede



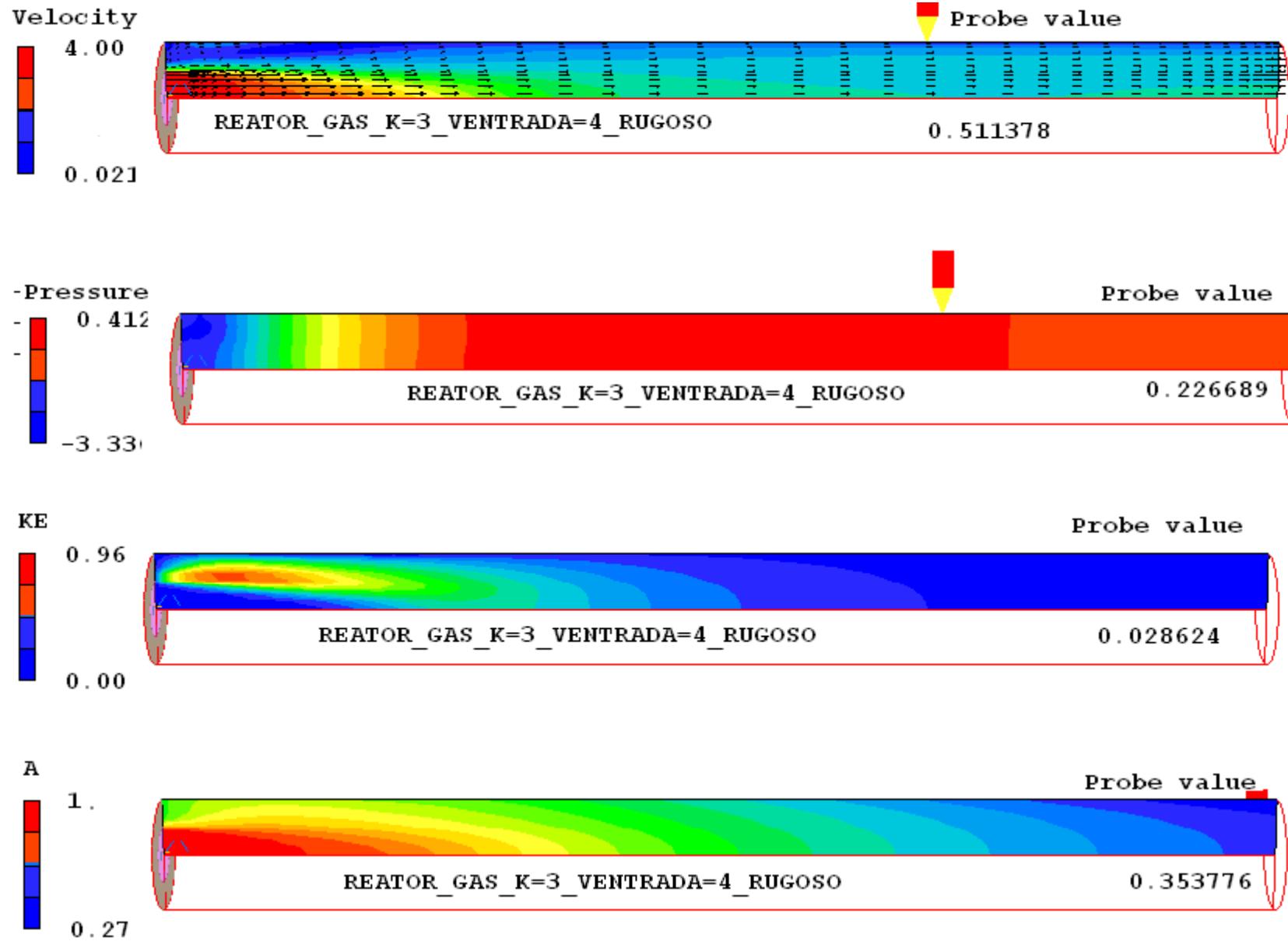
Caso 7



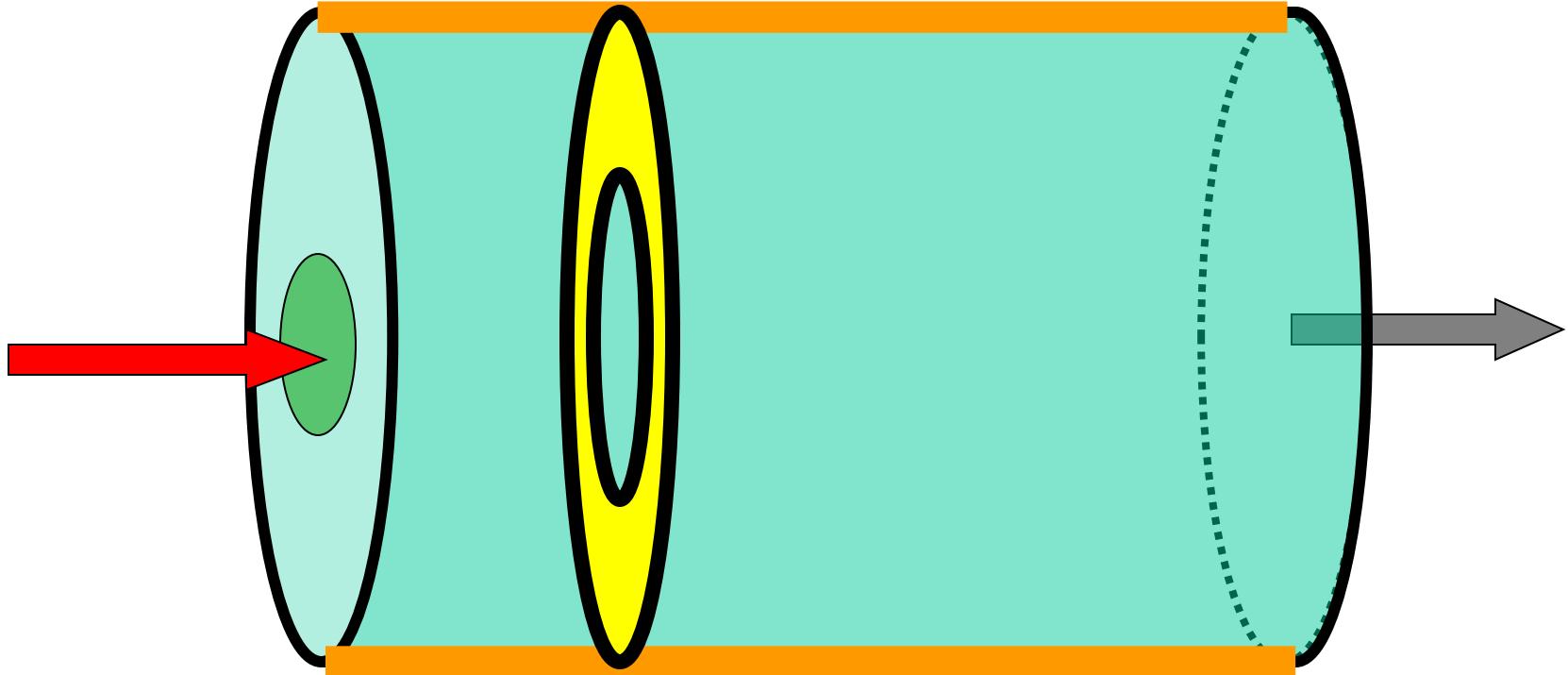
rugoso

C
a
s
o

7

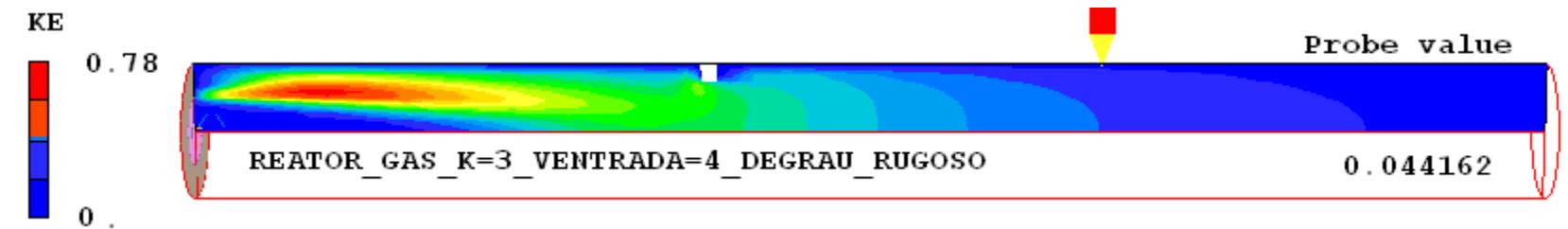
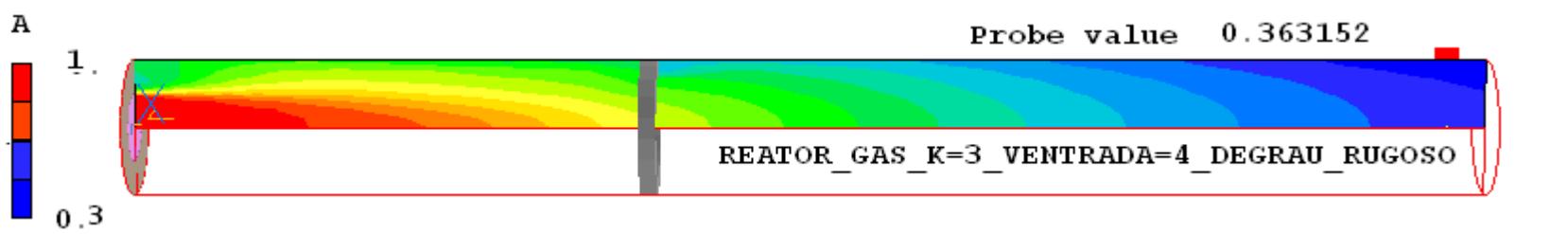
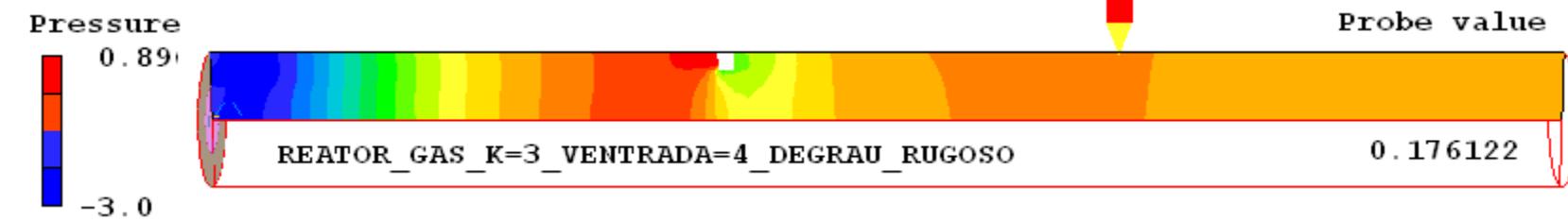
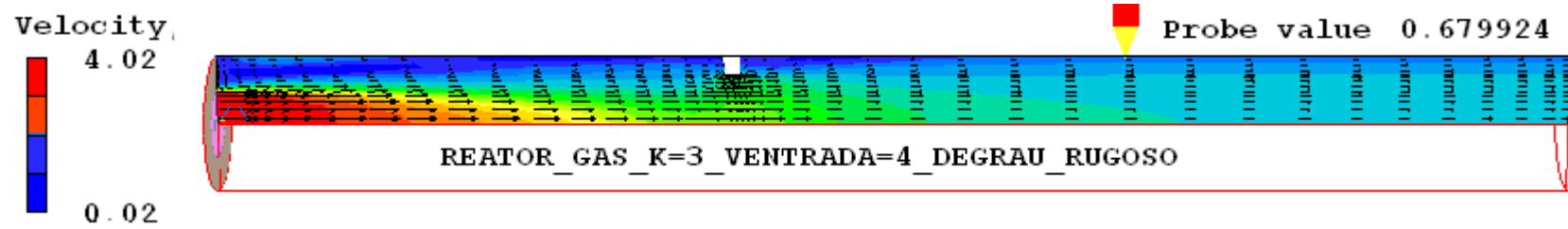


Caso 8

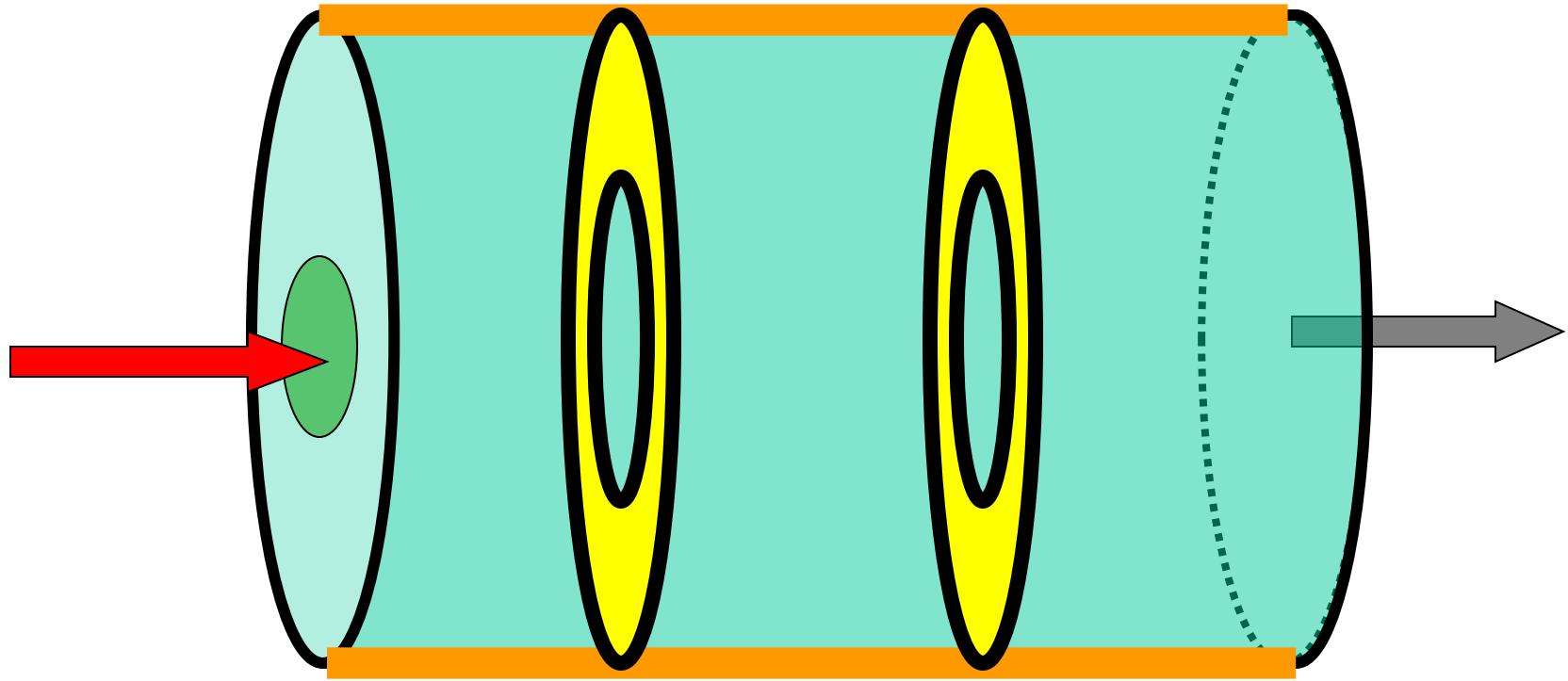


C
a
S
o

8

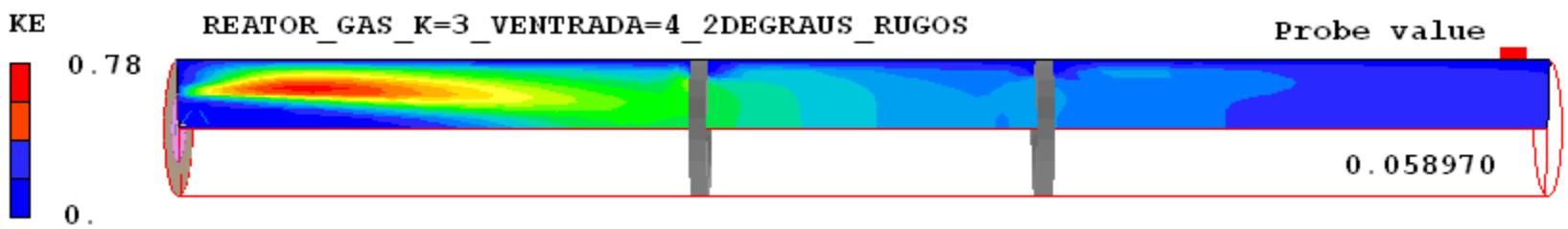
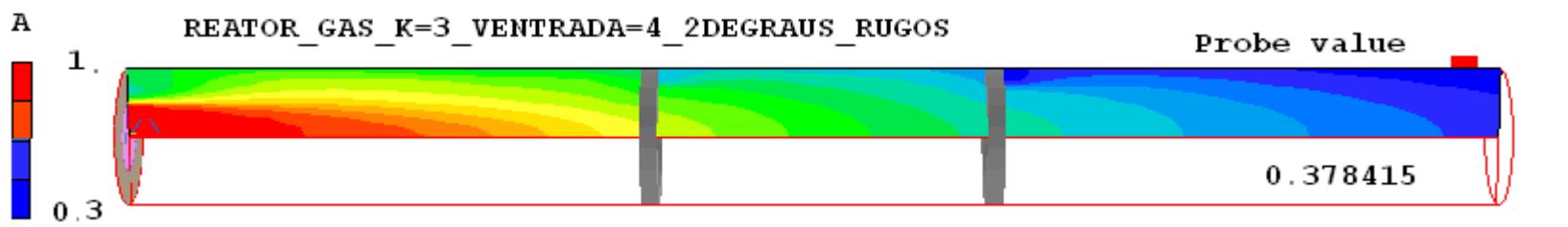
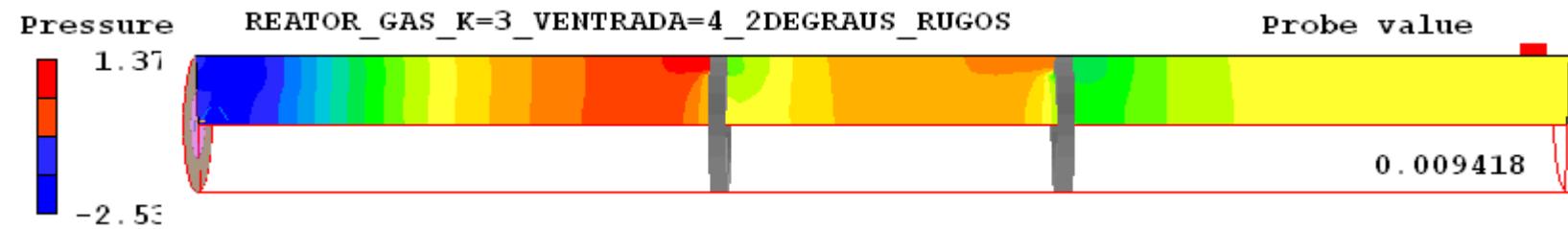
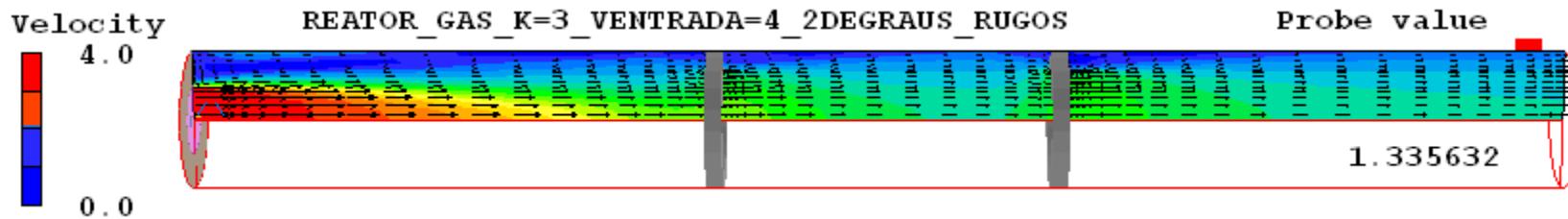


Caso 9



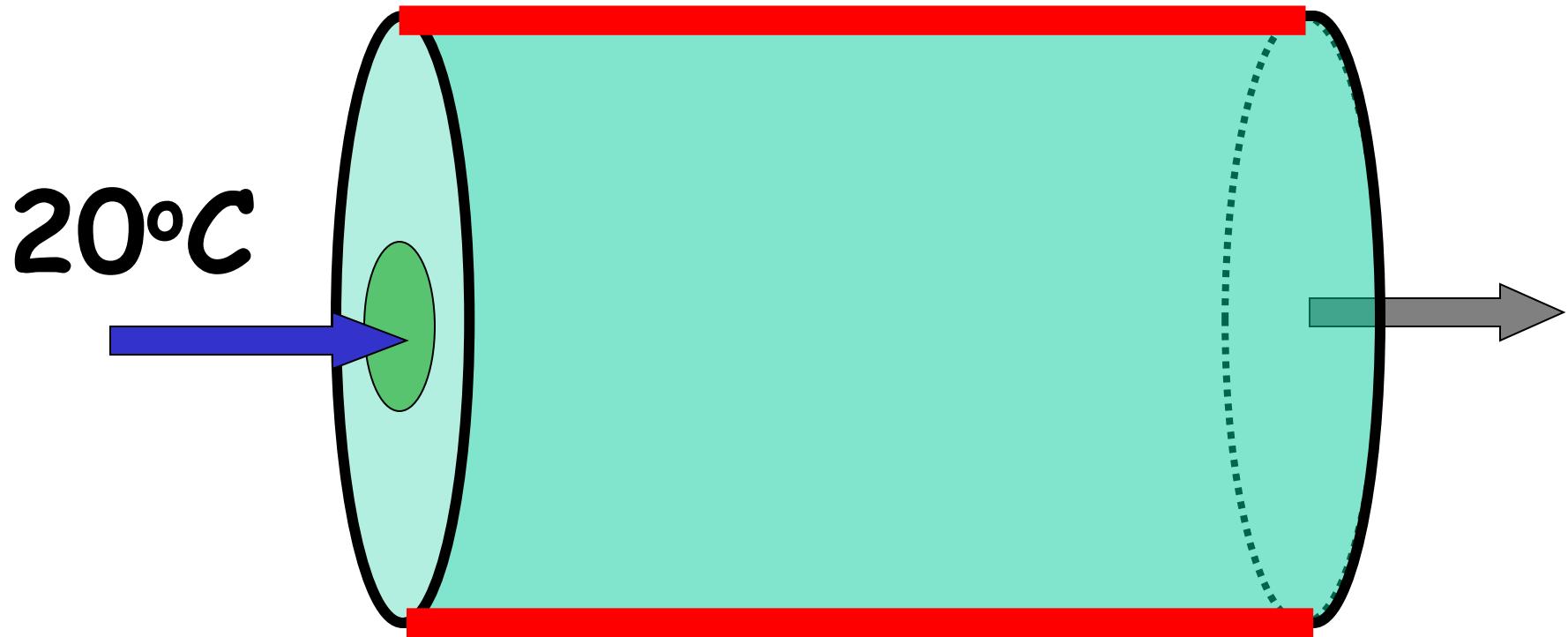
C
a
S
O

9



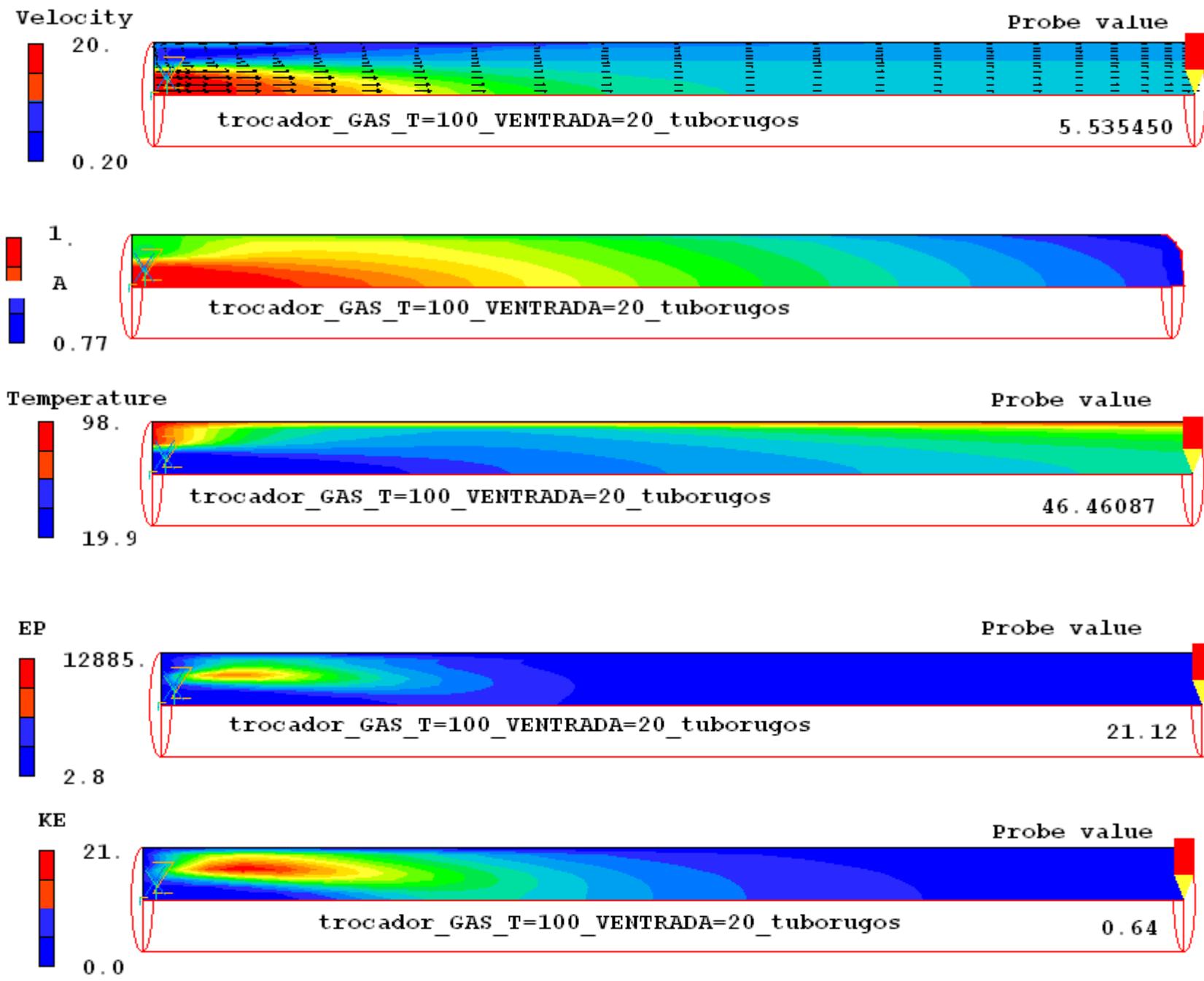
Caso 10

100°C

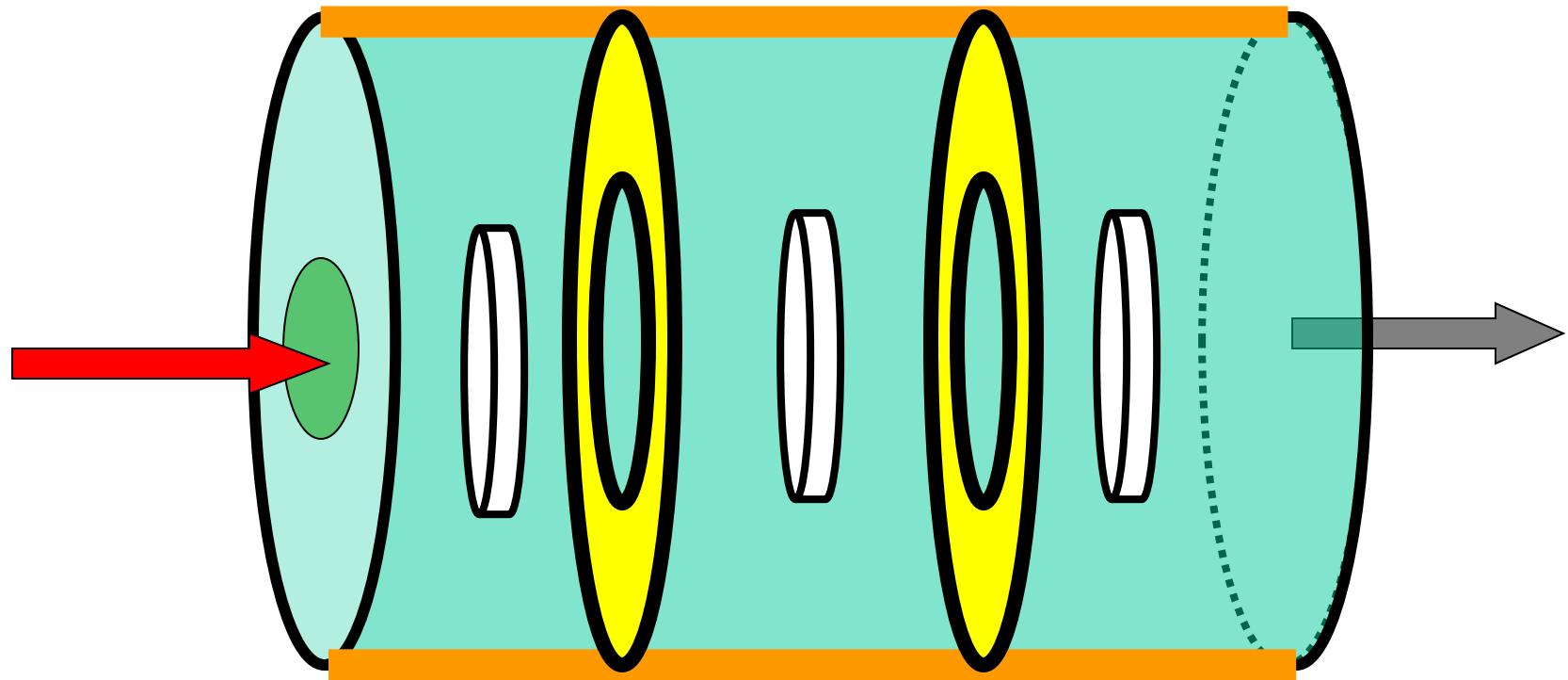


Rugosa & Quente

C
a
S
O
10

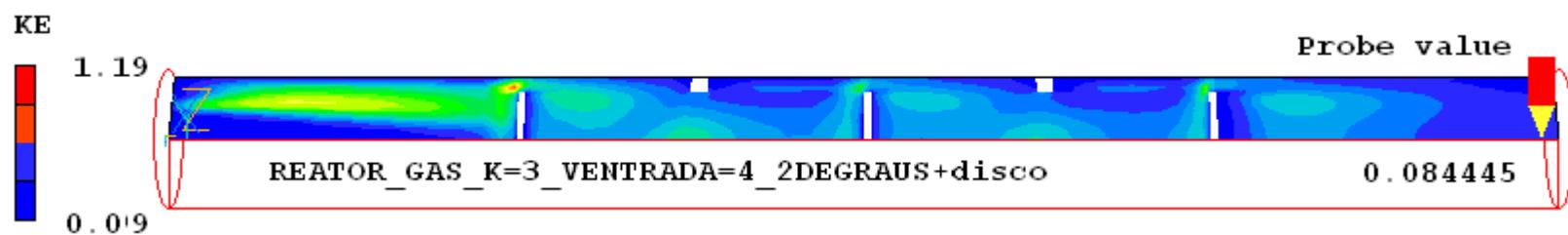
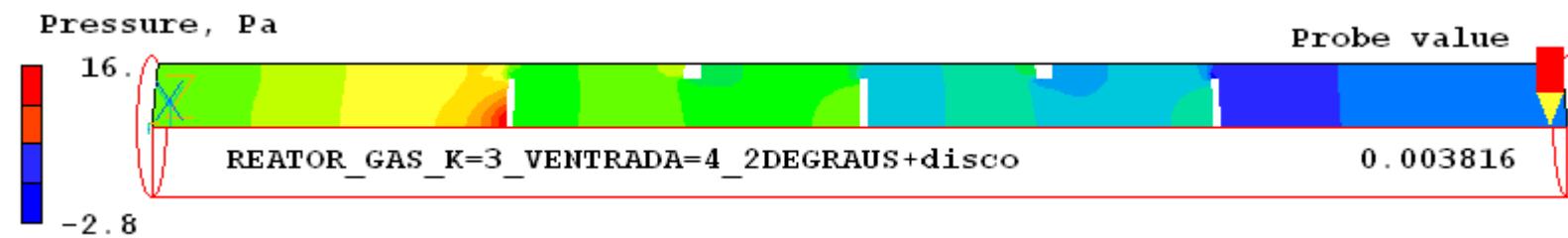
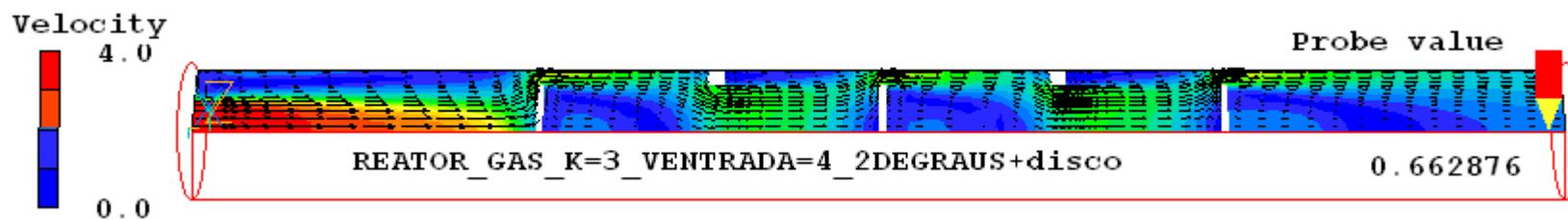


Caso 11

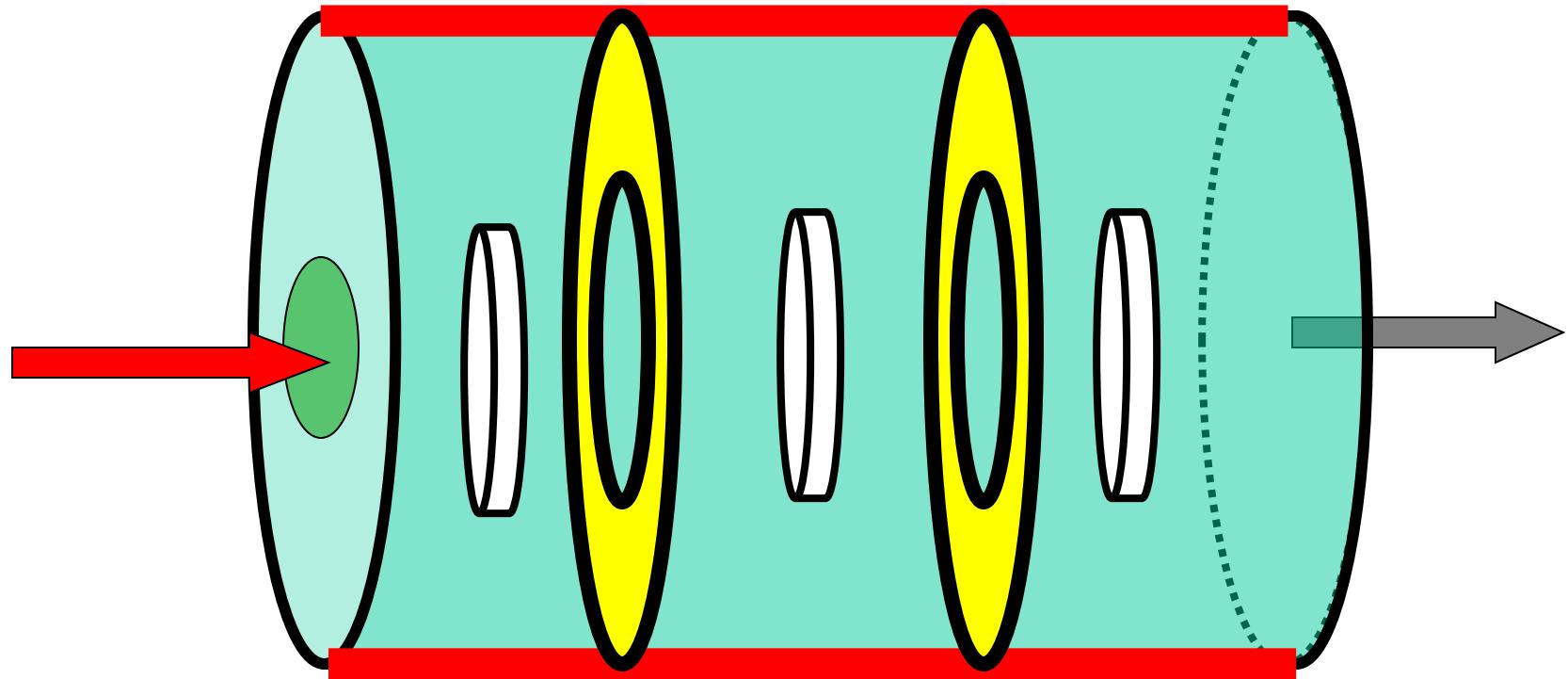


C
a
S
O

11



Caso 12



Rugosa & Quente

C
a
S
O
1
2

