

Quantum Theory of Many-Body systems in Condensed Matter (4302112) 2020

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Today's class: *Time-ordered Green's functions*

- Single-particle Green's functions.
- Retarded and advanced Green's functions.
- Many-particle Green's functions
- Lehmann representation.

Single-particle GFs: time-independent

- Great intro to Green's functions : <https://arxiv.org/abs/1604.02499>

Time-independent Schrödinger's equation:

$$\hat{H} = \hat{H}_0 + \hat{V}$$

$$\begin{cases} (E - \hat{H}) \psi_E(\vec{r}) = 0 \\ (E - \hat{H}_0) \psi_E^{(0)}(\vec{r}) = 0 \end{cases}$$

Associated GFs:

$$\begin{cases} (E - \hat{H}) G(\vec{r}, \vec{r}'; E) = \delta(\vec{r} - \vec{r}') \\ (E - \hat{H}_0) G_0(\vec{r}, \vec{r}'; E) = \delta(\vec{r} - \vec{r}') \end{cases}$$

Formally:

$$\begin{cases} G^{-1}(\vec{r}; E) = (E - \hat{H}) \\ (G_0)^{-1}(\vec{r}; E) = (E - \hat{H}_0) \end{cases} \Rightarrow \begin{cases} G^{-1}(\vec{r}; E) \psi_E(\vec{r}) = 0 \\ (G_0)^{-1}(\vec{r}; E) \psi_E^{(0)}(\vec{r}) = 0 \end{cases}$$

Single-particle GFs: time-independent

Integral equation (if we know G_0) (Lippmann-Schwinger): 

$$\psi_E(\vec{r}) = \psi_E^{(0)}(\vec{r}) + \int d\vec{r}' G_0(\vec{r}, \vec{r}'; E) V(\vec{r}') \psi_E(\vec{r}')$$

Also (formally):

$$\psi_E(\vec{r}) = \psi_E^{(0)}(\vec{r}) + \int d\vec{r}' G(\vec{r}, \vec{r}'; E) V(\vec{r}') \psi_E^{(0)}(\vec{r}')$$

Integral equation for G (if we know G_0) (Dyson's equation): 

$$G(\vec{r}, \vec{r}'; E) = G_0(\vec{r}, \vec{r}'; E) + \int d\vec{r}_1 G_0(\vec{r}, \vec{r}_1; E) V(\vec{r}_1) G(\vec{r}_1, \vec{r}'; E)$$

“Pictoric view” :Diagrams!

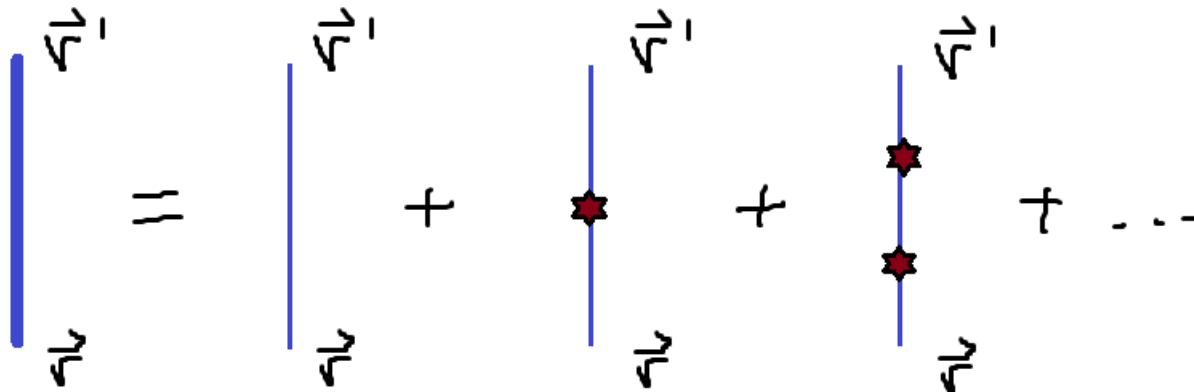
Single-particle GFs: time-independent

$$G(\vec{r}, \vec{r}'; E) = G_0(\vec{r}, \vec{r}'; E) + \int d\vec{r}_1 G_0(\vec{r}, \vec{r}_1; E) V(\vec{r}_1) G(\vec{r}_1, \vec{r}'; E)$$



$$G = G_0 + G_0 V G$$

$$G = G_0 + G_0 V G_0 + G_0 V G_0 V G_0 + \dots$$



“Pictoric view” :Diagrams!

Single-particle GFs: time-dependent

Time-dependent Schrödinger's equation:

$$\hat{H} = \hat{H}_0 + \hat{V}$$
$$\begin{cases} \left(i\hbar \frac{\partial}{\partial t} - \hat{H} \right) \psi(\vec{r}, t) = 0 \\ \left(i\hbar \frac{\partial}{\partial t} - \hat{H}_0 \right) \psi^{(0)}(\vec{r}, t) = 0 \end{cases}$$

Associated GFs:

$$\begin{cases} \left(i\hbar \frac{\partial}{\partial t} - \hat{H} \right) G(\vec{r}, t; \vec{r}', t') = \delta(\vec{r} - \vec{r}')\delta(t - t') \\ \left(i\hbar \frac{\partial}{\partial t} - \hat{H}_0 \right) G_0(\vec{r}, t; \vec{r}', t') = \delta(\vec{r} - \vec{r}')\delta(t - t') \end{cases}$$

Integral equations:

$$\psi(\vec{r}, t) = \psi^{(0)}(\vec{r}, t) + \int d\vec{r}' \int dt' G_0(\vec{r}, t; \vec{r}', t') V(\vec{r}') \psi(\vec{r}', t')$$

$$\psi(\vec{r}, t) = \psi^{(0)}(\vec{r}, t) + \int d\vec{r}' \int dt' G(\vec{r}, t; \vec{r}', t') V(\vec{r}') \psi^{(0)}(\vec{r}', t')$$

Retarded GF: propagator

Notice that: $\psi(\vec{r}, t) = \int d\vec{r}' G(\vec{r}, t; \vec{r}', t') \psi(\vec{r}', t')$ t ≠ t'

satisfies $\left(i\hbar \frac{\partial}{\partial t} - \hat{H} \right) \psi(\vec{r}, t) = 0$

“Propagator”:
(Retarded)

t > t' also: $G^R(\vec{r}, t; \vec{r}', t') = \theta(t - t') G(\vec{r}, t; \vec{r}', t')$

$G^R(\vec{r}, t; \vec{r}', t') = -i\theta(t - t') \langle \vec{r} | e^{-i\hat{H}(t-t')} | \vec{r}' \rangle \quad \hbar \equiv 1$

satisfies $\left(i\hbar \frac{\partial}{\partial t} - \hat{H} \right) G^R(\vec{r}, t; \vec{r}', t') = \delta(\vec{r} - \vec{r}') \delta(t - t')$ (Show!)

:

In general: $G^R(n, t; n', t') = -i\theta(t - t') \langle n | e^{-i\hat{H}(t-t')} | n' \rangle$

Many-particle GFs: time-ordered

Many-particle propagator:

$$G^R(\vec{r}, t; \vec{r}', t') = -i\theta(t - t') \left\langle \left[\hat{\psi}(\vec{r}, t), \hat{\psi}^\dagger(\vec{r}', t') \right]_{\mp} \right\rangle$$

$$\langle \dots \rangle = \frac{1}{Z} \text{Tr} (\hat{\rho} \dots)$$

field operators

-: Bosons

+: Fermions

Other definitions:

Advanced:

$$G^A(\vec{r}, t; \vec{r}', t') = +i\theta(t' - t) \left\langle \left[\hat{\psi}(\vec{r}, t), \hat{\psi}^\dagger(\vec{r}', t') \right]_{\mp} \right\rangle$$

“Greater”:

$$G^>(\vec{r}, t; \vec{r}', t') = -i \left\langle \hat{\psi}(\vec{r}, t) \hat{\psi}^\dagger(\vec{r}', t') \right\rangle$$

“Lesser”:

$$G^<(\vec{r}, t; \vec{r}', t') = -i(\pm 1) \left\langle \hat{\psi}^\dagger(\vec{r}', t') \hat{\psi}(\vec{r}, t) \right\rangle$$

+: Bosons

-: Fermions

Change of basis

$$\hat{a}_k^\dagger |0\rangle = |k\rangle \quad \hat{\psi}^\dagger(\vec{r}) = \sum_k (\langle \vec{r} | k \rangle)^* \hat{a}_k^\dagger \quad \hat{\psi}(\vec{r}) = \sum_k \langle \vec{r} | k \rangle \hat{a}_k$$

$$\Rightarrow G^R(\vec{r}, t; \vec{r}', t') = \sum_{kk'} \langle \vec{r} | k \rangle G_{kk'}^R(t, t') (\langle \vec{r}' | k' \rangle)^*$$

$$G_{kk'}^R(t, t') = \theta(t - t') (G_{kk'}^>(t, t') - G_{kk'}^<(t, t'))$$

$$G_{kk'}^A(t, t') = \theta(t' - t) (G_{kk'}^<(t, t') - G_{kk'}^>(t, t'))$$

$$G_{kk'}^<(t, t') = -i(\pm 1) \langle \hat{a}_{k'}^\dagger(t') \hat{a}_k(t) \rangle \quad \begin{array}{l} +: \text{Bosons} \\ -: \text{Fermions} \end{array}$$

$$G_{kk'}^>(t, t') = -i \langle \hat{a}_k(t) \hat{a}_{k'}^\dagger(t') \rangle$$

In general:

Lehmann Representation

Many-particle
spectrum:

$$\hat{H}|\alpha\rangle = E_\alpha|\alpha\rangle$$

Matrix elements:

$$\langle\alpha'|\hat{a}_k^\dagger|\alpha\rangle = (\langle\alpha|\hat{a}_k|\alpha'\rangle)^\dagger$$

We can write a GF as:



$$G_{km}^>(t, t') = \frac{-i}{Z} \sum_{\alpha\alpha'} e^{-\beta E_\alpha} \langle\alpha|\hat{a}_k|\alpha'\rangle \langle\alpha'|\hat{a}_m^\dagger|\alpha\rangle e^{+i(E_\alpha - E_{\alpha'})(t-t')}$$

In frequency:

$$G_{km}^>(\omega) = -i \frac{2\pi}{Z} \sum_{\alpha\alpha'} e^{-\beta E_\alpha} \langle\alpha|\hat{a}_k|\alpha'\rangle \langle\alpha'|\hat{a}_m^\dagger|\alpha\rangle \delta(\omega - (E_{\alpha'} - E_\alpha))$$

Lehmann Representation:

Spectral function

Lehmann Representation for the retarded GF:

$$G_{km}^R(\omega^+) = \frac{1}{Z} \sum_{\alpha\alpha'} e^{-\beta E_\alpha} \left[\frac{\langle \alpha | \hat{a}_k | \alpha' \rangle \langle \alpha' | \hat{a}_m^\dagger | \alpha \rangle}{\omega + i\eta - (E_{\alpha'} - E_\alpha)} + \frac{\langle \alpha | \hat{a}_m^\dagger | \alpha' \rangle \langle \alpha' | \hat{a}_k | \alpha \rangle}{\omega + i\eta + (E_{\alpha'} - E_\alpha)} \right]$$

Spectral function (diagonal): $A_k(\omega^+) \equiv \frac{-\text{Im}}{\pi} G_{kk}^R(\omega^+)$

$A_k(\omega^+) = \frac{1}{Z} \sum_{\alpha\alpha'} (e^{-\beta E_\alpha} + e^{-\beta E_{\alpha'}}) |\langle \alpha' | \hat{a}_k^\dagger | \alpha \rangle|^2 \delta(\omega^+ - (E_{\alpha'} - E_\alpha))$

At T=0 (transitions to the GS):

$$A_k^{T=0}(\omega^+) = \frac{1}{Z} \sum_{\alpha'} |\langle \alpha' | \hat{a}_k^\dagger | 0 \rangle|^2 \delta(\omega^+ - E_{\alpha'})$$

