

Quantum Theory of Many-Body systems in Condensed Matter (4302112) 2020

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Today's class: *Equations of motion and time-ordered correlation functions*

- Equations of motion for creation and destruction operators.
- Retarded and advanced correlation functions.
- Fourier transform

Equations of motion for operators

Heisenberg Picture: equation of motion:
$$\left\{ \begin{array}{l} \frac{d\hat{a}_k(t)}{dt} = \frac{i}{\hbar} [\hat{H}, \hat{a}_k(t)]_- \\ \hat{a}_k(t) \equiv e^{+i\hat{H}t/\hbar} \hat{a}_k e^{-i\hat{H}t/\hbar} \end{array} \right.$$

Commutator/anti-commutator at *different times*? $[\hat{a}_k(t_2), \hat{a}_m^\dagger(t_1)]_{\pm} = ?$

Special case (diagonal Hamiltonian):

$$\hat{H} = \sum_k \epsilon_k \hat{a}_k^\dagger \hat{a}_k \quad \Rightarrow \quad \left\{ \begin{array}{l} \hat{a}_k(t) = e^{-i\epsilon_k t/\hbar} \hat{a}_k(0) \\ [\hat{a}_k(t_1), \hat{a}_m^\dagger(t_2)]_{\pm} = e^{-i(\epsilon_k t_1 - \epsilon_m t_2)/\hbar} \delta_{km} \end{array} \right.$$

Equations of motion for operators

“Tight-binding-like”?

$$\hat{H} = \sum_k \epsilon_k \hat{a}_k^\dagger \hat{a}_k - t \sum_{\substack{k_1, k_2 \\ (k_1 \neq k_2)}} \hat{a}_{k_1}^\dagger \hat{a}_{k_2} + \hat{a}_{k_2}^\dagger \hat{a}_{k_1}$$

In this case you can still solve (“close”) the equations of motion (**Assignment**):

$$\left\{ \begin{array}{l} \frac{d\hat{a}_k(t)}{dt} = \frac{i}{\hbar} \left[\hat{H}, \hat{a}_k(t) \right]_- \\ \frac{d\hat{a}_k^\dagger(t)}{dt} = \frac{i}{\hbar} \left[\hat{H}, \hat{a}_k^\dagger(t) \right]_- \end{array} \right.$$

In the *interacting* case, however, you cannot do it (**Assignment**):

$$\hat{H} = \sum_k \epsilon_k \hat{a}_k^\dagger \hat{a}_k + \frac{1}{2} \sum_{k_1 k_2 q} \hat{a}_{k_1+q}^\dagger \hat{a}_{k_2-q}^\dagger \hat{a}_{k_2} \hat{a}_{k_1}$$

Time-dependent correlation functions

$$\langle \hat{A}(t) \hat{B}(t') \rangle = \frac{1}{Z} \text{Tr} \left(\hat{\rho} \hat{A}(t) \hat{B}(t') \right)$$

depends only on (t-t')

$$\blacksquare = \frac{1}{Z} \text{Tr} \left(e^{-\beta \hat{H}} e^{i\hat{H}(t-t')/\hbar} \hat{A} e^{-i\hat{H}(t-t')/\hbar} \hat{B} \right)$$

$$\begin{cases} \text{“Retarded”}: & C_{\hat{A}, \hat{B}}^r(t-t') = -i\theta(t-t') \langle \hat{A}(t) \hat{B}(t') \rangle & t > t' \\ \text{“Advanced”}: & C_{\hat{A}, \hat{B}}^a(t-t') = +i\theta(t'-t) \langle \hat{A}(t) \hat{B}(t') \rangle & t < t' \end{cases}$$

Fourier transform: retarded case

$t > t'$


FT of the retarded correlation function:

$$C_{\hat{A},\hat{B}}^r(\omega) = \int_{-\infty}^{\infty} dt(t-t') e^{i\omega(t-t')} C_{\hat{A},\hat{B}}^r(t-t')$$

Ok, as long as:

$$\lim_{(t-t') \rightarrow \infty} C_{\hat{A},\hat{B}}^r(t-t') = 0$$

Physically: some relaxation mechanism

Or: $\omega \rightarrow \omega^+ = \omega + i\eta$ 

$$\left\{ \begin{array}{l} C_{\hat{A},\hat{B}}^r(\omega^+) = \int_{-\infty}^{\infty} dt e^{i\omega^+ t} C_{\hat{A},\hat{B}}^r(t) \\ C_{\hat{A},\hat{B}}^r(t) = \lim_{\eta \rightarrow 0} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} C_{\hat{A},\hat{B}}^r(\omega^+) \end{array} \right.$$

Fourier transform: advanced case

$t < t'$

FT of the advanced correlation function:

$$C_{\hat{A},\hat{B}}^a(\omega) = \int_{-\infty}^{\infty} dt(t-t') e^{i\omega(t-t')} C_{\hat{A},\hat{B}}^a(t-t')$$

Ok, as long as:

$$\lim_{(t-t') \rightarrow -\infty} C_{\hat{A},\hat{B}}^a(t-t') = 0 \quad \text{Physically: some relaxation mechanism}$$

Or: $\omega \rightarrow \omega^- = \omega - i\eta$

$$\left\{ \begin{array}{l} C_{\hat{A},\hat{B}}^a(\omega^-) = \int_{-\infty}^{\infty} dt e^{i\omega^- t} C_{\hat{A},\hat{B}}^a(t) \\ C_{\hat{A},\hat{B}}^a(t) = \lim_{\eta \rightarrow 0} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} C_{\hat{A},\hat{B}}^a(\omega^-) \end{array} \right.$$