

# Quantum Theory of Many-Body systems in Condensed Matter (4302112) 2020

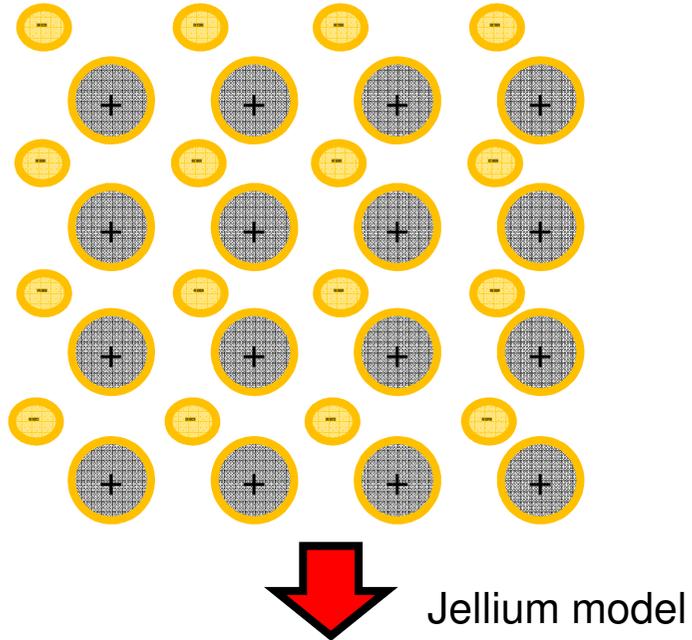
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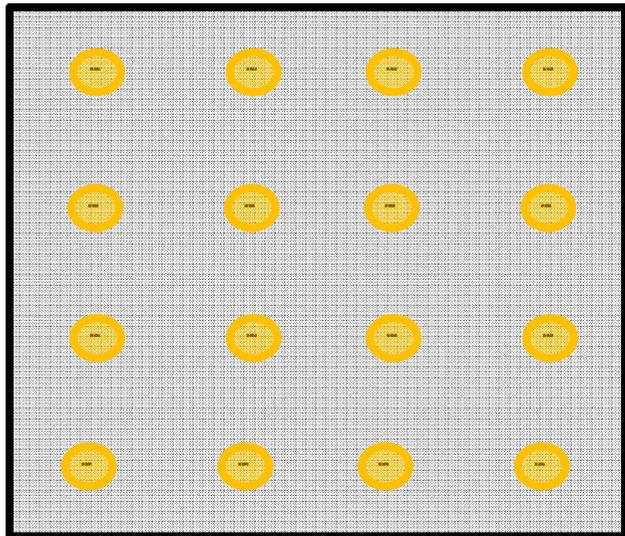
Today's class: *Non-interacting electron gas*

- Jellium model
- Non-interacting Fermi gas.
- Ground-state: Fermi energy.
- Density of states.

# Electrons in a solid: the jellium model



Jellium model



- Model for electrons in a metal.
- “Core ions”: form a homogenous “fluid”, positively charged (positive background).
- Conduction electrons: move “freely” in this background, weakly interacting with each other.
- The whole system is charge neutral.
- Electrons behave as in a “Fermi gas”.

# Free fermion “gas”: single-particle picture

- Single-particle Hamiltonian:  $\hat{h}|\varphi_{\mathbf{k}}\rangle = \varepsilon_{\mathbf{k}}|\varphi_{\mathbf{k}}\rangle$        $\mathbf{k} = (k_x, k_y, k_z)$

- Schrödinger’s equation for a particle of mass  $m^*$ :

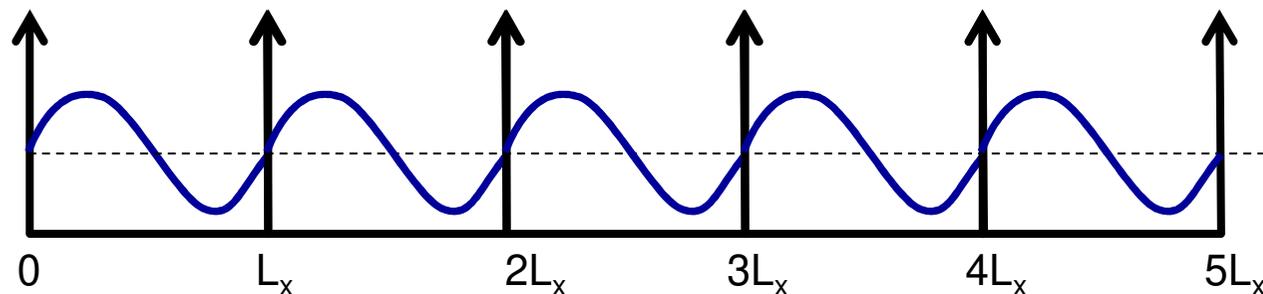
$$\frac{\hbar^2}{2m^*} \nabla^2 \varphi(\mathbf{r}) = \varepsilon \varphi(\mathbf{r}) \Rightarrow \nabla^2 \varphi(\mathbf{r}) = -k^2 \varphi(\mathbf{r})$$

Solution (plane waves):

$$\varphi_{\mathbf{k}}(\mathbf{r}) = A e^{i\mathbf{k}\cdot\mathbf{r}} \quad |\mathbf{k}|^2 = \frac{2m^* \varepsilon}{\hbar^2}$$

- Let us consider periodic boundary conditions:

$$\varphi_{\mathbf{k}}(x + L_x, y, z) = \varphi_{\mathbf{k}}(x, y + L_y, z) = \varphi_{\mathbf{k}}(x, y, z + L_z) = \varphi_{\mathbf{k}}(x, y, z)$$



$$\varphi_{\mathbf{k}}(\mathbf{r} + \mathbf{L}) = A e^{i\mathbf{k}\cdot\mathbf{r}} e^{i\mathbf{k}\cdot\mathbf{L}}$$

$$e^{i\mathbf{k}\cdot\mathbf{L}} = 1$$

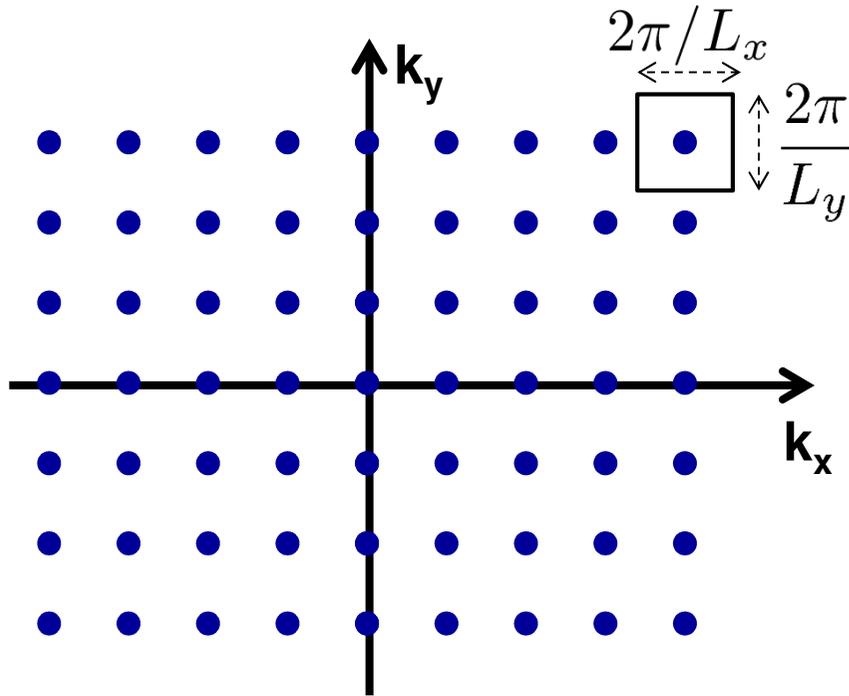
- Solution: energy quantization:

$$k_x L_x + k_y L_y + k_z L_z = 2m\pi$$

$$m = n_x + n_y + n_z = 0, \pm 1, \pm 2, \dots$$

# “Counting” the states.

- Possible  $\mathbf{k}$  values:



$$k_{x(y,z)} = \frac{2\pi n_{x(y,z)}}{L_{x(y,z)}} \text{ with } n_{x(y,z)} = 0, \pm 1, \pm 2, \dots$$

- State and energy for each  $\mathbf{k}$ :

$$\varphi_{\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{L_x L_y L_z}} e^{i\mathbf{k} \cdot \mathbf{r}} \quad \boxed{\varepsilon_{\mathbf{k}} = \frac{\hbar^2 |\mathbf{k}|^2}{2m^*}}$$

- Number of states in  $\mathbf{k}$  space (3D) :

“Volume” of each state in  $\mathbf{k}$

$$V_{\mathbf{k}}^{1\text{st}} = \left(\frac{2\pi}{L_x}\right) \left(\frac{2\pi}{L_y}\right) \left(\frac{2\pi}{L_z}\right) = \frac{(2\pi)^3}{V_{\mathbf{r}}}$$

Tip: to go from discrete  $\rightarrow$  continuum!

$$\boxed{\frac{1}{V_{\mathbf{r}}} \sum_{\mathbf{k}} \rightarrow \int \frac{d^3 \mathbf{k}}{(2\pi)^3}}$$

Number of states in a given volume in  $\mathbf{k}$ :

$$N_{\text{st}} = \sum_{\mathbf{k}} = \int \frac{d^3 \mathbf{k}}{V_{\mathbf{k}}^{1\text{st}}} = V_{\mathbf{r}} \int \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

# N-particle Hamiltonian

- Need to account for spin ( $\sigma=\uparrow,\downarrow$ ). The single-particle states are then:

$$\hat{h}|\varphi_{\mathbf{k},\sigma}\rangle = \varepsilon_{\mathbf{k}}|\varphi_{\mathbf{k},\sigma}\rangle \quad |\varphi_{\mathbf{k},\sigma}\rangle = c_{\mathbf{k},\sigma}^{\dagger}|0\rangle$$

- N-body Hamiltonian:

$$\hat{H} = \sum_{n=1}^N \hat{h}^{(n)} = \sum_{\mathbf{k},\sigma=\uparrow,\downarrow} \varepsilon_{\mathbf{k}} \hat{c}_{\mathbf{k},\sigma}^{\dagger} \hat{c}_{\mathbf{k},\sigma} = \sum_{\mathbf{k},\sigma=\uparrow,\downarrow} \varepsilon_{\mathbf{k}} \hat{n}_{\mathbf{k},\sigma} \quad \left\{ \begin{array}{l} \varepsilon_{\mathbf{k}} = \frac{\hbar^2 |\mathbf{k}|^2}{2m^*} \\ n_{\mathbf{k},\sigma} = 0, 1 \end{array} \right.$$

- Eigenstates and eigenenergies:

$$\hat{H}|\dots n_{\mathbf{k},\uparrow}, n_{\mathbf{k},\downarrow} \dots\rangle = \left[ \sum_{\mathbf{k}} \varepsilon_{\mathbf{k}} (n_{\mathbf{k},\uparrow} + n_{\mathbf{k},\downarrow}) \right] |\dots n_{\mathbf{k},\uparrow}, n_{\mathbf{k},\downarrow} \dots\rangle$$

- Let's consider N even (non-polarized gas): (N/2 spin  $\uparrow$ , N/2 spin  $\downarrow$ ).

$$N = \sum_{\mathbf{k}} n_{\mathbf{k},\uparrow} + n_{\mathbf{k},\downarrow} \quad \text{with} \quad \frac{N}{2} = \sum_{\mathbf{k}} n_{\mathbf{k},\uparrow} = \sum_{\mathbf{k}} n_{\mathbf{k},\downarrow}$$

# N-particle ground state

- For a given set of occupations:  $E_{n_{0,\uparrow}, n_{0,\downarrow} \dots n_{\mathbf{k},\uparrow}, n_{\mathbf{k},\downarrow} \dots} = \sum_{\mathbf{k}} \varepsilon_{\mathbf{k}} (n_{\mathbf{k},\uparrow} + n_{\mathbf{k},\downarrow})$
- Ground-state of the N-particle system:  $\varepsilon_{\mathbf{k}} = \frac{\hbar^2 |\mathbf{k}|^2}{2m^*}$

$$\left\{ \begin{array}{l} \hat{H} |\text{GS}\rangle = E_{\text{GS}} |\text{GS}\rangle \\ E_{\text{GS}} = \langle \text{GS} | \hat{H} | \text{GS} \rangle \end{array} \right. \quad \langle \text{GS} | \hat{N} | \text{GS} \rangle = \left\langle \text{GS} \left| \sum_{\mathbf{k},\sigma} \hat{n}_{\mathbf{k},\sigma} \right| \text{GS} \right\rangle = N$$

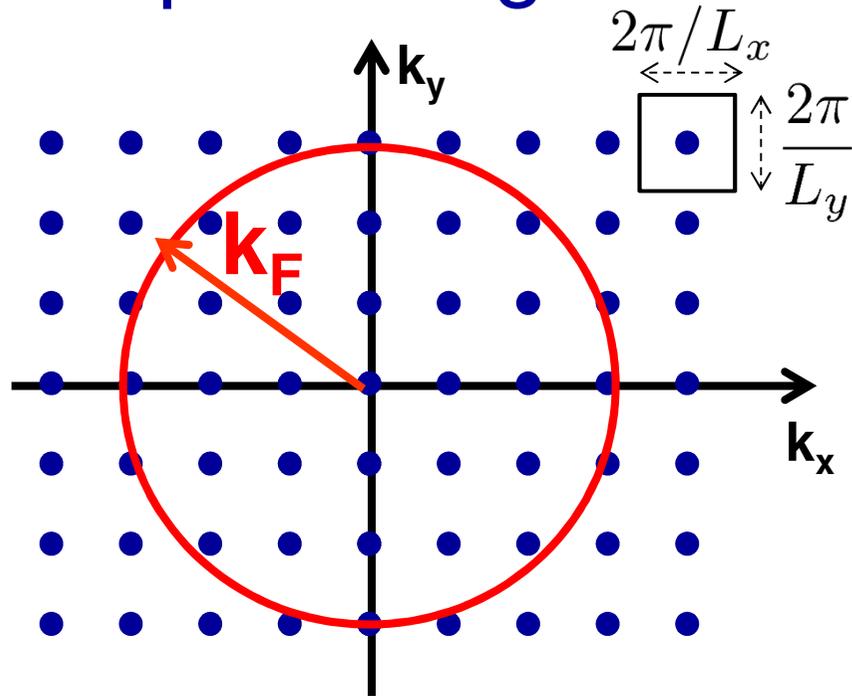
- Ground state energy :

$$E_{\text{GS}} = \sum_{|\mathbf{k}| < |\mathbf{k}_F|} \varepsilon_{\mathbf{k}} (n_{\mathbf{k},\uparrow} + n_{\mathbf{k},\downarrow})$$

$\mathbf{k}_F$  : Fermi wave vector

$$\left\{ \begin{array}{l} n_{|\mathbf{k}| \leq |\mathbf{k}_F|, \sigma = \uparrow, \downarrow} = 1 \\ n_{|\mathbf{k}| > |\mathbf{k}_F|, \sigma = \uparrow, \downarrow} = 0 \\ \frac{N}{2} = \sum_{|\mathbf{k}| \leq |\mathbf{k}_F|} \langle \text{GS} | \hat{n}_{\mathbf{k},\uparrow} | \text{GS} \rangle = \sum_{|\mathbf{k}| \leq |\mathbf{k}_F|} \langle \text{GS} | \hat{n}_{\mathbf{k},\downarrow} | \text{GS} \rangle . \end{array} \right.$$

# N-particle ground state: Fermi Energy.



- Fermi energy  $\mathcal{E}_F$ : single-particle energy of the latest occupied states.
- $\mathcal{E}=\mathcal{E}_F$ : **sphere in k space.**
- The radius of this sphere is  $k_F$

$$k_F = \frac{\sqrt{2m^* \mathcal{E}_F}}{\hbar}$$

- In each occupied state ( $k \leq k_F$ ) there are 2 Fermions (Pauli).

- For N fermions:

$$N = 2 \sum_{|\mathbf{k}| \leq |\mathbf{k}_F|} \langle \text{GS} | \hat{n}_{\mathbf{k}, \uparrow} | \text{GS} \rangle = 2 \sum_{|\mathbf{k}| < k_F} = 2V_r \int_{V_{\mathbf{k}}} \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

Integrating over the Fermi sphere in k space:

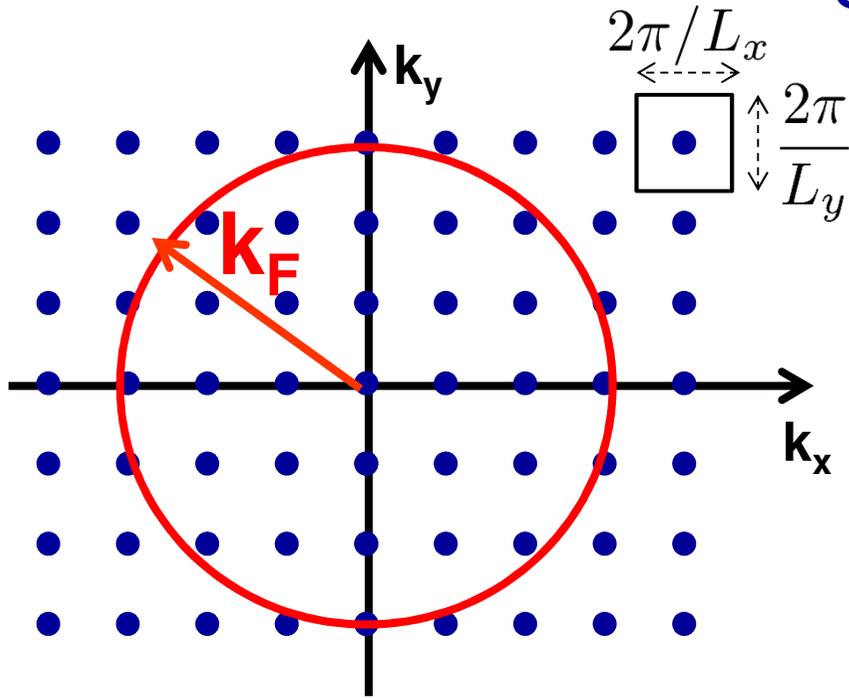
$$\blacksquare \quad \frac{N}{V_r} = \frac{2}{8\pi^3} \frac{4}{3} \pi k_F^3 = \frac{k_F^3}{3\pi^2}$$

- In terms of the density  $n=N/V_r$ :

$$\mathcal{E}_F = \left( \frac{\hbar^2 (3\pi^2)^{2/3}}{2m^*} \right) n^{2/3}$$

**In 3D,  $\mathcal{E}_F \sim n^{2/3}$  !**

# Ground-state energy of the N-particle state.



- Adding the energies of each occupied state:

$$E_{GS} = 2 \sum_{|\mathbf{k}| < |\mathbf{k}_F|} \varepsilon_{\mathbf{k}} n_{\mathbf{k}, \uparrow} = 2 \sum_{|\mathbf{k}| < k_F} \varepsilon_{\mathbf{k}} = 2V_{\mathbf{r}} \int_{V_{\mathbf{k}}} \frac{\hbar^2 k^2}{2m^*} \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

- Integrating e rearranging:

$$E = \frac{3}{5} V_{\mathbf{r}} n \varepsilon_F \Rightarrow \frac{E}{N} = \frac{3}{5} \varepsilon_F$$

Energy *per particle* in a Fermion gas.

- Typical values in metals (Copper) :

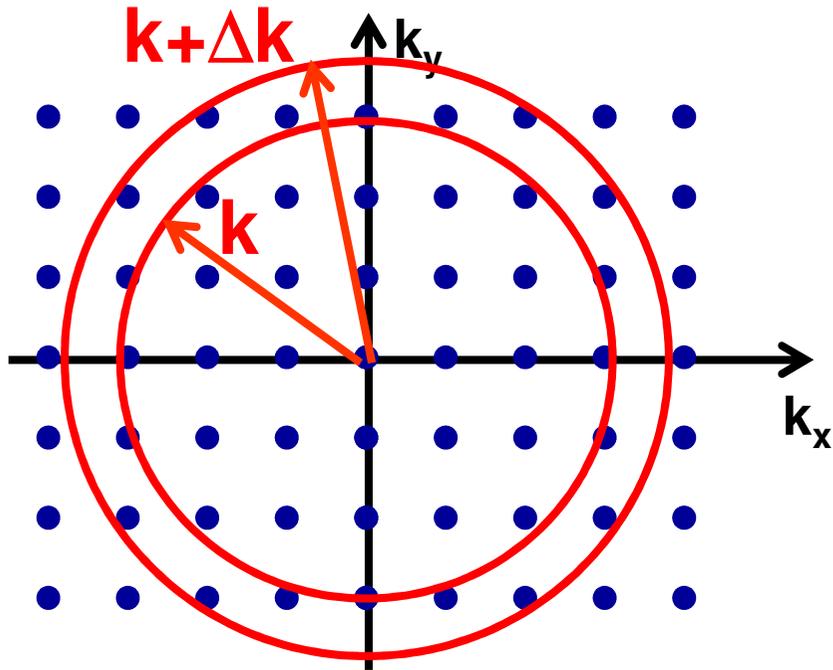
$$\varepsilon_F \approx 7.03 \text{ eV} = 81600 \text{ K}/k_B$$

$$\lambda_F = \frac{2\pi}{k_F} \approx 4.6 \text{ \AA}$$

$$v_F = \frac{\hbar k_F}{m^*} \approx 1.57 \times 10^6 \text{ m/s} \approx 0.005c$$

- Useful! (1 meV = 11.6 K/ $k_B$ )

# Density of states.



- No of states between  $k$  and  $k+\Delta k$  (3D):

$$\Delta N_{st}^{\Delta k} = 2 \frac{\Delta V_{\mathbf{k}}}{V_{\mathbf{k}}^{1st}} = V_{\mathbf{r}} \frac{8\pi k^2 dk}{(2\pi)^3}$$

$$\Delta N_{st}^{\Delta k} = V_{\mathbf{r}} \frac{k^2 dk}{\pi^2}$$

- No. of states between  $\mathcal{E}$  and  $\mathcal{E}+\Delta\mathcal{E}$  :

$$\Delta N_{st}^{\Delta \mathcal{E}} = \rho(\mathcal{E}) d\mathcal{E}$$

where  $\rho(\mathcal{E})$  is the **density of states**:

- Converting from  $k$  to  $\mathcal{E}$ :

$$\mathcal{E}_k = \frac{\hbar^2 k^2}{2m^*} \Rightarrow d\mathcal{E} = \frac{\hbar^2 k dk}{m^*}$$

Since the number of states is the same:

$$\Delta N_{st}^{\Delta k} = \Delta N_{st}^{\Delta \mathcal{E}} = \frac{V_{\mathbf{r}}}{2\pi^2} \left( \frac{2m^*}{\hbar^2} \right)^{\frac{3}{2}} \mathcal{E}^{\frac{1}{2}} d\mathcal{E}$$

- Density of states in 3D:

$$\rho_{3D}(\mathcal{E}) = \frac{V_{\mathbf{r}}}{2\pi^2} \left( \frac{2m^*}{\hbar^2} \right)^{\frac{3}{2}} \mathcal{E}^{\frac{1}{2}}$$