

Quantum Theory of Many-Body systems in Condensed Matter (4302112) 2020

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Today's class: *Review of Quantum Mechanics*

- Single-particle quantum systems
- *Example: Harmonic oscillator*
- *Assignments: Ladder operators.*

Single-particle systems ("First quantization")

- Single-particle Hamiltonian:
$$\hat{H} = \frac{(\vec{p} - q\vec{A})^2}{2m^*} + U(\vec{r})$$

Non-relativistic particle under a potential $U(\vec{r})$ and magnetic field $\vec{B} = \nabla \times \vec{A}$

- Schrödinger's equation (Dirac's notation):

$$i\hbar \frac{\partial |\psi(t)\rangle}{\partial t} = \hat{H} |\psi(t)\rangle$$

"ket"

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$$(|\psi(t)\rangle)^\dagger = \langle\psi(t)| \quad \text{"bra"}$$

$$\langle\vec{r}|\psi(t)\rangle = \psi(\vec{r}, t)$$

$$\langle\phi|\psi\rangle = \int d^3\vec{r} \phi^*(\vec{r})\psi(\vec{r})$$

Basis in the Hilbert space

$\{|\varphi_\alpha\rangle\} \rightarrow$ orthonormal basis (complete set) of the single-particle Hilbert's space.

$$\langle\varphi_\alpha|\varphi_\beta\rangle = \delta_{\alpha\beta} \quad \text{and} \quad \sum_{\alpha} |\varphi_\alpha\rangle\langle\varphi_\alpha| = \mathbb{1}$$

We can write: $|\psi(t)\rangle = \sum_{\alpha} C_{\alpha}(t) |\varphi_{\alpha}\rangle$ with $C_{\alpha}(t) = \langle\varphi_{\alpha}|\psi(t)\rangle$

Operators! $\left\{ \begin{array}{l} \hat{A}|\psi\rangle = |\phi\rangle \\ \langle\psi|\hat{A}^{\dagger} = \langle\phi| \end{array} \right.$ Representation in the $\{|\varphi_{\alpha}\rangle\}$ basis:

$$\left\{ \begin{array}{l} A_{\alpha\beta} = \langle\varphi_{\alpha}|\hat{A}|\varphi_{\beta}\rangle = \langle\varphi_{\alpha}| \left(\hat{A}|\varphi_{\beta}\rangle \right) \\ A_{\beta\alpha}^{*} = \langle\varphi_{\beta}|\hat{A}^{\dagger}|\varphi_{\alpha}\rangle = \left(\langle\varphi_{\beta}|\hat{A}^{\dagger} \right) |\varphi_{\alpha}\rangle \end{array} \right.$$

Observables (Hermitian operators)

$$\left\{ \begin{array}{l} \hat{A} = \hat{A}^\dagger \rightarrow \text{Hermitian operator.} \\ \hat{A}|\varphi_a\rangle = a|\varphi_a\rangle \rightarrow a \in \mathbb{R} \quad \text{Real eigenvalues (Physical observable).} \\ P_\psi(a) \equiv |\langle\varphi_a|\psi\rangle|^2 \rightarrow \text{Probability of measuring the value } a \text{ for the} \\ \text{observable } \hat{A} \text{ if the particle is in state } |\psi\rangle \end{array} \right.$$

Obs: if the spectrum is continuous, this becomes a *probability density*.

Example: position operator:

$$dP_\psi(\vec{r}) \equiv |\langle\vec{r}|\psi(t)\rangle|^2 d^3\vec{r} = |\psi(\vec{r}, t)|^2 d^3\vec{r}$$

Probability of the particle in state $|\psi(t)\rangle$ being in a volume $d^3\vec{r}$ around position $\vec{r}$ at a time t .

Example: 1D Harmonic oscillator

Hamiltonian:

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2 \quad \text{with} \quad [\hat{x}, \hat{p}] = i\hbar$$

Position rep.

$$\langle \vec{r} | \hat{p} | \psi(t) \rangle = \frac{\hbar}{i} \frac{\partial}{\partial x} \psi(x, t)$$

Let us define:

$$\left\{ \begin{array}{l} \hat{X} = \sqrt{\frac{m\omega}{\hbar}} \hat{x} \\ \hat{P} = \frac{1}{\sqrt{m\hbar\omega}} \hat{p} \end{array} \right. \Rightarrow \hat{H} = \frac{\hbar\omega}{2} (\hat{P}^2 + \hat{X}^2)$$

and

$$\left\{ \begin{array}{l} \hat{a} = \frac{1}{\sqrt{2}} (\hat{X} + i\hat{P}) \\ \hat{a}^\dagger = \frac{1}{\sqrt{2}} (\hat{X} - i\hat{P}) \end{array} \right. \longleftrightarrow \left\{ \begin{array}{l} \hat{X} = \frac{1}{\sqrt{2}} (\hat{a}^\dagger + \hat{a}) \\ \hat{P} = \frac{i}{\sqrt{2}} (\hat{a}^\dagger - \hat{a}) \end{array} \right.$$

Example: 1D Harmonic oscillator

Let us show that: $[\hat{X}, \hat{P}] = i$ and $[\hat{a}, \hat{a}^\dagger] = 1$

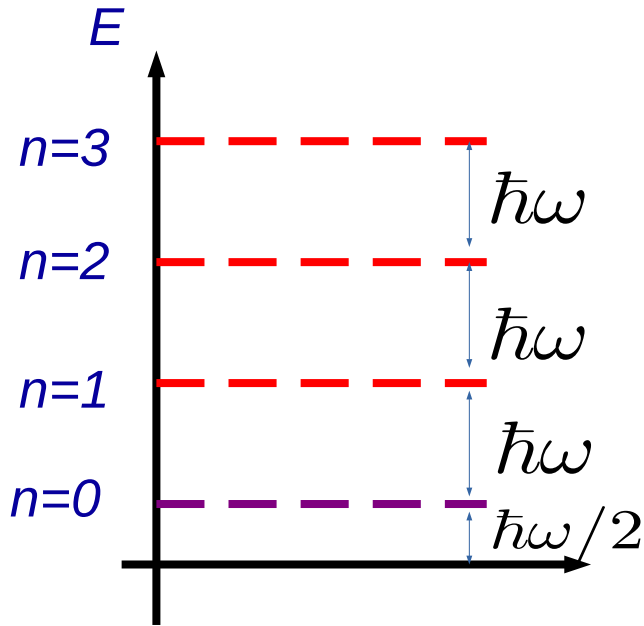
Also: $\hat{H} = \hbar\omega \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right)$ You will show (Assignment) that:
 $[\hat{a}, \hat{a}] = [\hat{a}^\dagger, \hat{a}^\dagger] = 0$

Defining: $\hat{N} \equiv \hat{a}^\dagger \hat{a}$ (Show it commutes with H (Assignment))

and its eigenstates : $\hat{N}|n\rangle = n|n\rangle$ where n is an integer (Show).

It follows that:
(Show)
$$\left\{ \begin{array}{l} \hat{H}|n\rangle = E_n|n\rangle \\ a^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle \\ a|n\rangle = \sqrt{n}|n-1\rangle \end{array} \right. \quad \begin{array}{l} E_n = \hbar\omega \left(n + \frac{1}{2} \right) \quad n = 0, 1, \dots \\ |n\rangle = \frac{(\hat{a}^\dagger)^n}{\sqrt{n!}} |0\rangle \end{array}$$

Single-particle spectrum



Single-particle
state:

$$\hat{H}|n\rangle = E_n|n\rangle$$

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right) \quad n = 0, 1, \dots \quad \text{Ground state: } E_0 = \hbar\omega$$

$\{|n\rangle\} \rightarrow$ orthonormal basis of the
(single-particle) Hilbert's space

$$\langle n|n'\rangle = \delta_{nn'} \quad \text{and} \quad \sum_n |n\rangle\langle n| = \mathbb{1}$$

$$|\psi(t)\rangle = \sum_n C_n(t) |n\rangle \quad \text{with} \quad C_n(t) = \langle n|\psi(t)\rangle$$

$$P_\psi(n) \equiv |\langle n|\psi\rangle|^2 \rightarrow \text{Probability of measuring the value } E_n \text{ for the energy if the particle is in state } |\psi\rangle$$

Assignments: Ladder operators

Problem 1

Consider the 1D harmonic oscillator Hamiltonian $H = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$ and the operators $\hat{a}$ e $\hat{a}^\dagger$ discussed in class.

1. Show that $[\hat{a}, \hat{a}^\dagger] = 1$ and that $[\hat{a}, \hat{a}] = [\hat{a}^\dagger, \hat{a}^\dagger] = 0$.
2. Show that $\hat{N} \equiv \hat{a}^\dagger \hat{a}$ commutes with H .
3. Show that, if $\hat{N}|\lambda\rangle = \lambda|\lambda\rangle$ then λ is a positive integer. Suggestion: show that $\hat{a}|\lambda\rangle$ is also an eigenstate of $\hat{N}$ and that $\hat{a}|0\rangle = 0$.
4. If $\hat{N}|n\rangle = n|n\rangle$, show that $|n\rangle \propto (\hat{a}^\dagger)^n|0\rangle$ where " $|0\rangle$ " is such that $\hat{a}|0\rangle = 0$.
5. Show that:

$$\begin{aligned}\hat{a}|n\rangle &= \sqrt{n} |n-1\rangle \\ \hat{a}^\dagger|n\rangle &= \sqrt{n+1} |n+1\rangle .\end{aligned}$$

6. Show that:

$$\hat{a}\hat{a}^\dagger (\hat{a}^\dagger)^{p-1}|0\rangle = p(\hat{a}^\dagger)^{p-1}|0\rangle.$$

(p is an integer)

Problem 2

Consider the single particle Hamiltonian:

$$H = \hbar\omega \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right) + \hbar\omega_0 \left(\hat{a}^\dagger + \hat{a} \right) ,$$

where $\hat{a}$ and $\hat{a}^\dagger$ are (bosonic) ladder operators satisfying $[\hat{a}, \hat{a}^\dagger] = 1$ and ω and ω_0 are positive constants.

1. Discuss the physical origin of the second term in H , by comparing it with the harmonic oscillator
2. Diagonalize H and find its eigenenergies. Suggestion: introduce an operator $\hat{\alpha} = \hat{a} + \omega_0/\omega$.