

Aula 9

Função de onda de duas partículas

Adição de spins

$$4 = 3 + 1$$

Função de onda de duas partículas

Partículas não interagentes: $\chi = \chi^{(1)} \chi^{(2)}$

Vamos verificar se é uma boa hipótese:

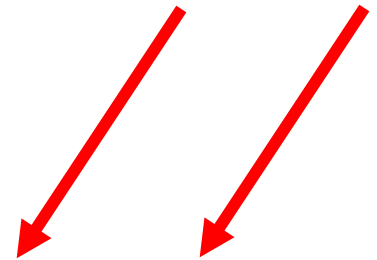
Equação de Schrödinger independente do tempo

$$\begin{aligned} \hat{H} |\chi\rangle &= E |\chi\rangle \\ \text{ou } \hat{H} \chi &= E \chi \end{aligned} \quad \left\{ \begin{array}{l} \hat{H}_1 \chi^{(1)} = E_1 \chi^{(1)} \\ \hat{H}_2 \chi^{(2)} = E_2 \chi^{(2)} \end{array} \right.$$

As partículas 1 e 2 interagem com um agente externo
mas **não entre si !**

Notação :

$$\psi(1, 2) = \psi_s(1, 2) = |\uparrow\uparrow\rangle = |\uparrow\rangle |\uparrow\rangle = \chi_+^{(1)} \chi_+^{(2)}$$

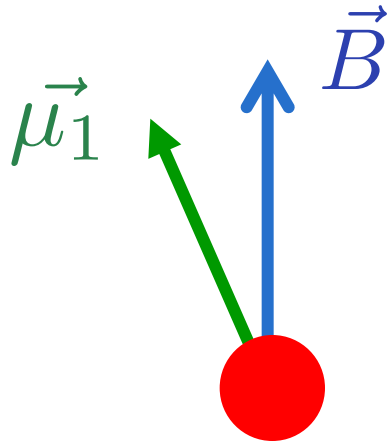


Estão em espaços diferentes !!!

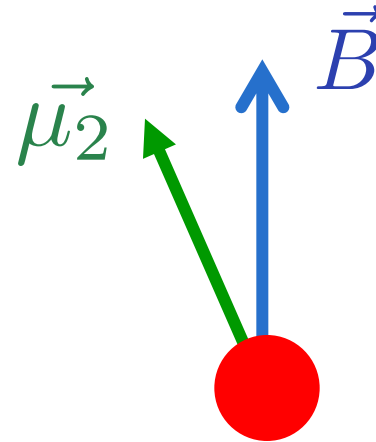
(espaço 1 da partícula 1 ; espaço 2 da partícula 2)

Se eu quiser falar com o spin ...

Preciso de um campo magnético



$$H_1 = -\vec{\mu}_1 \cdot \vec{B}$$



$$H_2 = -\vec{\mu}_2 \cdot \vec{B}$$

$$\hat{H} = \hat{H}_1 + \hat{H}_2$$

$$\chi = \chi^{(1)} \chi^{(2)}$$

$$\hat{H} \chi = (\hat{H}_1 + \hat{H}_2) \chi^{(1)} \chi^{(2)}$$

$$\hat{H} \chi = (\hat{H}_1 + \hat{H}_2) \chi^{(1)} \chi^{(2)} = \hat{H}_1 \chi^{(1)} \chi^{(2)} + \hat{H}_2 \chi^{(1)} \chi^{(2)}$$

$$= \hat{H}_1 \chi^{(1)} \chi^{(2)} + \chi^{(1)} \hat{H}_2 \chi^{(2)}$$

$$= E_1 \chi^{(1)} \chi^{(2)} + \chi^{(1)} E_2 \chi^{(2)}$$

$$= (E_1 + E_2) \chi^{(1)} \chi^{(2)} = E \chi^{(1)} \chi^{(2)}$$



estão em
espaços
diferentes !!!

$$\hat{H} \chi = E \chi$$

$$\left\{ \begin{array}{l} \hat{H} = \hat{H}_1 + \hat{H}_2 \\ \chi = \chi^{(1)} \chi^{(2)} \end{array} \right.$$

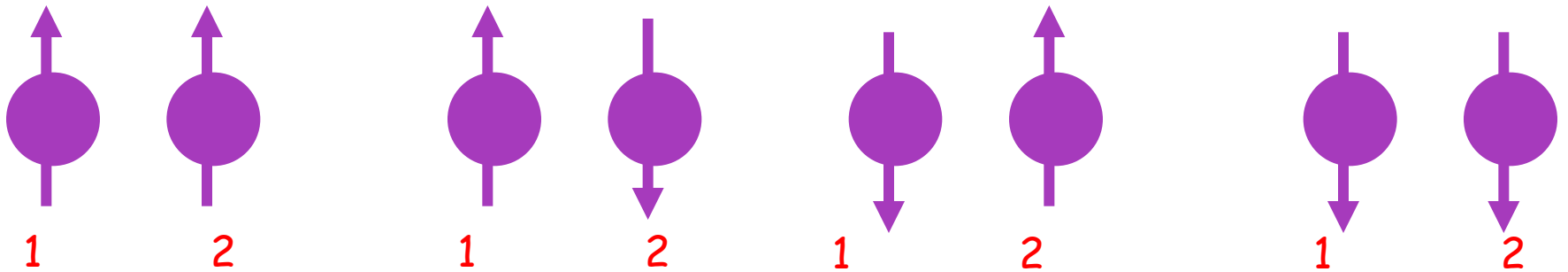


$$\hat{H} \chi = E \chi$$



Produto é solução da ESIT!

Projeção do spin total na direção z



$$4 = 3 + 1$$

$$S_z \chi_1 \chi_2 = (S_z^{(1)} + S_z^{(2)}) \chi_1 \chi_2$$

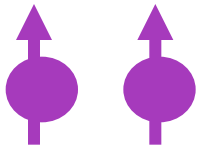
$$= S_z^{(1)} \chi_1 \chi_2 + S_z^{(2)} \chi_1 \chi_2$$

$$= m_1 \hbar \chi_1 \chi_2 + m_2 \hbar \chi_1 \chi_2$$

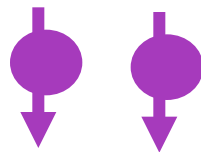
$$= (m_1 + m_2) \hbar \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = m \hbar \chi_1 \chi_2$$

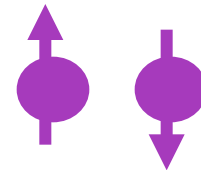
$$m = (m_1 + m_2)$$



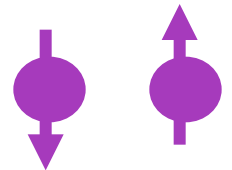
$$m = 1$$



$$m = -1$$



$$m = 0$$

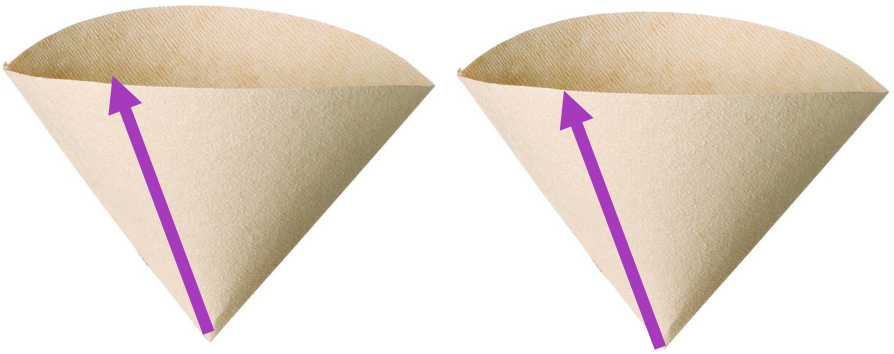


$$m = 0$$

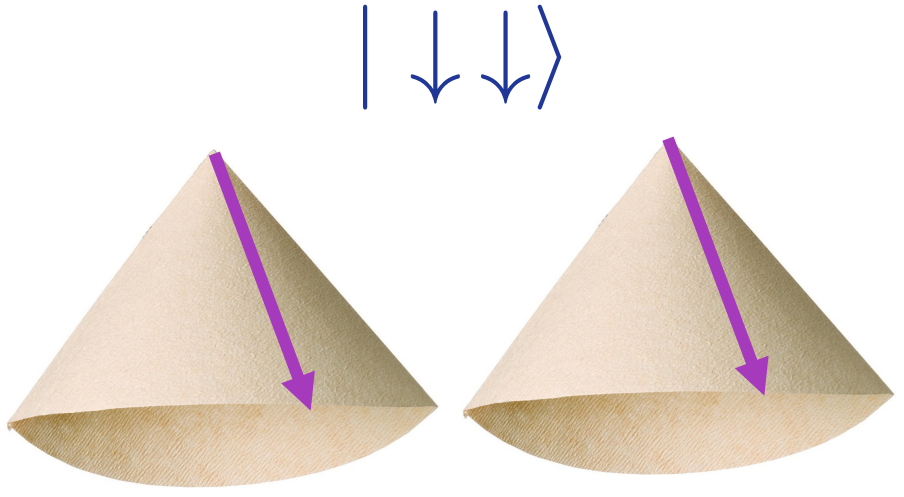
Vamos "ver" os spins das duas partículas

Faça você mesmo em casa !

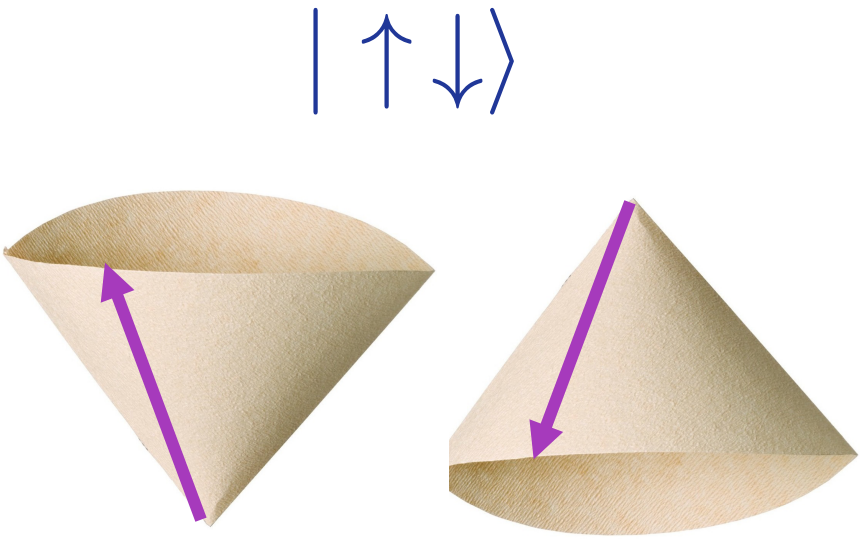




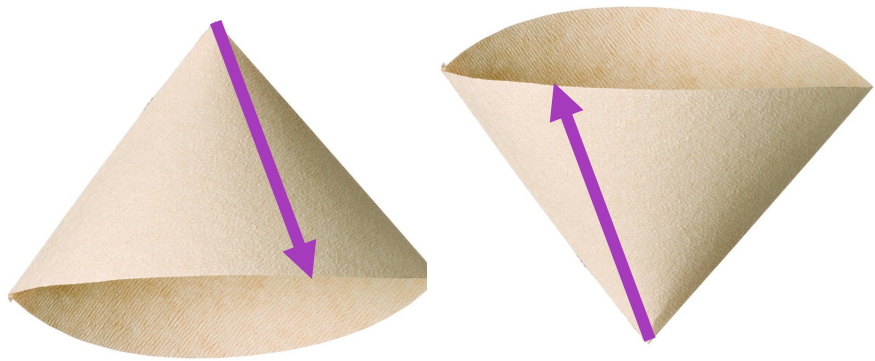
$|\uparrow\uparrow\rangle$



$|\downarrow\downarrow\rangle$



$|\uparrow\downarrow\rangle$



$|\downarrow\uparrow\rangle$

Vamos "abaixar" o spin do estado

$$|\uparrow\uparrow\rangle = \chi_{+}^{(1)} \chi_{+}^{(2)}$$

Lembramos que :

$$\hat{S}_+ |\downarrow\rangle = \hbar \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \hbar \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\hat{S}_+ |\downarrow\rangle = \hbar |\uparrow\rangle$$

$$\hat{S}_+ |\uparrow\rangle = \hbar \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \hbar \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$\hat{S}_+ |\uparrow\rangle = 0$$

$$\hat{S}_- |\uparrow\rangle = \hbar \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \hbar \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\hat{S}_- |\uparrow\rangle = \hbar |\downarrow\rangle$$

$$\hat{S}_- |\downarrow\rangle = \hbar \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \hbar \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$\hat{S}_- |\downarrow\rangle = 0$$

$$S_- = S_-^{(1)} + S_-^{(2)}$$

$$S_- \chi_+^{(1)} \chi_+^{(2)} = (S_-^{(1)} + S_-^{(2)}) \chi_+^{(1)} \chi_+^{(2)}$$

$$= S_-^{(1)} \chi_+^{(1)} \chi_+^{(2)} + S_-^{(2)} \chi_+^{(1)} \chi_+^{(2)}$$


$$= \hbar \chi_-^{(1)} \chi_+^{(2)} + \hbar \chi_+^{(1)} \chi_-^{(2)}$$

$$S_- |\uparrow\uparrow\rangle = \hbar (|\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle)$$

Vamos "abaixar" o spin do estado

$$|\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle$$

$$S_- \left[\hbar \chi_-^{(1)} \chi_+^{(2)} + \hbar \chi_+^{(1)} \chi_-^{(2)} \right] =$$

$$= S_-^{(1)} \left[\cancel{\hbar \chi_-^{(1)}} \chi_+^{(2)} + \hbar \chi_+^{(1)} \chi_-^{(2)} \right] +$$

$$+ S_-^{(2)} \left[\hbar \chi_-^{(1)} \chi_+^{(2)} + \cancel{\hbar \chi_+^{(1)}} \chi_-^{(2)} \right]$$

$$S_- \left[\hbar \chi_-^{(1)} \chi_+^{(2)} + \hbar \chi_+^{(1)} \chi_-^{(2)} \right] =$$

$$= \hbar^2 \chi_-^{(1)} \chi_-^{(2)} + \hbar^2 \chi_-^{(1)} \chi_-^{(2)} = 2 \hbar^2 \chi_-^{(1)} \chi_-^{(2)}$$

$$S_- \hbar (|\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle) = 2 \hbar^2 |\downarrow\downarrow\rangle$$

Três estados $|s, m_s\rangle$ conectados por S_- :

$$|1, 1\rangle \rightarrow |1, 0\rangle \rightarrow |1, -1\rangle$$

Este é o tripleto de spin 1

Vamos mostrar que o tripleto tem spin 1

$$S^2 = \vec{S} \cdot \vec{S} = \left(\vec{S}^{(1)} + \vec{S}^{(2)} \right) \cdot \left(\vec{S}^{(1)} + \vec{S}^{(2)} \right)$$

$$S^2 = \left(S^{(1)} \right)^2 + \left(S^{(2)} \right)^2 + 2 \vec{S}^{(1)} \cdot \vec{S}^{(2)}$$

$$= \left(S^{(1)} \right)^2 + \left(S^{(2)} \right)^2 + 2 \left(S_x^{(1)} \cdot S_x^{(2)} + S_y^{(1)} \cdot S_y^{(2)} + S_z^{(1)} \cdot S_z^{(2)} \right)$$

$$S^2 |1\ 1\rangle = S^2 \chi_+^{(1)} \chi_+^{(2)}$$

$$= \left(S^{(1)} \right)^2 \chi_+^{(1)} \chi_+^{(2)} + \left(S^{(2)} \right)^2 \chi_+^{(1)} \chi_+^{(2)} +$$

$$+ 2 \left(S_x^{(1)} \cdot S_x^{(2)} + S_y^{(1)} \cdot S_y^{(2)} + S_z^{(1)} \cdot S_z^{(2)} \right) \chi_+^{(1)} \chi_+^{(2)}$$

Vamos usar :

$$S_z \chi_+ = \frac{\hbar}{2} \chi_+$$

$$S_x \chi_- = \frac{\hbar}{2} \chi_+$$

$$S_y \chi_+ = -\frac{\hbar}{2i} \chi_-$$

$$S_z \chi_- = -\frac{\hbar}{2} \chi_-$$

$$S_x \chi_+ = \frac{\hbar}{2} \chi_-$$

$$S_y \chi_- = \frac{\hbar}{2i} \chi_+$$

$$S^2 \chi_+ = \frac{3}{4} \hbar^2 \chi_+$$

$$S^2 \chi_- = \frac{3}{4} \hbar^2 \chi_-$$

Exercício

$$= \left(S^{(1)} \right)^2 \chi_+^{(1)} \chi_+^{(2)} + \left(S^{(2)} \right)^2 \chi_+^{(1)} \chi_+^{(2)} +$$

$$+ 2 \left(S_x^{(1)} \cdot S_x^{(2)} + S_y^{(1)} \cdot S_y^{(2)} + S_z^{(1)} \cdot S_z^{(2)} \right) \chi_+^{(1)} \chi_+^{(2)}$$

$$= \left(S^{(1)}\right)^2 \chi_+^{(1)} \chi_+^{(2)} + \left(S^{(2)}\right)^2 \chi_+^{(1)} \chi_+^{(2)} +$$

$$+ 2 \left(S_x^{(1)} \cdot S_x^{(2)} + S_y^{(1)} \cdot S_y^{(2)} + S_z^{(1)} \cdot S_z^{(2)} \right) \chi_+^{(1)} \chi_+^{(2)}$$

$$= \frac{3}{4} \hbar^2 \chi_+^{(1)} \chi_+^{(2)} + \frac{3}{4} \hbar^2 \chi_+^{(1)} \chi_+^{(2)} +$$

$$+ 2 \cancel{\frac{\hbar}{2} \chi_-^{(1)} \frac{\hbar}{2} \chi_-^{(2)}} + 2(-1) \cancel{\frac{\hbar}{2i} \chi_-^{(1)} (-1) \frac{\hbar}{2i} \chi_-^{(2)}} + 2 \frac{\hbar}{2} \chi_+^{(1)} \frac{\hbar}{2} \chi_+^{(2)}$$

$$= \frac{3}{2} \hbar^2 \chi_+^{(1)} \chi_+^{(2)} + \frac{\hbar^2}{2} \chi_+^{(1)} \chi_+^{(2)} = 2\hbar^2 \chi_+^{(1)} \chi_+^{(2)}$$

$$\boxed{S^2 \chi_+^{(1)} \chi_+^{(2)} = 2\hbar^2 \chi_+^{(1)} \chi_+^{(2)}}$$

$$S^2 \chi_+^{(1)} \chi_+^{(2)} = 2\hbar^2 \chi_+^{(1)} \chi_+^{(2)}$$

Lembrando que

$$\left\{ \begin{array}{l} \hat{S}^2 \psi = s(s+1)\hbar^2 \psi \\ \hat{S}_z \psi = m_s \hbar \psi \end{array} \right.$$

$$S^2 \chi_+^{(1)} \chi_+^{(2)} = s(s+1)\hbar^2 \chi_+^{(1)} \chi_+^{(2)}$$

$$s(s+1) = 2$$

$$s = 1$$

Exercício: verifique que os estados abaixo tem spin 1

$$|\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle = \chi_{-}^{(1)} \chi_{+}^{(2)} + \chi_{+}^{(1)} \chi_{-}^{(2)}$$

$$|\downarrow\downarrow\rangle = \chi_{-}^{(1)} \chi_{-}^{(2)}$$

Exercício: verifique que o estado abaixo tem spin 0

$$|\downarrow\uparrow\rangle - |\uparrow\downarrow\rangle = \chi_{-}^{(1)} \chi_{+}^{(2)} - \chi_{+}^{(1)} \chi_{-}^{(2)}$$

Não são 4 combinações ! São 3 + 1 !

$$\left\{ \begin{array}{l} |\uparrow\uparrow\rangle \\ |\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle \\ |\downarrow\downarrow\rangle \end{array} \right.$$

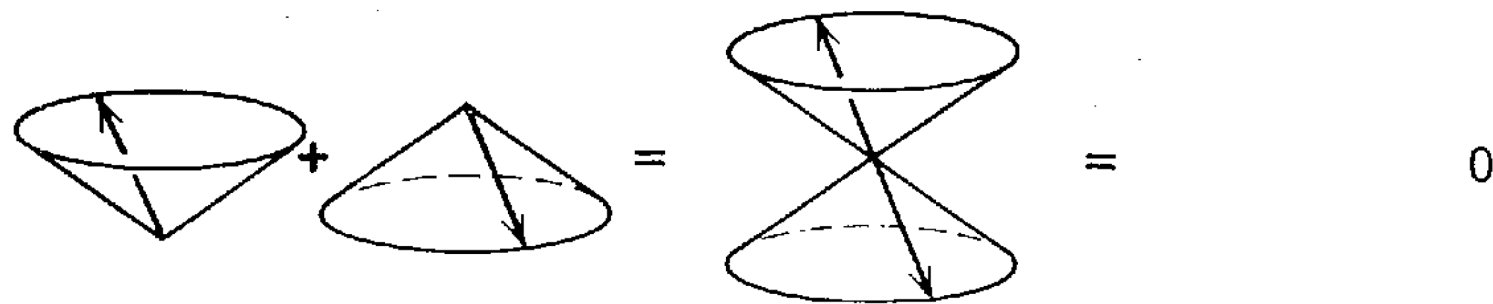
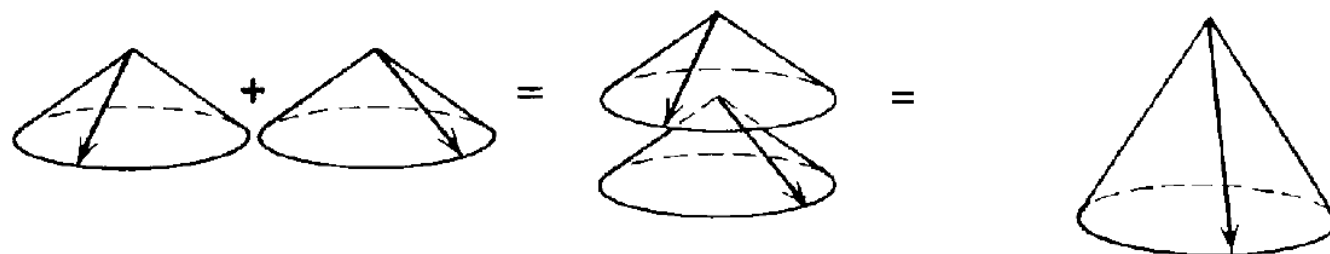
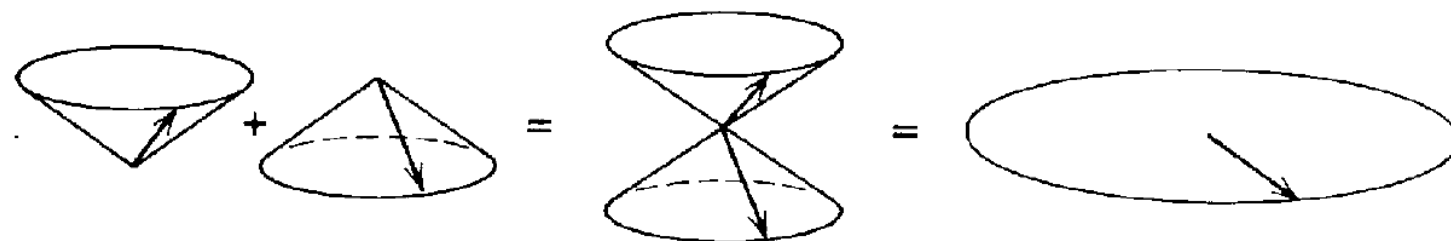
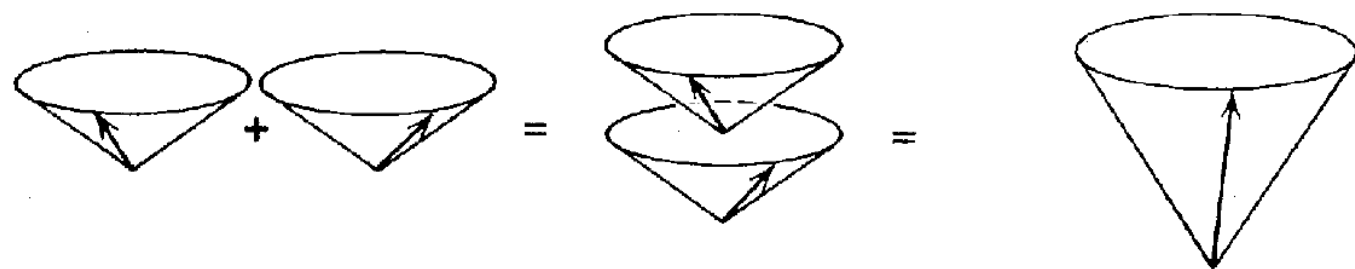
Tripleto de spin 1

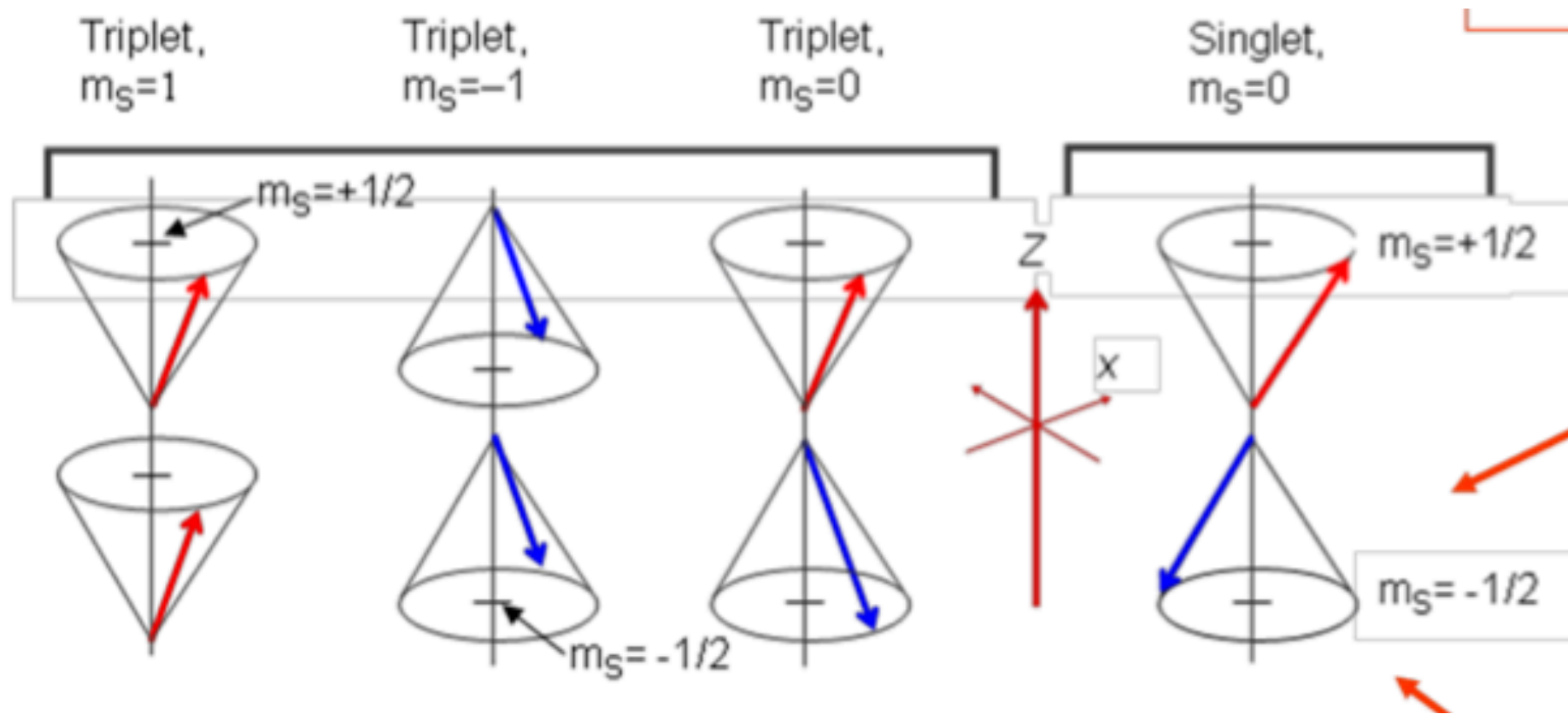
Conectados pelos
operadores escada

$$|\downarrow\uparrow\rangle - |\uparrow\downarrow\rangle$$

Singleto de spin 0

"aniquilado" pela escada





Simetria de troca da função de onda

$$1 \rightarrow 2 \qquad 2 \rightarrow 1$$

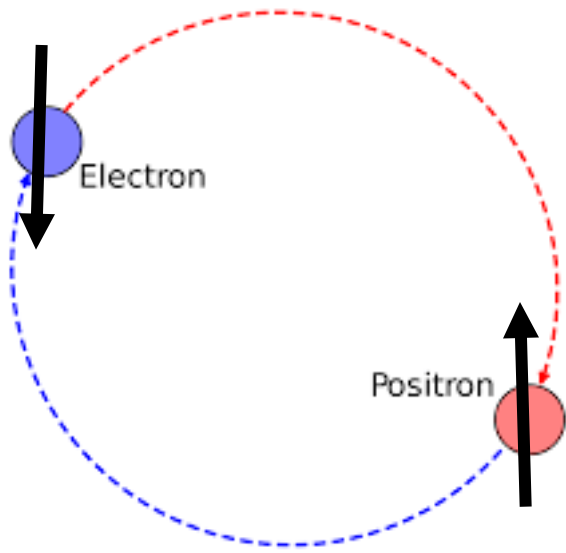
$$\left\{ \begin{array}{l} |\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle \rightarrow |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle \\ |\downarrow\downarrow\rangle \rightarrow |\downarrow\downarrow\rangle \\ |\uparrow\uparrow\rangle \rightarrow |\uparrow\uparrow\rangle \end{array} \right.$$

Tripleto é
simétrico !!!

$$|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \rightarrow |\downarrow\uparrow\rangle - |\uparrow\downarrow\rangle = - [|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle]$$

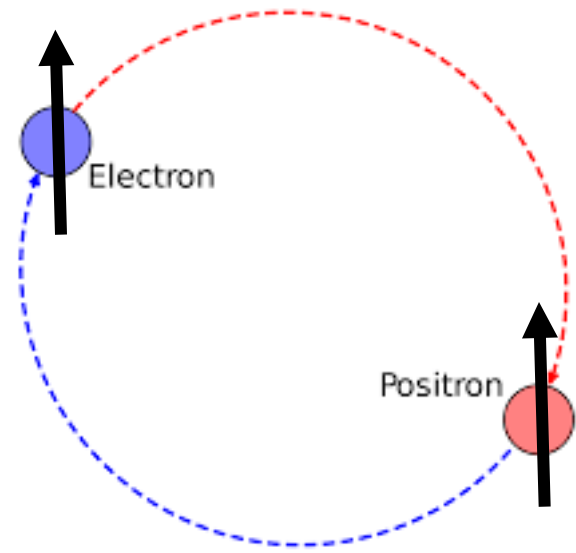
Singleto é antissimétrico !!!

Exemplo: positronium



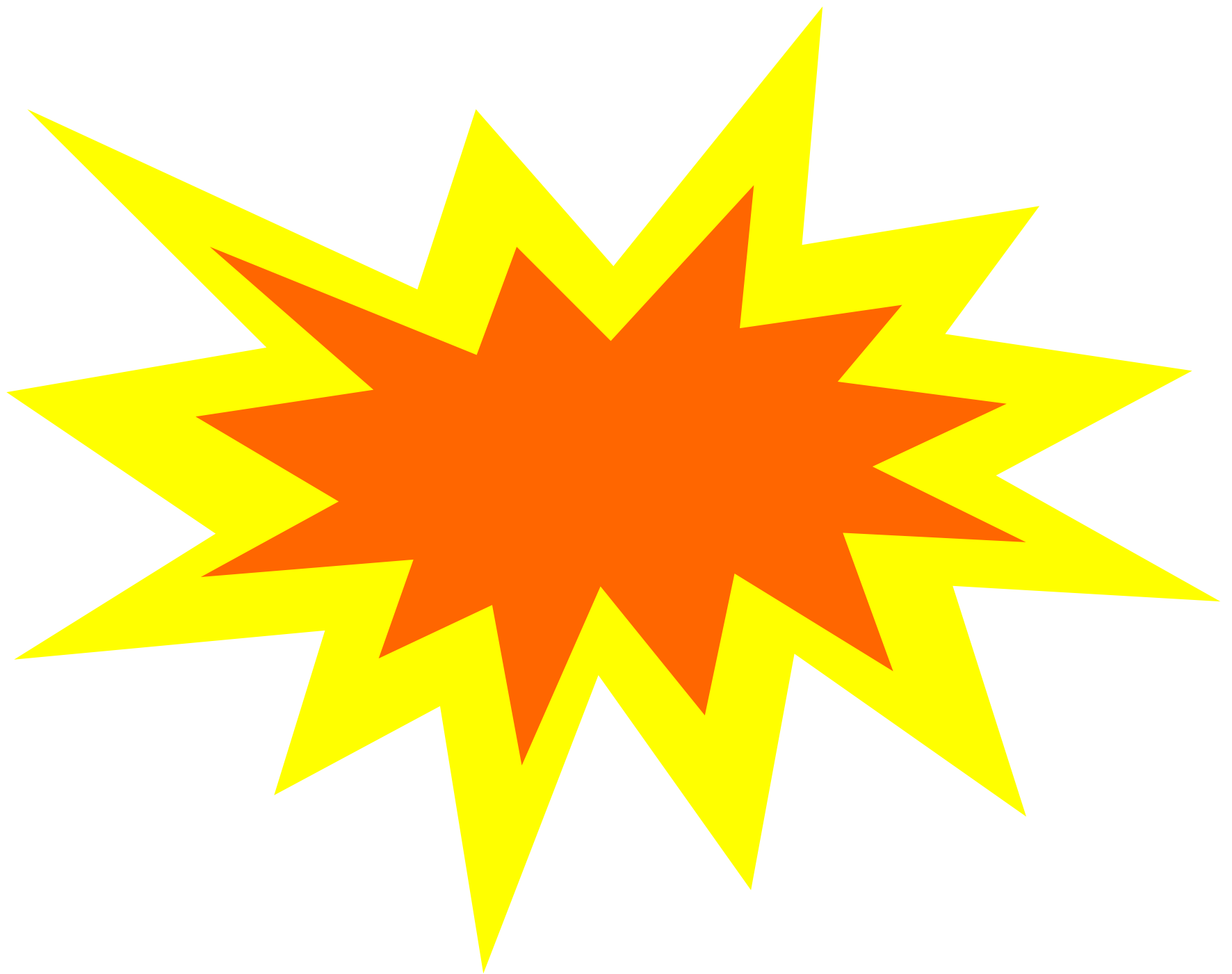
Para-positronium

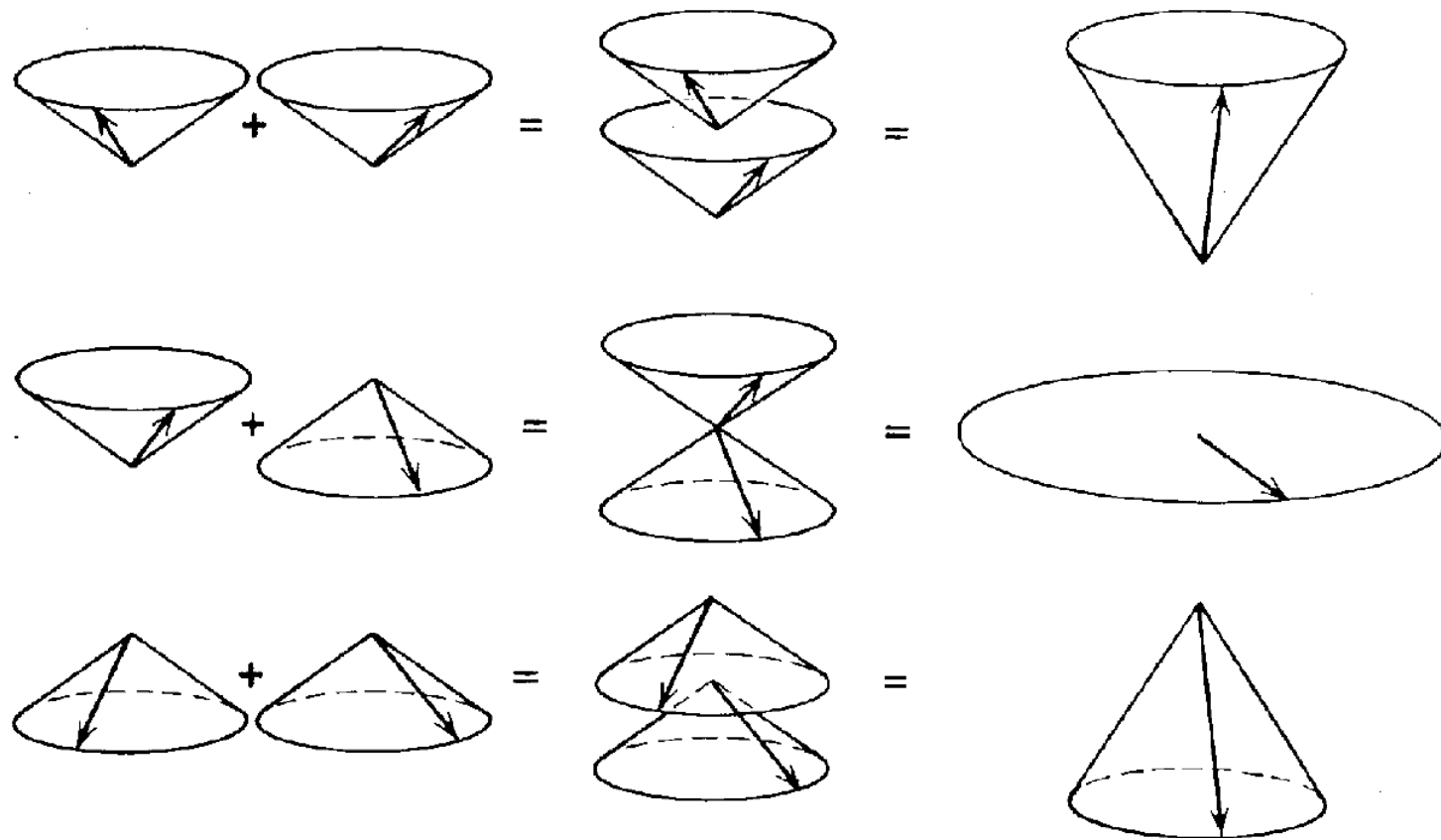
Singleto : $S = 0$



Orto-positronium

Tripleto : $S = 1$



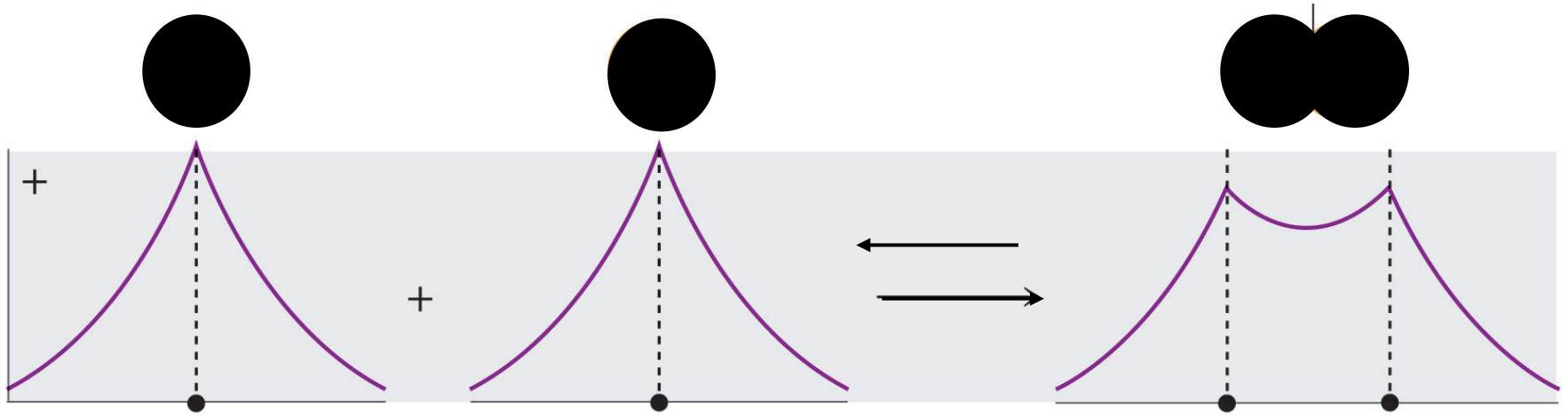


As três combinações estão conectadas pelo operador escada !

O conjunto é chamado **triplete** !

Função de onda de duas partículas

Parte espacial :

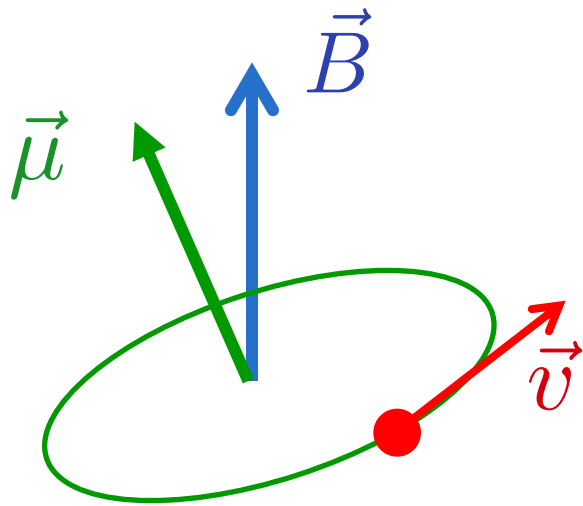


Cada partícula tem uma onda (extensa) associada

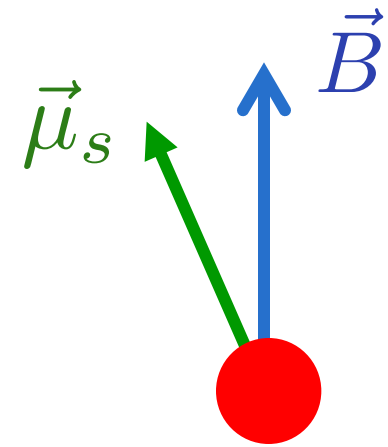
Quando dizer que dois corpos estão realmente separados ?

Partículas não interagentes: $\psi(x_1, x_2) = \psi_1(x_1) \psi_2(x_2)$





$$\vec{B} = B_z \hat{k}$$



$$U = -\vec{\mu} \cdot \vec{B}$$

Carga positiva

$$U = -g_l \mu_b m_l B_z$$

$$U = -g_s \mu_b m_s B_z$$

Carga negativa

$$U = +g_l \mu_b m_l B_z$$

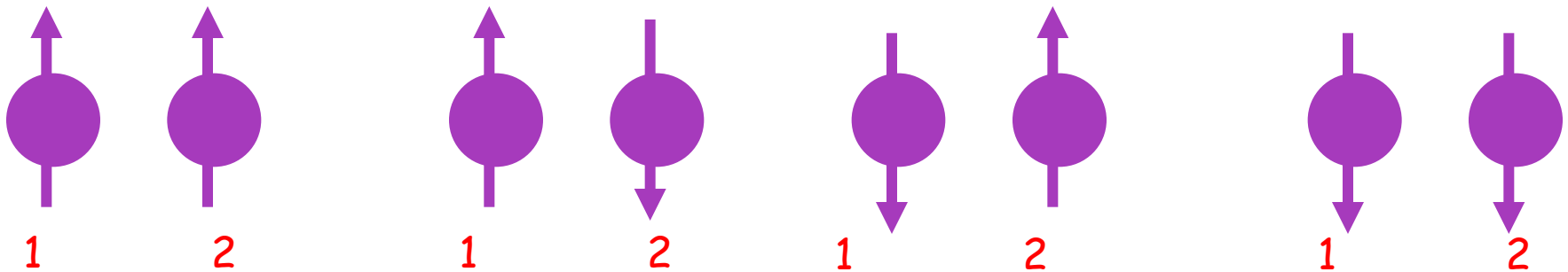
$$U = +g_s \mu_b m_s B_z$$

Adição de spins

Vamos somar os spins de duas partículas de spin 1/2

Por exemplo: spin total do átomo de hidrogênio

Projeções na direção z



Classicamente somamos os dois vetores de spin:

$$\vec{S} = \vec{S}^{(1)} + \vec{S}^{(2)}$$

Função de onda (autoestado) de spin de duas partículas :

$$|\chi\rangle = |\chi_1\rangle |\chi_2\rangle = \chi_1 \chi_2$$

Autovalor de S_z no estado composto :

$$\vec{S} = \vec{S}^{(1)} + \vec{S}^{(2)} \quad \longrightarrow \quad \hat{S} = \hat{S}^{(1)} + \hat{S}^{(2)}$$

(operadores)

$$\hat{S}_z = \hat{S}_z^{(1)} + \hat{S}_z^{(2)} \quad \longrightarrow \quad S_z = S_z^{(1)} + S_z^{(2)}$$

$$S_z \chi_1 \chi_2 = (S_z^{(1)} + S_z^{(2)}) \chi_1 \chi_2$$



$$S_z \chi_1 \chi_2 = (S_z^{(1)} + S_z^{(2)}) \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = S_z^{(1)} \chi_1 \chi_2 + S_z^{(2)} \chi_1 \chi_2$$

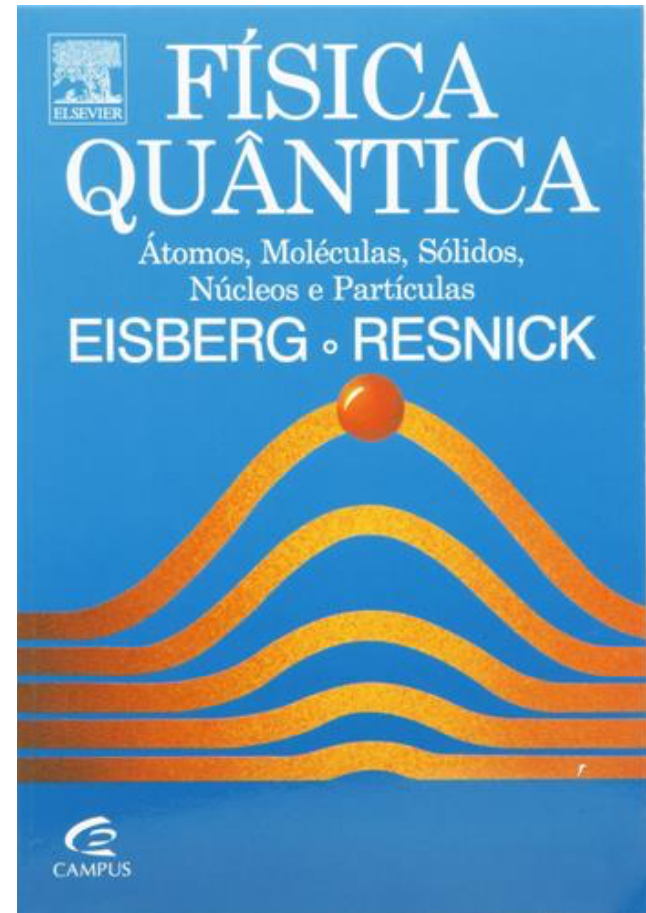
$$S_z \chi_1 \chi_2 = m_1 \hbar \chi_1 \chi_2 + m_2 \hbar \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = (m_1 + m_2) \hbar \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = m \hbar \chi_1 \chi_2$$

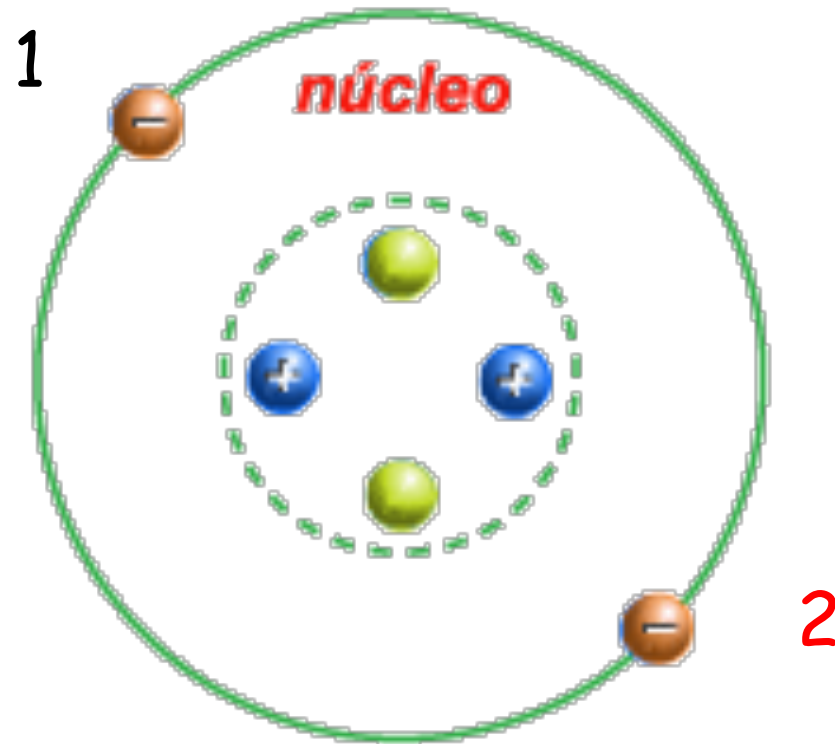
$\uparrow \uparrow$	$m = 1$	$\downarrow \uparrow$	$m = 0$
$\uparrow \downarrow$	$m = 0$	$\downarrow \downarrow$	$m = -1$

Sistemas de duas partículas



Capítulo 9

Sistemas de duas partículas



átomo de hélio

T_1 = energia cinética de 1

V_1 = energia potencial entre 1 e N

T_2 = energia cinética de 2

V_2 = energia potencial entre 2 e N

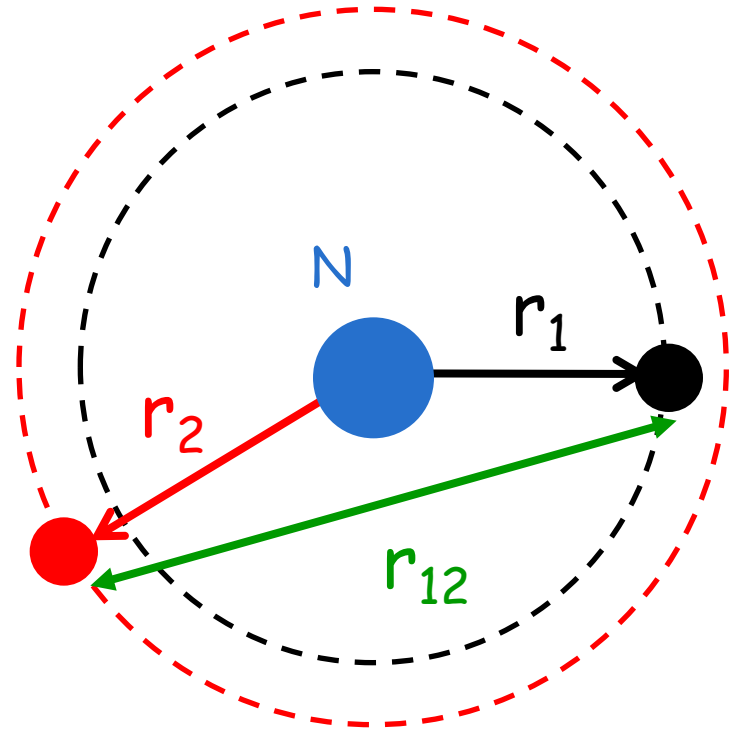
V_{12} = energia potencial entre 1 e 2

$$H_1 = T_1 + V_1$$

$$H_2 = T_2 + V_2$$

$$H_{12} = V_{12}$$

$$H = H_1 + H_2 + H_{12}$$



T_1 = energia cinética de 1

V_1 = energia potencial entre 1 e N

T_2 = energia cinética de 2

V_2 = energia potencial entre 2 e N

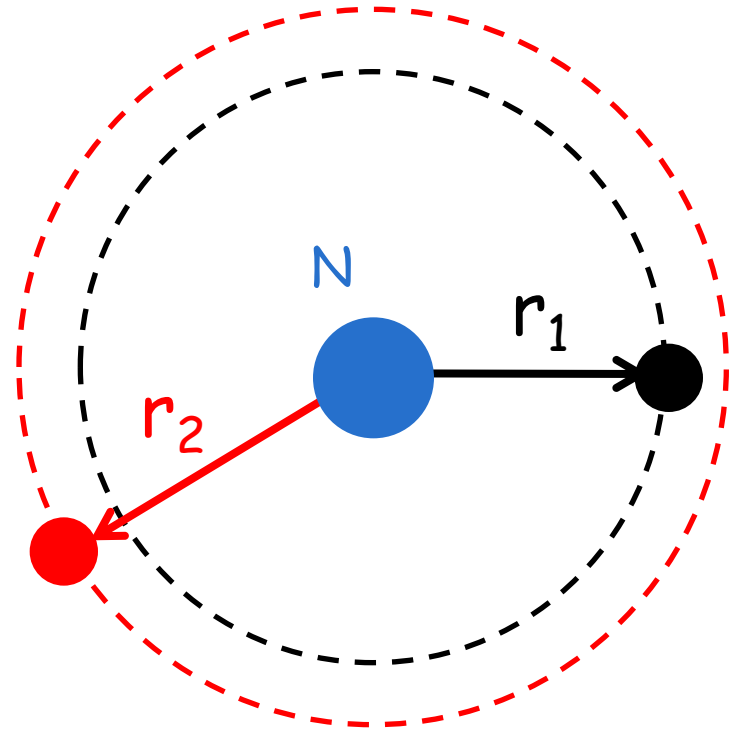
V_{12} = energia potencial entre 1 e 2

$$H_1 = T_1 + V_1$$

$$H_2 = T_2 + V_2$$

$$H_{12} = 0$$

$$H = H_1 + H_2$$



Desprezamos a interação
entre os dois eletrons !!!

Função Hamiltoniana



Operador Hamiltoniano

$$H = H_1 + H_2$$



$$\left\{ \begin{array}{l} \hat{H}_1 = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_1^2} + V_1 \\ \hat{H}_2 = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_2^2} + V_2 \end{array} \right.$$

$$\hat{H} \psi = E \psi$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x_1, x_2)}{\partial x_1^2} - \frac{\hbar^2}{2m} \frac{\partial^2 \psi(x_1, x_2)}{\partial x_2^2} +$$

$$+ V_1 \psi(x_1, x_2) + V_2 \psi(x_1, x_2) = E \psi(x_1, x_2)$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x_1, x_2)}{\partial x_1^2} - \frac{\hbar^2}{2m} \frac{\partial^2 \psi(x_1, x_2)}{\partial x_2^2} + V_1 \psi(x_1, x_2) + V_2 \psi(x_1, x_2) = E \psi(x_1, x_2)$$

Ansatz : $\psi(x_1, x_2) = \psi_1(x_1) \psi_2(x_2)$



Substituimos $\psi(x_1, x_2) = \psi_1(x_1) \psi_2(x_2)$ **e dividimos por** $\psi_1(x_1) \psi_2(x_2)$

$$-\frac{\hbar^2}{2m} \frac{1}{\psi_1(x_1)} \frac{\partial^2 \psi_1(x_1)}{\partial x_1^2} + V_1 - \frac{\hbar^2}{2m} \frac{1}{\psi_2(x_2)} \frac{\partial^2 \psi_2(x_2)}{\partial x_2^2} + V_2 = E$$

Passamos as funções de x_2 para o lado direito :

$$-\frac{\hbar^2}{2m} \frac{1}{\psi_1(x_1)} \frac{\partial^2 \psi_1(x_1)}{\partial x_1^2} + V_1(x_1) = E + \frac{\hbar^2}{2m} \frac{1}{\psi_2(x_2)} \frac{\partial^2 \psi_2(x_2)}{\partial x_2^2} - V_2(x_2)$$

$$-\frac{\hbar^2}{2m} \frac{1}{\psi_1(x_1)} \frac{\partial^2 \psi_1(x_1)}{\partial x_1^2} + V_1(x_1) = E + \frac{\hbar^2}{2m} \frac{1}{\psi_2(x_2)} \frac{\partial^2 \psi_2(x_2)}{\partial x_2^2} - V_2(x_2)$$

$f(x_1) = g(x_2)$ verdade para qualquer x_1 e x_2 se $f(x_1) = g(x_2) = cte$

$$-\frac{\hbar^2}{2m} \frac{1}{\psi_1(x_1)} \frac{\partial^2 \psi_1(x_1)}{\partial x_1^2} + V_1(x_1) = E + \frac{\hbar^2}{2m} \frac{1}{\psi_2(x_2)} \frac{\partial^2 \psi_2(x_2)}{\partial x_2^2} - V_2(x_2) = E_1$$

$$\left\{ \begin{array}{l} -\frac{\hbar^2}{2m} \frac{1}{\psi_1(x_1)} \frac{\partial^2 \psi_1(x_1)}{\partial x_1^2} + V_1(x_1) = E_1 \\ -\frac{\hbar^2}{2m} \frac{1}{\psi_2(x_2)} \frac{\partial^2 \psi_2(x_2)}{\partial x_2^2} + V_2(x_2) = E - E_1 = E_2 \end{array} \right.$$

estas equações
têm solução
conhecida !!!

Encontramos $\psi(x_1, x_2) = \psi_1(x_1) \psi_2(x_2)$ e $E = E_1 + E_2$