9.2 – Achieving Load Carrying Capacity: Theoretical and Practical Stability



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Coworker: S. Lenci, M. Thompson

D	AY	TIME	LECTURE
Monday	05/11	14.00 -14.45	Historical Framework - A Global Dynamics Perspective in the Nonlinear Analysis of Systems/Structures
		15.00 -15.45	Achieving Load Carrying Capacity: Theoretical and Practical Stability
		16.00 -16.45	Dynamical Integrity: Concepts and Tools_1
Wednesday	07/11	14.00 -14.45	Dynamical Integrity: Concepts and Tools_2
		15.00 -15.45	Global Dynamics of Engineering Systems
		16.00 -16.45	Dynamical integrity: Interpreting/Predicting Experimental Response
Monday	12/11	14.00 -14.45	Techniques for Control of Chaos
		15.00 -15.45	A Unified Framework for Controlling Global Dynamics
		16.00 -16.45	Response of Uncontrolled/Controlled Systems in Macro- and Micro-mechanics
Wednesday	14/11	14.00 -14.45	A Noncontact AFM: (a) Nonlinear Dynamics and Feedback Control (b) Global Effects of a Locally-tailored Control
		15.00 -15.45	Exploiting Global Dynamics to Control AFM Robustness
		16.00 -16.45	Dynamical Integrity as a Novel Paradigm for Safe/Aware Design

Outline

- 1. Main stability concepts at a glance
- 2. Local versus global safety in statics and dynamics
- 3. Solution/attractor robustness in phase space; the relevant 'safe' basins (through an archetypal model)
- 4. Solution/attractor robustness and basin compactness in control parameter space (through an archetypal model)
- 5. Robustness/erosion profiles
- 6. Moving from theoretical to practical stability

Achieving load carrying capacity

- Load carrying capacity: an old issue associated with the concept of loss of stability
- Stability: to be discussed by also considering the effects of (*static* or *dynamic*) imperfections, always present in nature/technology → A system must be able to sustain changes in both initial conditions and control parameters, without changing its desired outcome
- Robustness: a fundamental issue in analysis and design
- Dynamic integrity: a global safety concept essential to secure practical stability of systems
- Historical concepts and contributions at a glance

Leonhard Euler (1707-1783)

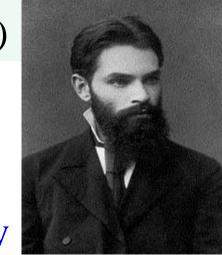
• First fundamental contribution: Euler buckling load of a column



- Loss of load carrying capacity identified as the system instability occurring at the local bifurcation point of an equilibrium path when changing a control parameter (axial load) talking, of course, in modern language
- A substantially **static** notion of stability

Aleksander Lyapunov (1857-1918)

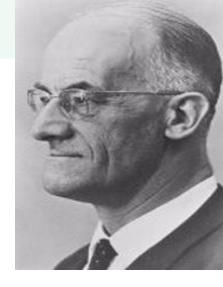
 Rigorous formulation within a more dynamically oriented notion of stability



- Lyapunov (or classical) local stability roughly states that under infinitesimal changes in initial conditions
 - the system must keep the reference response
- Major role in the solutions of a variety of engineering problems ensuing from modern technological developments

Warner Koiter (1914-1997)

Within the mechanical community, looking at the effects of changes of control parameters:



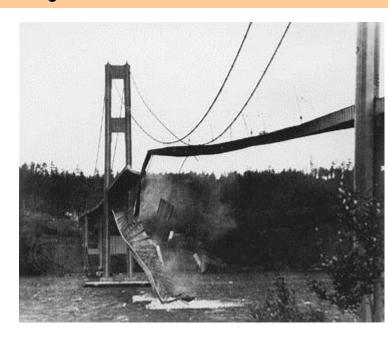
- Koiter realized that model imperfections are crucial in lowering the critical load
 - Due to imperfections, the branching point becomes a snap point, which (in the dangerous cases) occurs at a lower load threshold
- Dynamical character of stability was clear, but the reference framework was still 'static'

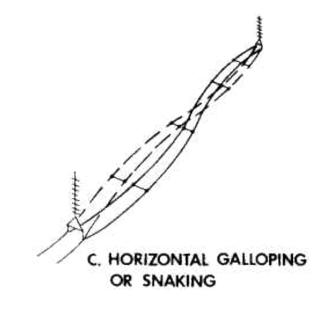
Structural stability

- Later on, **bifurcation theory** provided a mathematical background to this engineering intuition:
 - transcritical and pitchfork bifurcations (branching) are structurally unstable (i.e., unobservable in the real world, unless somehow forcing them) and become saddle-node bifurcations (snap) after system perturbations (imperfections in mechanical language)
- Structural stability: studying the effect of perturbations of the system with respect to parameters and not w.r.t. initial conditions, as in classical local stability

Dynamic stability

- When 'flutter' or 'galloping'
 of real systems came into
 play, dynamics definitely
 entered the concept of loss of
 stability
- In bifurcation theory language, the Hopf bifurcation was 'discovered' and experimentally observed, according to the fact that it is structurally stable





From theoretical to practical stability

Classical stability: small changes of initial conditions do not affect substantially the system response

Key point: how small have to be perturbations?

From a mathematical point of view the magnitude of perturbations is not important (e.g. 10^{-50} is ok)

But from a **practical** point of view it is important, since in our real world **imperfections** have a **finite** magnitude



Local (or classical, Lyapunov) stability is not enough for practical applications!!

Michael Thompson (1937-)

Practical stability of attractors to be addressed in an actually dynamical environment



Around the 90s:

- By considering a global approach, notion of dynamic integrity introduced, which is fundamental for properly pursuing the safety of structures
- Basins of attraction and their variation with a varying control parameter - become fundamental tools

Solution robustness in phase space

Properly complementing the solely local theoretical character of the classical concept of stability with a global practical one

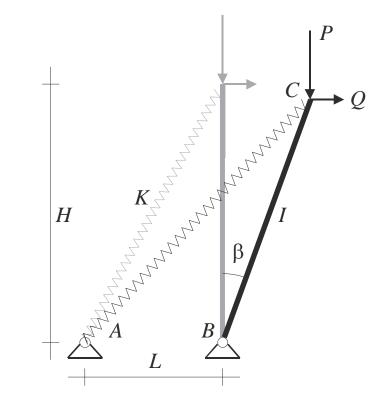
Already in the **static** case:

approaching a (local) bifurcation, the **basin** of reference solution shrinks to zero and becomes **unsafely small**, although the solution is still **stable** in the sense of Lyapunov

⇒ pursued response non-robust with respect to finite dynamic perturbations, though being its basin integer (no fractality)

An archetypal asymmetric model

- Single-dof mechanical model typically used to illustrate post-buckling behavior and imperfection sensitivity of structural systems liable to unstable buckling
- Q = "static" imperfection



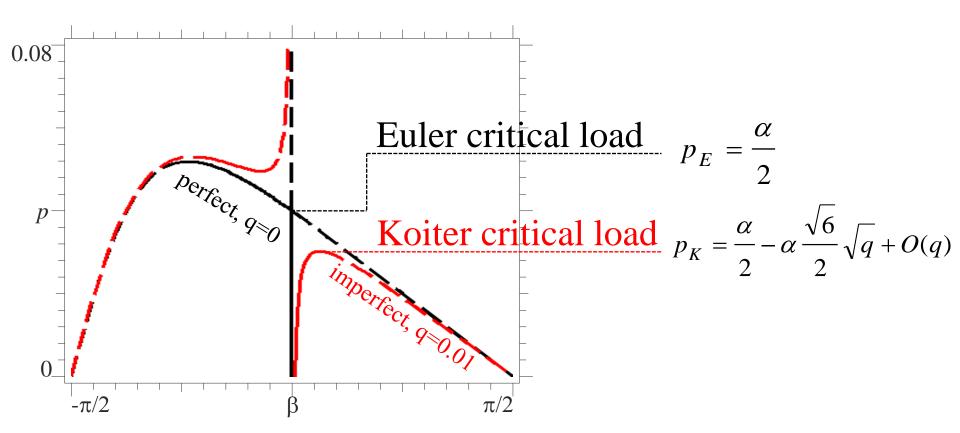
No damping and no dynamic excitation

$$\ddot{\beta} - p\sin(\beta) + \left[1 - \frac{1}{\sqrt{1 + \alpha\sin(\beta)}} - q\right] \cos(\beta) = 0$$

$$\alpha = \frac{2LH}{L^2 + H^2} \in [0,1]$$
In the following $\alpha = 0.8$

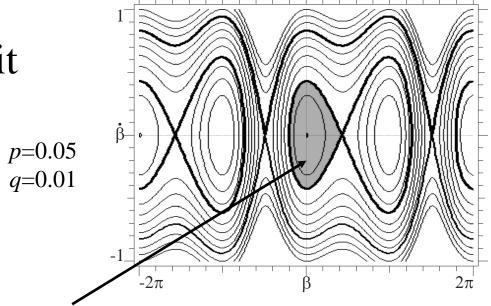
Equilibrium points and critical loads

$$p\sin(\beta) = \left[1 - \frac{1}{\sqrt{1 + \alpha \sin(\beta)}} - q\right] \cos(\beta)$$



Global safety

Phase portrait



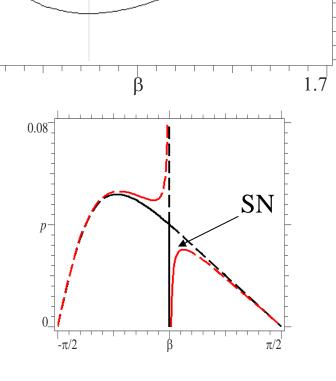
- 'basin of attraction' of equilibrium point
- The larger the area, the larger the ability of the system to support *finite* changes in i.c. \rightarrow

the larger the safety of the structure

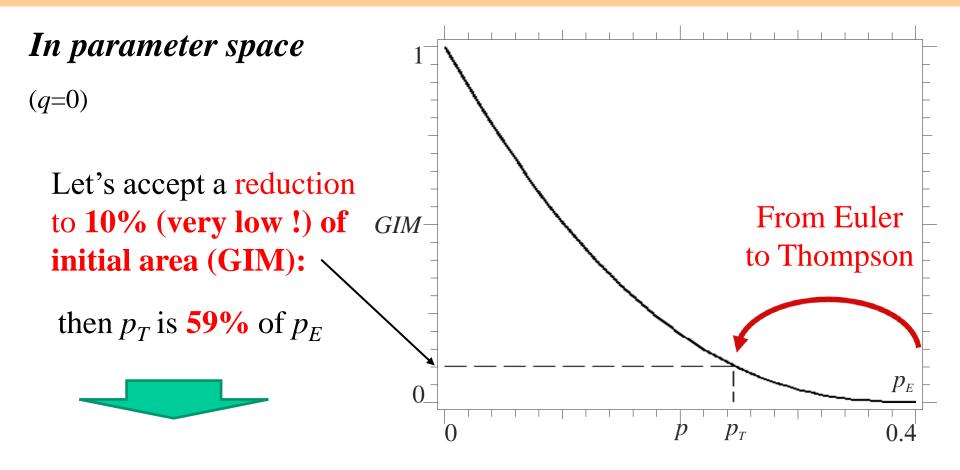
Basin and actual critical load reduction

-0.8

- Area shrinks as approaching critical load, with or without 'static' imperfections, and 'rapidly' becomes too small for real world, where *finite* dynamic imperfections exist
- Basin of attraction' under (even transient) dynamic perturbations shrinks to the attractor
 - → (Koiter) SN bifurcation <u>overestimates</u> the actual <u>critical load</u>



Area decrement without imperfections



• In the neighbourhood of p_E the safe region is merely residual and unsafe \longrightarrow practical p_T ('Thompson') critical load much less than p_E (Euler)

Area decrement with 'static' imperfections

Same qualitative behaviour

 $(q\neq 0)$

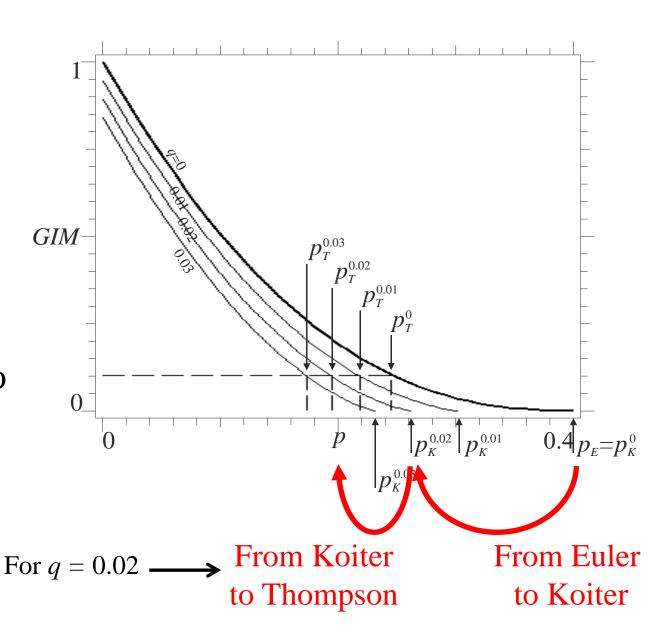
practical p_T

('Thompson')

critical load also

lower than p_K

(Koiter)



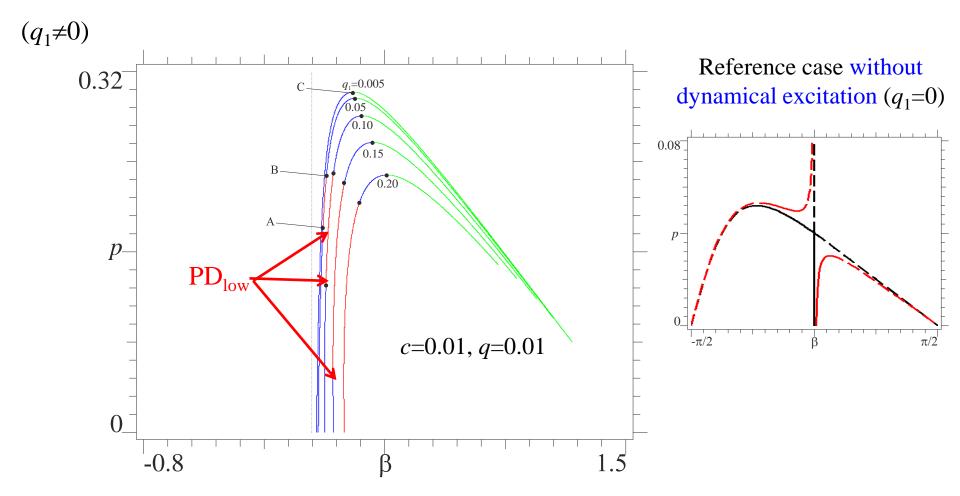
With dynamic excitation

• What happens when a dynamic excitation is applied, e.g., $q+q_1 \sin(\omega t)$?

$$\ddot{\beta} + c\dot{\beta} - p\sin(\beta) + \left[1 - \frac{1}{\sqrt{1 + \alpha\sin(\beta)}} - (q + q_1\sin(\omega t))\right]\cos(\beta) = 0$$

- The phase space augments of one dimension, but this is not a problem, and can be overcome, e.g., by considering Poincaré sections
- Also damping is added for realistic engineering analysis

Periodic solutions

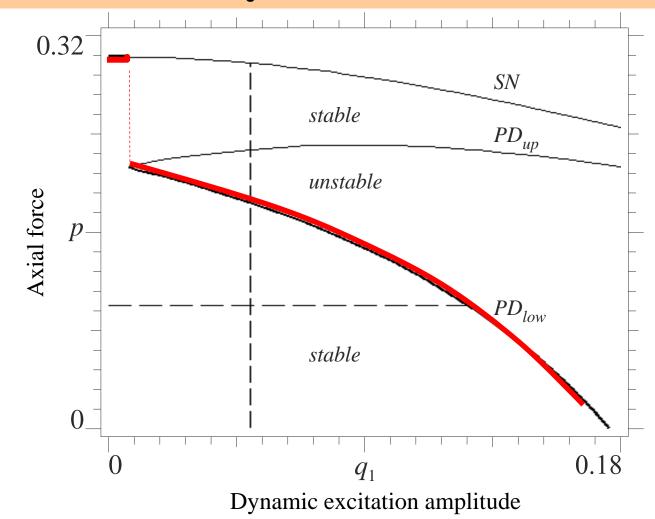


- The saddle-node (SN) decreases by increasing q_1
- A period doubling (PD) reduces the stability threshold (above PD_{low} the solution may jump out of well)

Stability threshold with dynamic excitation

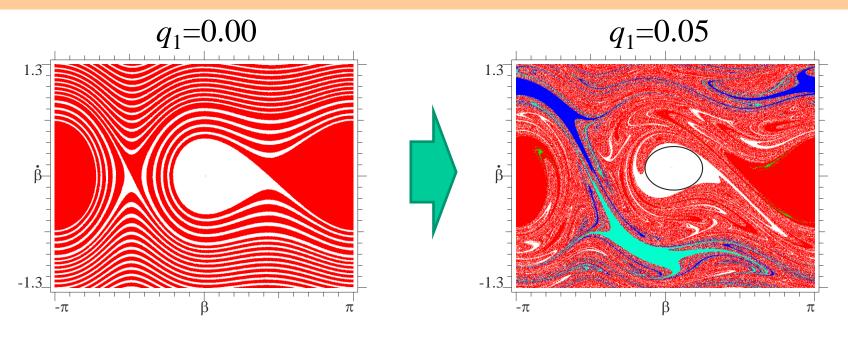
 $q+q_1\sin(\omega t)$

- periodic,
- quasiperiodic,
- chaotic attractors



• Interaction between static (p) and dynamic (q_1) loads causes meaningful loss of load carrying capacity (w.r.t. Koiter one)

Fractalization

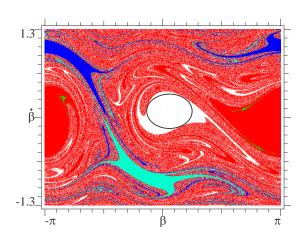


- Existence/competition of more attractors
- Basin of attraction **no longer safe** against small but **finite** incidental changes of i.c.
- Basin is eroded and loses its compactness/integrity
- Load carrying capacity depends on practical stability under imperfections/perturbations

Major effects of dynamic excitations

- Attractors are no longer equilibrium points, but periodic, quasi-periodic, ... chaotic orbits
- The topology of the basins of attraction changes significantly; **fractality** commonly appears
- Dynamic integrity:

 → a major role
 in determining the
 load carrying capacity



Practical stability under imperfections/perturbations

W.r. to <u>dynamic</u> imperfections: <u>initial conditions</u> in <u>phase</u> space solution/attractor robustness and basin properties

Static solution: robust if large safe basin

Dynamic attractor: - robust if large and compact (i.e. integer) basin

- non-robust if large but fractal basin

W.r. to **system** imperfections: **parameters** in **control space**

how solution/attractor robustness and basin compactness in phase space evolve with a varying control parameter

Static solution: robustness profile of safe basin

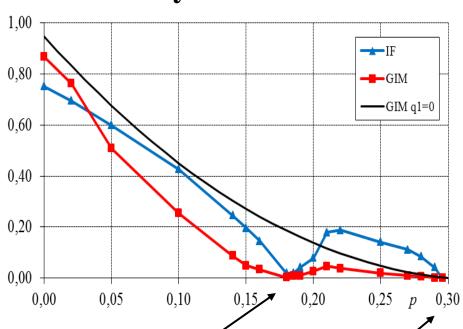
Dynamic attractor: - robustness profiles of (integer) competing basins

- erosion profile with integrity reduction

Robustness profiles: size reduction/increase of integer basin vs competing one **Erosion profiles**: reduction of basin integrity, to be explained also in terms of global bifurcation phenomena (homo/heteroclinic tangencies, crises, etc.)

Robustness profile

Increasing axial load, fixed dynamic excitation

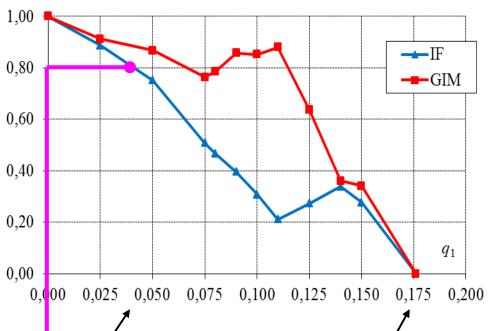


Practical (Thompson)
 stability threshold about 1/3
 of theoretical (Koiter)

critical load

Erosion profile

Increasing dynamic excitation, fixed axial load



Practical (Thompson) load carrying capacity much lower than Koiter one, e.g.:

lower than Koller one, e.g.:

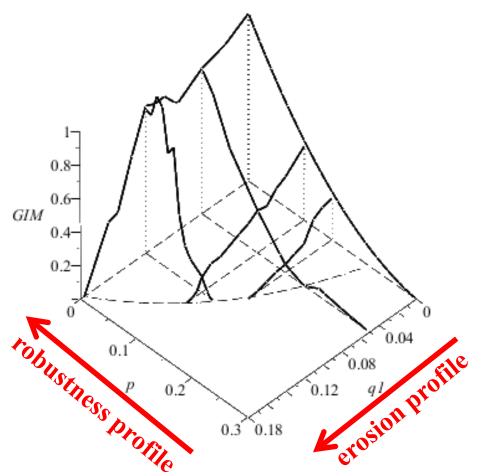
esidual F=80% (practically uneroded basins) \rightarrow

Thompson threshold = 22 % Koiter threshold

Hints for design

- Koiter load can be determined upon fixing the value of the expected static imperfection q
- Thompson load can be determined upon fixing the acceptable minimal integrity
 - (which corresponds to fixing the maximum allowed change in i.c. that can be safely supported by the system; in other words, this corresponds to fixing the "safety factor")
- Both Koiter and Thompson theories are thus 'applicable' with the knowledge of *q* and *GIM*

A summary interaction picture



- dynamic excitation reduces Koiter practical critical load
- static axial load
 reduces Thompson
 escape dynamic
 excitation

- Interaction of static axial load and dynamic excitation
- Dangerously residual robustness/compactness occurs well before disappearance of solution/attractor

Theoretical vs practical stability

end of **robustness/erosion** profile corresponds to attractor disappearance, i.e. to **loss of stability**

