

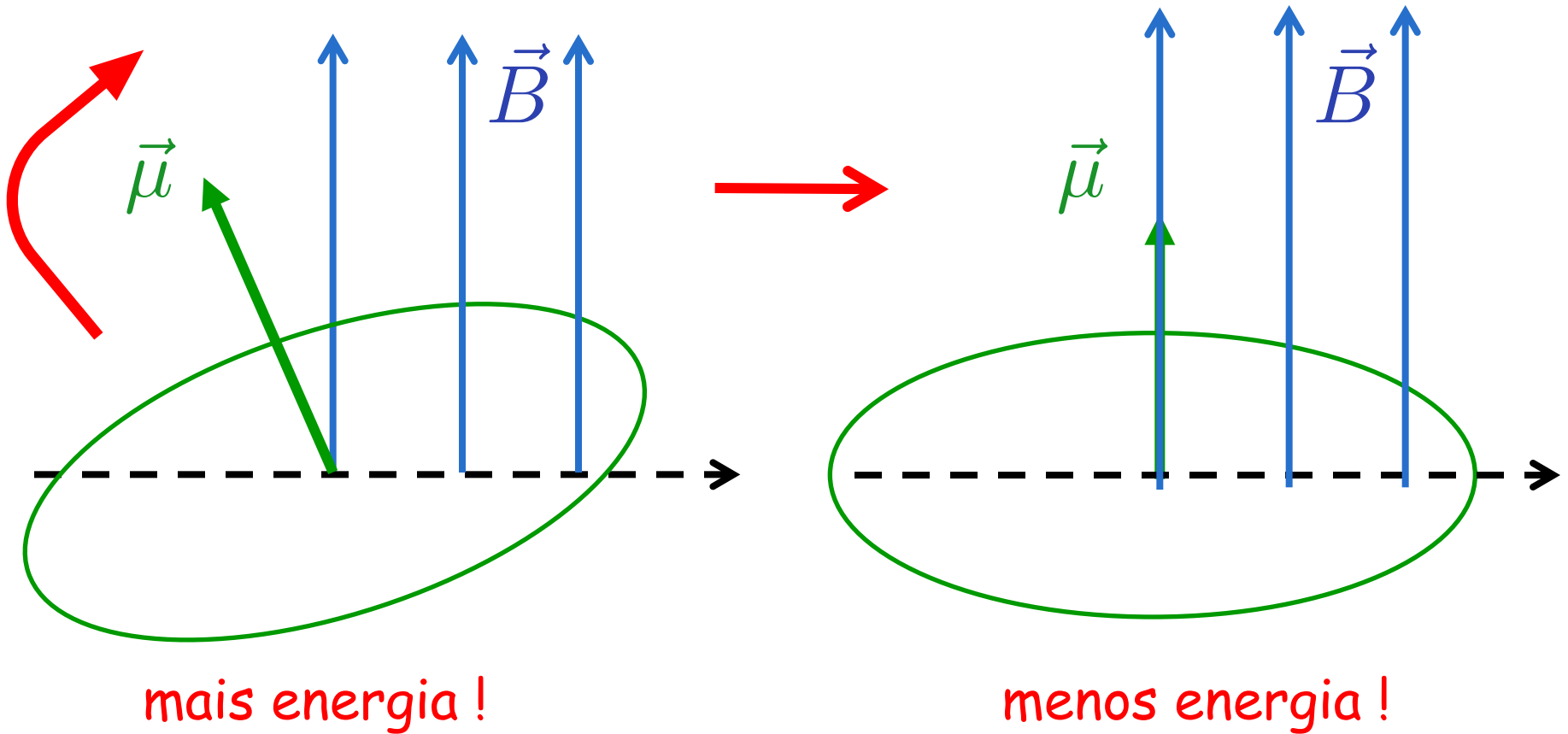
Aula 8

Valor médio: precessão de Larmor

Adição de spins

Exemplos

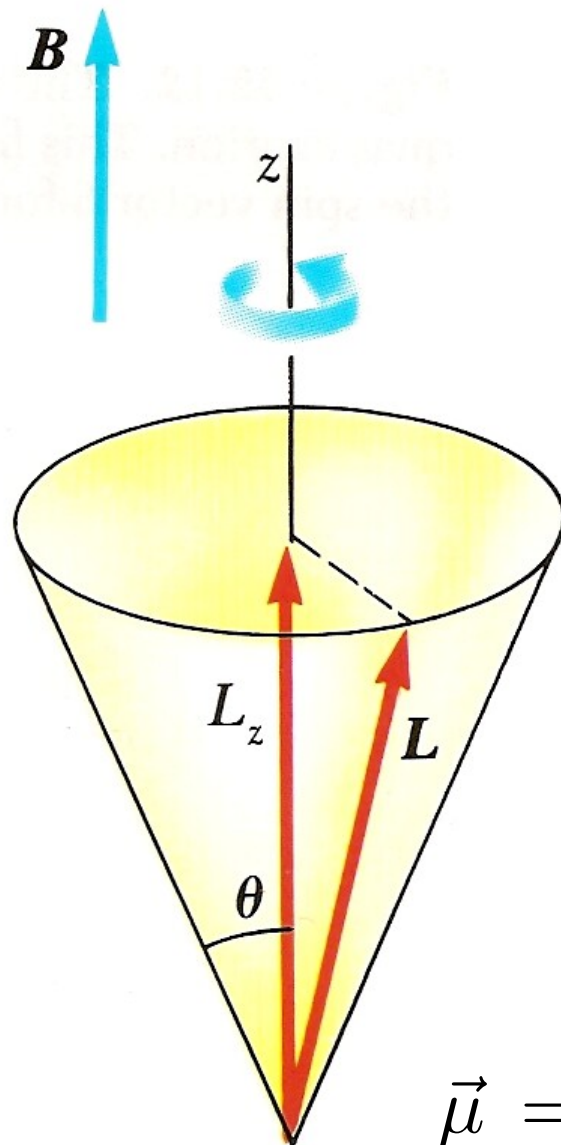
Espira tenta se alinhar como uma bússola !



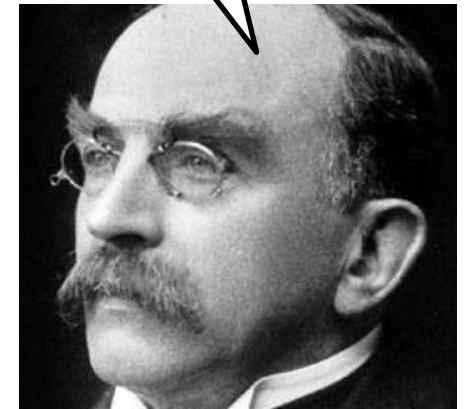
Energia potencial de orientação

$$U = -\vec{\mu} \cdot \vec{B}$$

Mas nunca se alinha ! Precessiona !



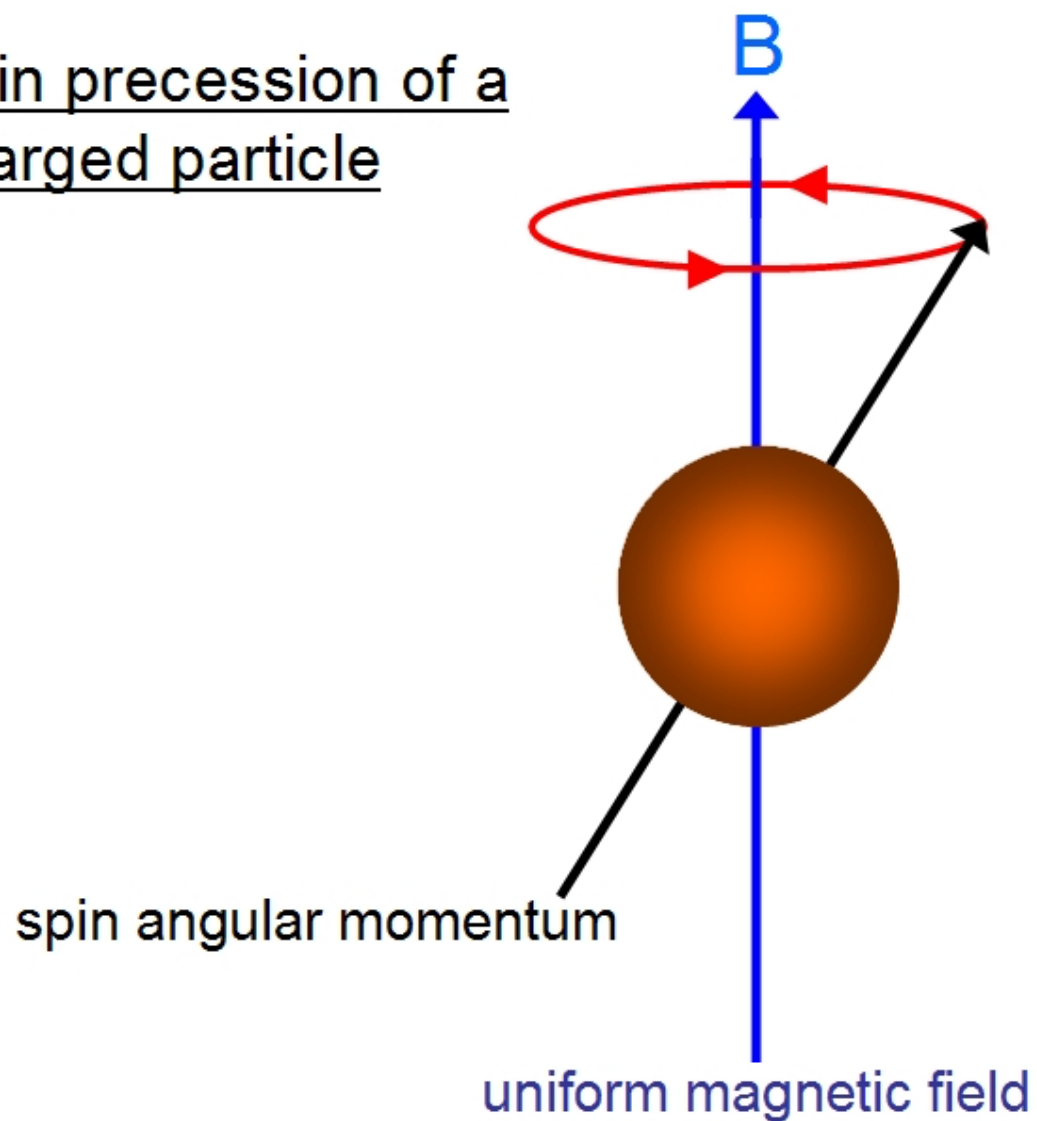
Precessão
de Larmor



Joseph Larmor

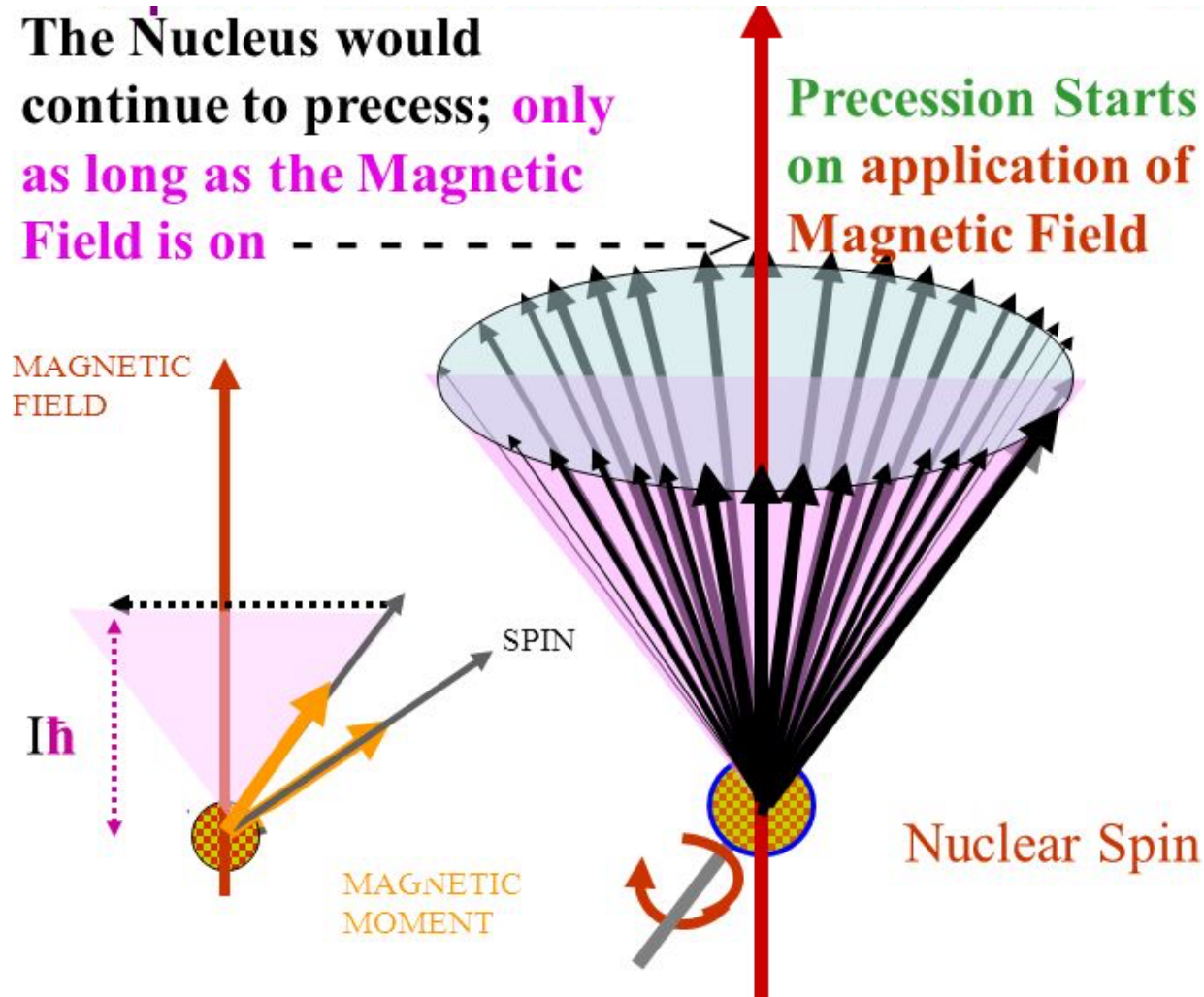
$$\vec{\mu} = \frac{g_l \mu_b}{\hbar} \vec{L}$$

Spin precession of a
charged particle



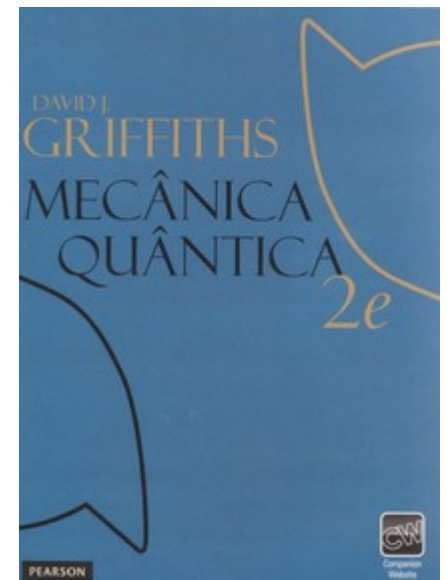
The Nucleus would
continue to precess; **only**
as long as the Magnetic
Field is on - - - - -

Precession Starts
on application of
Magnetic Field



Isto é uma descrição clássica !
...poderia ser um pião...

Como fazer um estudo quântico
da precessão do spin ?



$$U = -\vec{\mu}_s \cdot \vec{B}$$

$$\left\{ \begin{array}{l} \vec{B} = B_0 \hat{k} \\ \vec{\mu}_s = \frac{g_s \mu_b}{\hbar} \vec{S} = \gamma \vec{S} \end{array} \right.$$

$$U = -\gamma \vec{S} \cdot \vec{B} \quad \gamma \text{ é a razão giromagnética}$$

Função Hamiltoniana :

$$H = T + U = \text{e. cinética} + \text{e. potencial}$$

$$H = -\gamma \vec{S} \cdot \vec{B} = -\gamma B_0 S_z$$

Quantização : funções viram operadores

$$H = -\gamma \vec{S} \cdot \vec{B} = -\gamma B_0 S_z$$

$$H \rightarrow \hat{H} \quad \text{Operador Hamiltoniano}$$

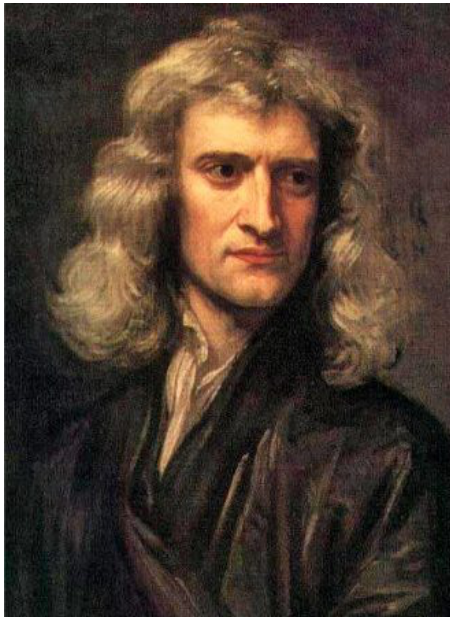
$$S_z \rightarrow \hat{S}_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\hat{H} = -\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Mecânica clássica

2ª lei de Newton

$$\frac{d^2 x}{dt^2} = - \frac{1}{m} \frac{\partial V}{\partial x}$$



Mecânica quântica

Equação de Schrödinger

$$i \hbar \frac{\partial \Psi}{\partial t} = - \frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V \Psi$$

$$\hat{H} = - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V$$

(Hamiltoniano)

$$i \hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$$

$$i \hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$$

Ansatz :

$$\Psi(t) = \psi_s e^{-i(E t)/\hbar} = |\chi\rangle e^{-i(E t)/\hbar}$$



$$\hat{H} |\chi\rangle = E |\chi\rangle$$

Já sabemos que só existem dois autoestados de spin:

$$|\chi\rangle = |\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad |\chi\rangle = |\downarrow\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Aplicando ao estado com spin pra cima

$$\hat{H} = -\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad |\chi\rangle = |\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\hat{H} |\uparrow\rangle = -\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = -\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\hat{H} |\uparrow\rangle = -\gamma B_0 \frac{\hbar}{2} |\uparrow\rangle$$

$$\hat{H} |\uparrow\rangle = E_+ |\uparrow\rangle$$

$$E_+ = -\gamma B_0 \frac{\hbar}{2}$$

Aplicando ao estado com spin para baixo

$$\hat{H} = -\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad |\chi\rangle = |\downarrow\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\hat{H} |\downarrow\rangle = -\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = +\gamma B_0 \frac{\hbar}{2} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\hat{H} |\downarrow\rangle = +\gamma B_0 \frac{\hbar}{2} |\downarrow\rangle$$

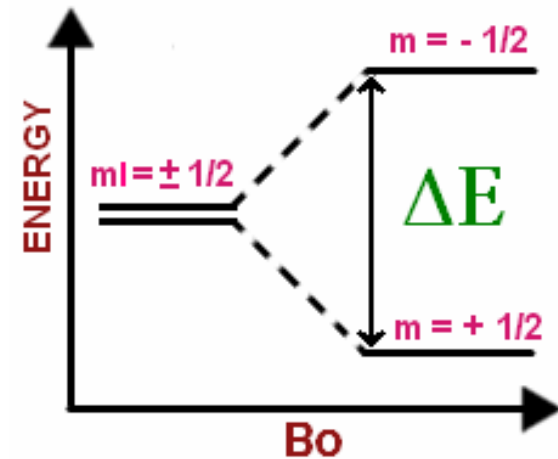
$$\hat{H} |\downarrow\rangle = E_- |\downarrow\rangle$$

$$E_- = +\gamma B_0 \frac{\hbar}{2}$$

$$\Psi(t) = |\chi\rangle e^{-i(E t)/\hbar}$$

Encontramos dois estados :

$$\left\{ \begin{array}{l} \Psi_+(t) = |\uparrow\rangle e^{-i(E_+ t)/\hbar} \\ \Psi_-(t) = |\downarrow\rangle e^{-i(E_- t)/\hbar} \end{array} \right.$$



Solução geral :

$$\Psi(t) = a \Psi_+ + b \Psi_- = a |\uparrow\rangle e^{-i(E_+ t)/\hbar} + b |\downarrow\rangle e^{-i(E_- t)/\hbar}$$

$$|a|^2 + |b|^2 = 1 \quad \longrightarrow \quad a = \cos\left(\frac{\alpha}{2}\right) \quad b = \sin\left(\frac{\alpha}{2}\right)$$

$$\Psi(t) = \cos\left(\frac{\alpha}{2}\right) \begin{pmatrix} 1 \\ 0 \end{pmatrix} e^{-i(E_+ t)/\hbar} + \sin\left(\frac{\alpha}{2}\right) \begin{pmatrix} 0 \\ 1 \end{pmatrix} e^{-i(E_- t)/\hbar}$$

$$\Psi(t) = \begin{pmatrix} \cos\left(\frac{\alpha}{2}\right) e^{-i(E_+ t)/\hbar} \\ \sin\left(\frac{\alpha}{2}\right) e^{-i(E_- t)/\hbar} \end{pmatrix} \quad \begin{aligned} E_+ &= -\gamma B_0 \frac{\hbar}{2} \\ E_- &= +\gamma B_0 \frac{\hbar}{2} \end{aligned}$$

$$\Psi(t) = \begin{pmatrix} \cos\left(\frac{\alpha}{2}\right) e^{+i(\gamma B_0 t)/2} \\ \sin\left(\frac{\alpha}{2}\right) e^{-i(\gamma B_0 t)/2} \end{pmatrix}$$

$$\Psi^\dagger(t) = \left(\cos\left(\frac{\alpha}{2}\right) e^{-i(\gamma B_0 t)/2} \quad \sin\left(\frac{\alpha}{2}\right) e^{+i(\gamma B_0 t)/2} \right)$$

Valor esperado de S_x

$$\langle \hat{S}_x \rangle = \langle \Psi | \hat{S}_x | \Psi \rangle = \Psi^\dagger \hat{S}_x \Psi$$

$$= \begin{pmatrix} \cos\left(\frac{\alpha}{2}\right) e^{-i(\gamma B_0 t)/2} & \sin\left(\frac{\alpha}{2}\right) e^{+i(\gamma B_0 t)/2} \end{pmatrix} \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \cos\left(\frac{\alpha}{2}\right) e^{+i(\gamma B_0 t)/2} \\ \sin\left(\frac{\alpha}{2}\right) e^{-i(\gamma B_0 t)/2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos\left(\frac{\alpha}{2}\right) e^{-i(\gamma B_0 t)/2} & \sin\left(\frac{\alpha}{2}\right) e^{+i(\gamma B_0 t)/2} \end{pmatrix} \begin{pmatrix} \sin\left(\frac{\alpha}{2}\right) e^{-i(\gamma B_0 t)/2} \\ \cos\left(\frac{\alpha}{2}\right) e^{+i(\gamma B_0 t)/2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \left(\cos\frac{\alpha}{2} \sin\frac{\alpha}{2} e^{-i\gamma B_0 t} + \sin\frac{\alpha}{2} \cos\frac{\alpha}{2} e^{+i\gamma B_0 t} \right)$$

$$= \frac{\hbar}{2} \cos\frac{\alpha}{2} \sin\frac{\alpha}{2} \left(e^{-i\gamma B_0 t} + e^{+i\gamma B_0 t} \right)$$

$$= \frac{\hbar}{2} \cos \frac{\alpha}{2} \sin \frac{\alpha}{2} \left(e^{-i \gamma B_0 t} + e^{+i \gamma B_0 t} \right)$$

$$= \frac{\hbar}{2} \frac{1}{2} \sin \alpha \ 2 \cos (\gamma B_0 t)$$

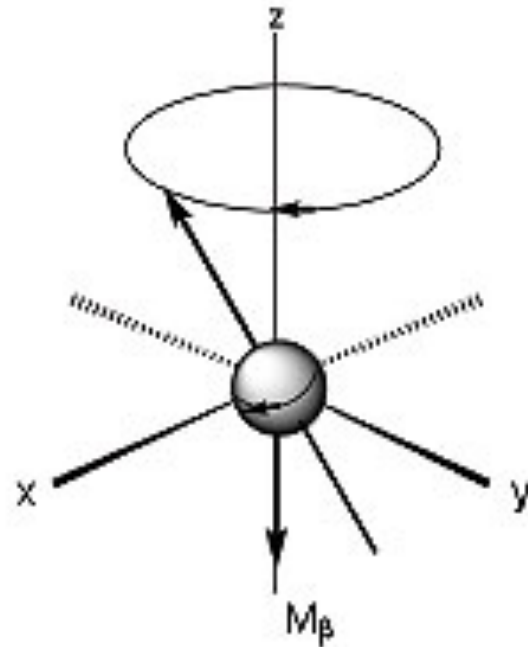
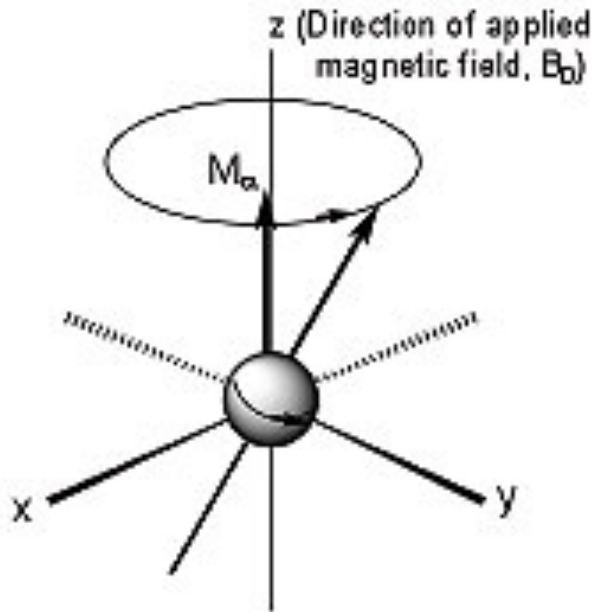
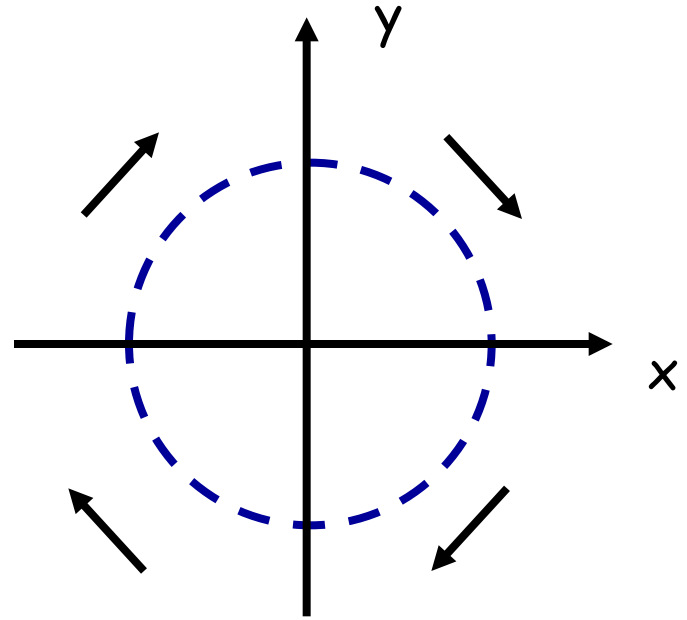
$$\left\{ \begin{array}{l} \langle \hat{S}_x \rangle = \frac{\hbar}{2} \frac{1}{2} \sin \alpha \ 2 \cos (\gamma B_0 t) \\ \langle \hat{S}_y \rangle = - \frac{\hbar}{2} \frac{1}{2} \sin \alpha \ 2 \sin (\gamma B_0 t) \end{array} \right.$$

(Exercício)

$$\omega = \gamma B_0 = \frac{g_s \mu_b B_0}{\hbar}$$

frequência de Larmor

$$\begin{cases} \langle \hat{S}_x \rangle = \frac{\hbar}{2} \frac{1}{2} \sin \alpha \cos (\gamma B_0 t) \\ \langle \hat{S}_y \rangle = -\frac{\hbar}{2} \frac{1}{2} \sin \alpha \sin (\gamma B_0 t) \end{cases}$$

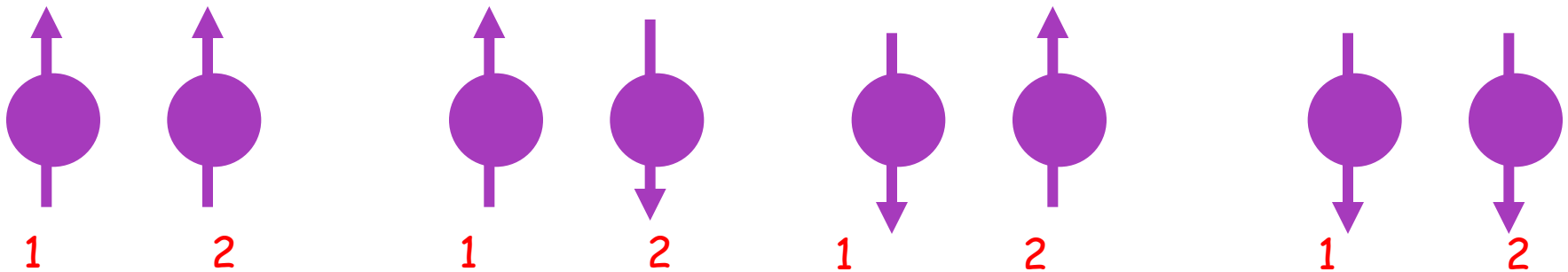


Adição de spins

Vamos somar os spins de duas partículas de spin 1/2

Por exemplo: spin total do átomo de hidrogênio

Projeções na direção z



Classicamente somamos os dois vetores de spin:

$$\vec{S} = \vec{S}^{(1)} + \vec{S}^{(2)}$$

Função de onda (autoestado) de spin de duas partículas :

$$|\chi\rangle = |\chi_1\rangle |\chi_2\rangle = \chi_1 \chi_2$$

Autovalor de S_z no estado composto :

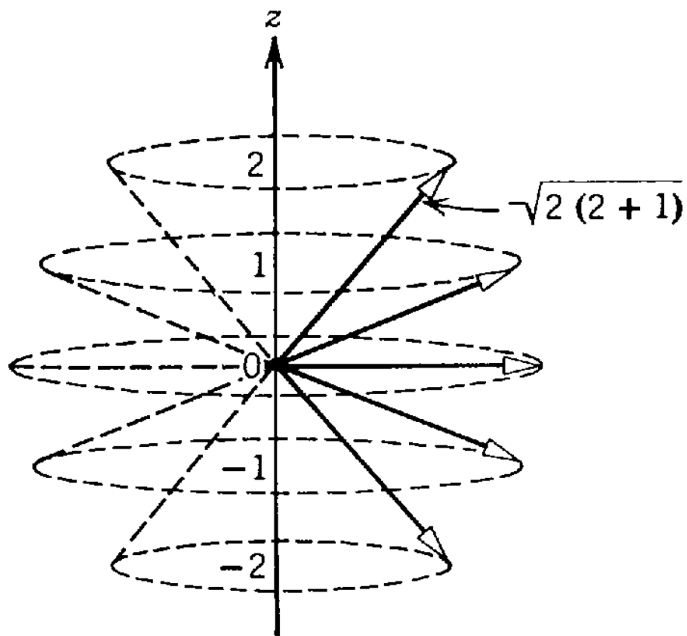
$$\vec{S} = \vec{S}^{(1)} + \vec{S}^{(2)} \quad \longrightarrow \quad \hat{S} = \hat{S}^{(1)} + \hat{S}^{(2)}$$

(operadores)

$$\hat{S}_z = \hat{S}_z^{(1)} + \hat{S}_z^{(2)} \quad \longrightarrow \quad S_z = S_z^{(1)} + S_z^{(2)}$$

$$S_z \chi_1 \chi_2 = (S_z^{(1)} + S_z^{(2)}) \chi_1 \chi_2$$



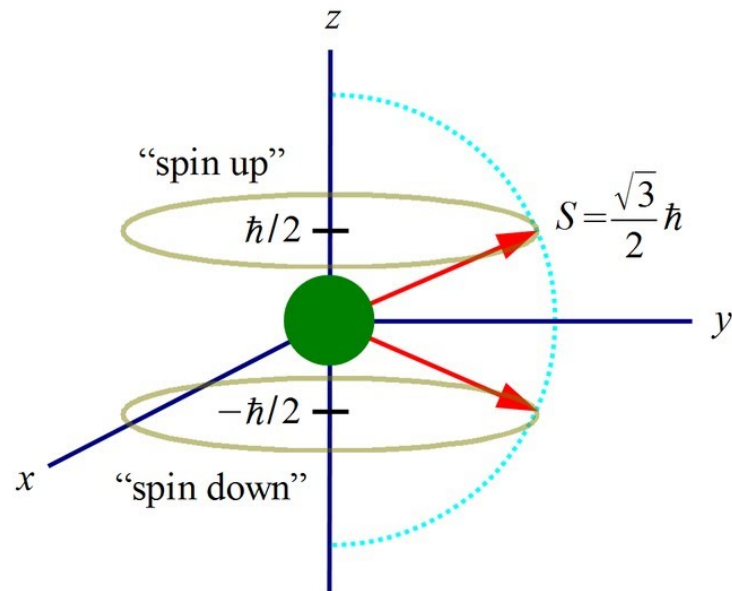


$$L = \sqrt{l(l+1)} \hbar$$

$$L_z = m_l \hbar$$

$$-l \leq m_l \leq l$$

$$l = 0, 1, 2, \dots$$



$$S = \sqrt{s(s+1)} \hbar$$

$$S_z = m_s \hbar$$

$$s = \frac{1}{2} \quad m_s = -\frac{1}{2}, +\frac{1}{2}$$



$$\hat{S}_z \psi = m_s \hbar \psi$$

$$S_z \chi_1 \chi_2 = (S_z^{(1)} + S_z^{(2)}) \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = S_z^{(1)} \chi_1 \chi_2 + S_z^{(2)} \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = m_1 \hbar \chi_1 \chi_2 + m_2 \hbar \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = (m_1 + m_2) \hbar \chi_1 \chi_2$$

$$S_z \chi_1 \chi_2 = m \hbar \chi_1 \chi_2$$

$\uparrow \uparrow$	$m = 1$	$\downarrow \uparrow$	$m = 0$
$\uparrow \downarrow$	$m = 0$	$\downarrow \downarrow$	$m = -1$

