

300303 - 2

Electromagnetismo II
Segunda Prova - Noturno
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1.

$$a) \vec{R} = \vec{r} - \vec{r}_q = d\hat{e}_x - z_0 \omega(\omega t_r) \hat{e}_z \quad \therefore R = [d^2 + z_0^2 \omega^2 (\omega t_r)^2]^{\frac{1}{2}}$$

$$t_r = t - \frac{R}{c} \quad \therefore c(t - t_r) = [d^2 + z_0^2 \omega^2 (\omega t_r)^2]^{\frac{1}{2}} \quad \therefore \boxed{(t - t_r)^2 - \frac{z_0^2}{c^2} \omega^2 (\omega t_r)^2 = \frac{d^2}{c^2}}$$

$$b) \vec{v} = \frac{d\vec{r}_q}{dt_r} = -z_0 \omega \sin(\omega t_r) \hat{e}_z; \quad \beta = \frac{v}{c} = \beta \ll 1 \quad \Rightarrow \quad \boxed{\frac{z_0 \omega}{c} \ll 1}$$

$$c) \beta \ll 1 \quad \Rightarrow \quad \frac{dP}{d\Omega} = \frac{1}{c^3} \frac{q^2}{16\pi^2 \epsilon_0} [\hat{n} \times (\hat{n} \times \vec{a})]^2$$

$$\frac{z_0}{r} \ll 1 \quad \Rightarrow \quad \frac{d\vec{r}_q}{dt_r} \approx \frac{\vec{r} - \vec{r}_q}{|\vec{r} - \vec{r}_q|} \approx \hat{e}_r$$

$$\therefore [\hat{n} \times (\hat{n} \times \vec{a})] = \hat{n} \times [\hat{e}_r \times a \hat{e}_z] = -a \sin\theta (\hat{n} \times \hat{e}_\phi);$$

$$\therefore [\hat{n} \times (\hat{n} \times \vec{a})]^2 = a^2 \sin^2\theta \quad ; \quad a = -z_0 \omega^2 \cos^2(\omega t_r)$$

$$\therefore \frac{dP}{d\Omega} = \frac{1}{c^3} \frac{q^2}{16\pi^2 \epsilon_0} z_0^2 \omega^4 \cos^2(\omega t_r) \sin^2\theta$$

$$\langle \frac{dP}{d\Omega} \rangle = \frac{q^2 z_0^2 \omega^4}{16\pi^2 \epsilon_0 c^3} \sin^2\theta \langle \cos^2(\omega t_r) \rangle \quad \therefore \quad \boxed{\langle \frac{dP}{d\Omega} \rangle = \frac{q^2 z_0^2 \omega^4}{32\pi^2 \epsilon_0 c^3} \sin^2\theta}$$

$$d) \langle P \rangle = \int_0^{2\pi} d\varphi \int_0^\pi \sin\theta d\theta \langle \frac{dP}{d\Omega} \rangle = \frac{q^2 z_0^2 \omega^4}{32\pi^2 \epsilon_0 c^3} \cdot 2\pi \cdot \int_0^\pi \sin^3\theta d\theta$$

$$\therefore \boxed{\langle P \rangle = \frac{q^2 z_0^2 \omega^4}{12\pi^2 \epsilon_0 c^3}}$$

$$2. \phi(\vec{r}, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi_{\omega}(\vec{r}) e^{-i\omega t} d\omega$$

$$\phi(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{1}{|\vec{r}-\vec{r}'|} \left[\frac{1}{\sqrt{2\pi}} \int \rho_{\omega}(\vec{r}') e^{\pm i\mathbf{k}|\vec{r}-\vec{r}'|} e^{-i\omega t} d\omega \right] d\tau'$$

$$= \frac{1}{4\pi\epsilon_0} \int \frac{d\tau'}{|\vec{r}-\vec{r}'|} \left\{ \frac{1}{\sqrt{2\pi}} \int \rho_{\omega}(\vec{r}') e^{-i\omega \left[t - \frac{|\vec{r}-\vec{r}'|}{c} \right]} d\omega \right\}$$

Mas $\frac{1}{\sqrt{2\pi}} \int \rho_{\omega}(\vec{r}') e^{-i\omega \left[t - \frac{|\vec{r}-\vec{r}'|}{c} \right]} d\omega = \rho(\vec{r}', t - \frac{|\vec{r}-\vec{r}'|}{c})$

Escolhemos o sinal (-) porque o argumento resultante representa o tempo retardado, t_r , ou seja, o tempo t menos o tempo que levou para a informação sobre o valor instantâneo da fonte se propagar de \vec{r}' a \vec{r} , com a velocidade da luz. Portanto

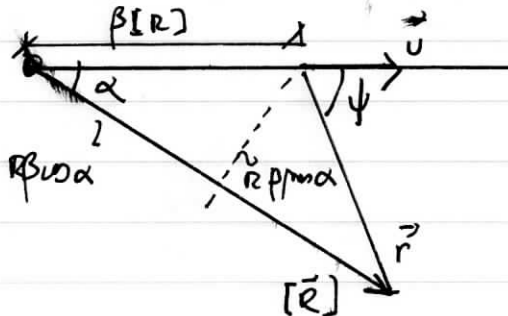
$$\boxed{\phi(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{|\vec{r}-\vec{r}'|} d\tau'; \quad t_r = t - \frac{|\vec{r}-\vec{r}'|}{c}}$$

3.

a) $\vec{a} = 0 \Rightarrow \vec{E}(\vec{r}, t) = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{s^3} (\vec{R} - R\vec{\beta})(1-\beta^2) \right]_{ret}$

$$s = R - \vec{R} \cdot \vec{\beta}$$

$$s = R - \vec{R} \cdot \vec{\beta} = R - R\beta \omega \alpha$$



Mas $R\beta\omega\alpha = R - \sqrt{r^2 - R^2\beta^2} \mu\beta\alpha$

$$\therefore s = R - (R - \sqrt{r^2 - R^2\beta^2} \mu\beta\alpha) = \sqrt{r^2 - R^2\beta^2} \mu\beta\alpha$$

Mas $R\mu\beta\alpha = r\mu\psi$ $\therefore s = \sqrt{r^2 - R^2\beta^2} \mu\psi$

$$\therefore s = r \sqrt{1 - \beta^2} \mu\psi$$

Por outro lado, da figura temos que

$$\vec{r} = R\vec{\beta} + \vec{r} \quad \therefore \vec{r} - \vec{\beta}R = \vec{r};$$

assim

$$\vec{E}(\vec{r}, t) = \frac{q\vec{r}}{4\pi\epsilon_0 r^3} \frac{1-\beta^2}{[1-\beta^2 \sin^2\psi]^{3/2}}$$

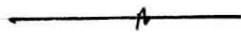
$$b) \quad \vec{B} = \frac{\vec{v}}{c^2} \times \vec{E}; \quad \vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B} = \frac{1}{\mu_0 c^2} (\vec{E} \times (\vec{v} \times \vec{E}))$$

$$\therefore \vec{S} = \frac{1}{\mu_0 c^2} [E^2 \vec{v} - (\vec{E} \cdot \vec{v}) \vec{E}] = \frac{1}{\mu_0 c^2} [c^2 \vec{v} - E v \cos\psi \vec{E}]$$

$$\therefore \vec{S} = \epsilon_0 [E^2 \vec{v} - E^2 v \cos\psi (\cos\psi \hat{v} + \sin\psi \hat{e}_\perp)]$$

$$\therefore \vec{S} = \epsilon_0 E^2 [\sin^2\psi \vec{v} - v \sin\psi \cos\psi \hat{e}_\perp]$$

onde \hat{e}_\perp é o vetor normal ao plano formado por \vec{v} e \vec{r} .



4.

$$a) \quad [R] \times \vec{E}_v = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{s^3} (-R (\vec{r} \times \vec{\beta})) (1-\beta^2) \right] = \frac{qR}{4\pi\epsilon_0} \left[\frac{\vec{\beta} \times \vec{r}}{s^2} (1-\beta^2) \right]$$

$$\therefore [R] \times \vec{E}_v = [R] c \vec{B}_v; \quad \vec{B}_v = \frac{[R] \alpha \vec{E}}{[R] c}$$

$$[R] \times \vec{E}_a = \frac{q}{4\pi\epsilon_0} \left[\frac{R}{2s^3} (\vec{r} \cdot \vec{a}) (\vec{\beta} \times \vec{r}) - \frac{R \vec{r} \times \vec{a}}{2s^2} \right] = [R] c \vec{B}_a \quad \checkmark$$

$$b) \quad \frac{\vec{r}}{R} \times \left\{ \vec{r} \times ((\vec{r} - R\vec{\beta}) \times \vec{a}) \right\} = \frac{\vec{r}}{R} \times [(\vec{r} \cdot \vec{a})(\vec{r} - R\vec{\beta}) - (R^2 - R\vec{r} \cdot \vec{\beta}) \vec{a}]$$

$$= \frac{\vec{r}}{R} \times [(\vec{r} \cdot \vec{a})(\vec{r} - R\vec{\beta}) - R s \vec{a}] = \frac{1}{R} [(\vec{r} \cdot \vec{a})(\vec{\beta} \times \vec{r}) - R s \vec{r} \times \vec{a}]$$

$$\therefore \vec{B}_a(\vec{r}, t) = \frac{q}{4\pi\epsilon_0 c} \left[\frac{\vec{r} \cdot \vec{a}}{2s^3} (\vec{\beta} \times \vec{r}) - \frac{\vec{r} \times \vec{a}}{2s^2} \right]$$

$$[\vec{r}^0] \cdot \vec{E}_a = \frac{q}{4\pi\epsilon_0} \frac{1}{(2s)^2} \left[\frac{1}{s} R S (\vec{r} \cdot \vec{a}) - R (\vec{r} \cdot \vec{a}) \right] = 0$$