

Prova Substitutiva
 Gabarito

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1. a) Como o fluxo através do circuito formado pela barra e resistor começa a diminuir, a corrente tem que circular no sentido de que o campo por ela produzido aumente o fluxo. Portanto a corrente irá circular no sentido indicado na figura.

$$b) \quad \mathcal{E} = - \int \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} + \oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell}$$

$$\vec{v} = 0; \quad \frac{\partial \vec{B}}{\partial t} = - \frac{B_0}{T} \hat{z} \quad \therefore \quad \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} = - \frac{B_0}{T} dS \quad \therefore \quad \mathcal{E} = \frac{B_0}{T} a z_0$$

$$\therefore \quad \boxed{I = \frac{B_0}{RT} a z_0}$$

$$\vec{f} = I \vec{\ell} \times \vec{B} = I a (-\hat{x}) \times B_0 \hat{z} \quad \therefore \quad \vec{f} = I a B_0 \hat{y} = \frac{B_0^2 a^2 z_0}{RT} \hat{y}$$

Para que a barra suba, esta força tem que ser maior que a da gravidade;

$$|\vec{f}| \geq |\vec{f}_g| \quad \therefore \quad \boxed{\frac{B_0^2 a^2 z_0}{RT} \geq mg}$$

$$c) \quad \mathcal{E} = - \int \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} + \oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = + \frac{B_0}{T} a (t - z_0) + \int_0^a (\hat{v} \hat{y} \times B_0 \hat{z}) \cdot (-dx \hat{x})$$

$$\therefore \quad \mathcal{E} = + \frac{B_0}{T} a (t - z_0) - v B_0 (1 - \frac{t}{T}) a \quad \therefore \quad \boxed{I = \frac{B_0 a}{RT} [t - z_0 - v(T - t)]}$$

$$d) \quad \vec{f} = I \vec{\ell} \times \vec{B} = \frac{B_0^2 a^2}{RT} [t - z_0 - v(T - t)] \hat{y}$$

$$m \frac{d\vec{v}}{dt} = \vec{f} + \vec{f}_g \quad \therefore$$

$$\boxed{m \frac{dv}{dt} = \frac{B_0 a}{RT} [t - z_0 - v(T - t)] - mg}$$

2. a) $\vec{B}(\vec{r}) = \frac{\mu_0 I}{4\pi} \int \frac{d\vec{r}' \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$

$\vec{r} = 0; \vec{r}' = x' \hat{e}_x + a \hat{e}_y; d\vec{r}' = dx' \hat{e}_x$

$\therefore d\vec{r}' \times (\vec{r} - \vec{r}') = +a \hat{e}_z \therefore \vec{B}_1(0) = \frac{\mu_0 I a}{4\pi} \int_{-\infty}^0 \frac{dx'}{[a^2 + x'^2]^{3/2}} \hat{e}_z$

$\therefore \vec{B}_1(0) = \frac{\mu_0 I a}{4\pi} \left[\frac{x'}{a^2 \sqrt{a^2 + x'^2}} \right]_{-\infty}^0 \hat{e}_z$

$\therefore \boxed{\vec{B}_1(0) = \frac{\mu_0 I}{4\pi a} \hat{e}_z}$

b) $\vec{r} = 0; \vec{r}' = a \hat{e}_r; d\vec{r}' = a d\theta' \hat{e}_\theta$

$\therefore d\vec{r}' \times (\vec{r} - \vec{r}') = -a^2 d\theta' \hat{e}_\theta \times \hat{e}_r = a^2 d\theta' \hat{e}_z$

$\therefore \vec{B}_2(0) = \frac{\mu_0 I a^2}{4\pi} \int_{-\pi/2}^0 \frac{d\theta'}{a^3} \hat{e}_z \therefore \boxed{\vec{B}_2(0) = \frac{\mu_0 I}{8a} \hat{e}_z}$

c) $\vec{r} = 0; \vec{r}' = a \hat{e}_x + y' \hat{e}_y; d\vec{r}' = dy' \hat{e}_y$

$\therefore d\vec{r}' \times (\vec{r} - \vec{r}') = a dy' \hat{e}_z \therefore \vec{B}_3 = \frac{\mu_0 I a}{4\pi} \int_0^\infty \frac{dy'}{[a^2 + y'^2]^{3/2}} \therefore \boxed{\vec{B}_3 = \frac{\mu_0 I}{4\pi a} \hat{e}_z}$

$\boxed{\vec{B}_+ = \frac{\mu_0 I}{2\pi a} \left[1 + \frac{\pi}{4} \right] \hat{e}_z}$

3. a) $\nabla \cdot \vec{B} = \nabla \cdot (\nabla \times \vec{A}) = 0$

$\nabla \times \vec{E} = -\nabla \times \frac{\partial \vec{A}}{\partial t} \therefore \nabla \times \left(\vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0 \therefore \vec{E} = -\nabla \phi - \frac{\partial \vec{A}}{\partial t}$

b) $\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \therefore \nabla \times (\nabla \times \vec{B}) = \mu_0 \vec{J} + \mu_0 \epsilon_0 \left[-\nabla \frac{\partial \phi}{\partial t} - \frac{\partial^2 \vec{A}}{\partial t^2} \right]$

$\therefore \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \left[-\nabla \frac{\partial \phi}{\partial t} - \frac{\partial^2 \vec{A}}{\partial t^2} \right]$

$$3. a) \quad \nabla \cdot \vec{B} = \nabla \cdot (\nabla \times \vec{A}) = 0$$

$$\nabla \times \vec{E} = -\nabla \times \frac{\partial \vec{A}}{\partial t} \therefore \nabla \times \left(\vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0 \therefore \vec{E} = -\nabla \phi - \frac{\partial \vec{A}}{\partial t} \quad (3)$$

$$b) \quad \nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \therefore \nabla \times (\nabla \times \vec{A}) = \mu_0 \vec{J} + \mu_0 \epsilon_0 \left[-\nabla \frac{\partial \phi}{\partial t} - \frac{\partial^2 \vec{A}}{\partial t^2} \right]$$

$$\therefore \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \left[-\nabla \frac{\partial \phi}{\partial t} - \frac{\partial^2 \vec{A}}{\partial t^2} \right]$$

$$\therefore \nabla^2 \vec{A} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J} + \nabla \left[\nabla \cdot \vec{A} + \mu_0 \epsilon_0 \frac{\partial \phi}{\partial t} \right]$$

$$\therefore \text{Impondo } \nabla \cdot \vec{A} + \mu_0 \epsilon_0 \frac{\partial \phi}{\partial t} = 0 \Rightarrow \nabla^2 \vec{A} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J}$$

c) Para fontes constantes no tempo, o potencial vetor é dado por

$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{|\vec{r} - \vec{r}'|} d\tau'$$

Se a fonte variar no tempo, a informação de que ela está variando na posição $\vec{r}'$ se propaga com a velocidade da luz até a posição $\vec{r}$, levando um tempo $|\vec{r} - \vec{r}'|/c$. Portanto, o potencial medido na posição $\vec{r}$, no instante t , é correspondente ao valor que a fonte tinha na posição $\vec{r}'$ no instante retardado $t - \frac{|\vec{r} - \vec{r}'|}{c}$.

$$4. a) \quad R_3 + T_3 = 1; R_3 = \frac{1}{3} T_3 \Rightarrow R_3 = \frac{1}{4} \quad T_3 = \frac{1}{2} = \frac{n_1 \sin \theta_1 - n_2 \sin \theta_2}{n_1 \sin \theta_1 + n_2 \sin \theta_2}$$

$$\therefore n_1 \sin \theta_1 = 3 n_2 \sin \theta_2; \left(\frac{n_1}{3 n_2} \right)^2 (1 - \mu^2 \theta_1) = 1 - \left(\frac{n_1}{n_2} \right)^2 \mu^2 \theta_1$$

$$\therefore \mu^2 \theta_1 = \frac{1 - \frac{n_1^2}{9 n_2^2}}{\frac{8 n_1^2}{9 n_2^2}} = \frac{3}{4} \quad \mu \theta_1 = \frac{\sqrt{3}}{2} \quad \boxed{\theta_1 = \frac{\pi}{3}}$$

$$b) \quad E'_k = r_3 |E_k|$$

$$4.b) \bar{E}_{1s} = r_s \bar{E}_1 \quad \therefore \boxed{\bar{E}_{1s} = \frac{\bar{E}_{1s}}{2}}; \quad \bar{E}_{2s} = t_s \bar{E}_{1s}$$

$$t_s = \frac{2n_1 \cos \theta_1}{n_1 \cos \theta_1 + n_2 \cos \theta_2} = \frac{2n_1 \cos \theta_1}{n_1 \cos \theta_1 \left(1 + \frac{n_2 \cos \theta_2}{n_1 \cos \theta_1}\right)}$$

Como do item a) temos:

$$\frac{n_2 \cos \theta_2}{n_1 \cos \theta_1} = \frac{1}{3}$$

$$\therefore t_s = \frac{2}{1 + \frac{1}{3}} = \frac{2}{4/3} \Rightarrow \boxed{t_s = \frac{3}{2}}$$

Daí:

$$\boxed{\bar{E}_{2s} = \frac{3}{2} \bar{E}_{1s}}$$